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HEAT TRANSFER ANALYSIS FOR  
UNSTEADY HIGH VELOCITY PIPE FLOW

CHING JEN CHEN

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JUL 14 1977  
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Prepared By

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→ inner surface of the pipe by inverting the temperature response measured by an interior probe close to the heated surface. The refinement is achieved by using the double precision format in the program and adapting the dimensionless formulation. The third is to study the inversion solution for a large time duration of a time dependent surface heat flux. The solution is obtained by the method of Laplace transformation with the convolution integral.

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HEAT TRANSFER ANALYSIS FOR UNSTEADY  
HIGH VELOCITY PIPE FLOW

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## ABSTRACT

This report consists of three parts. The first is the analysis for minimization of the temperature distortion due to the thermocouple cavity. The error is minimized or reduced to zero by optimizing the combination of cavity diameter and depth and the thermocouple transport properties and size. The second is the refinement of the presently available computer program for prediction of the surface temperature and heat flux at the inner surface of the pipe by inverting the temperature response measured by an interior probe close to the heated surface. The refinement is achieved by using the double precision format in the program and adapting the dimensionless formulation. The third is to study the inversion solution for a large time duration of a time dependent surface heat flux. The solution is obtained by the method of Laplace transformation with the convolution integral.

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## FOREWORD

This is the final report of the project entitled "Heat Transfer Analysis for Unsteady High Velocity Pipe Flow." The project was funded by Army Research Office Grant DAA-G29-76-G-0123 for a six month period from January 1976 to June 1976.

The work done for project is reported here in three parts. The first is the analysis for minimization of the temperature distortion due to the thermocouple cavity. The error is minimized or reduced to zero by optimizing a combination of cavity diameter and depth and the thermocouple material and size. The second is the refinement of the presently available computer programs for prediction of the surface temperature and heat flux at the inner surface of the pipe by inverting the temperature measured by an interior probe close to the heated surface. The refinement is achieved by using the double precision format in the program and adapting the dimensionless formulation. The third is to study the inversion solution for a large time duration of a time dependent surface heat flux. The solution is obtained by the method of Laplace transform and the convolution integral. Each of the above three subjects is reported as a part of the present report.

The investigator would like to thank Mr. D. M. Thomsen of General Rodman Laboratory, Rock Island Arsenal, Rock Island, Illinois for his constant participation in the course of the research.

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## NOMENCLATURE

c	specific heat
$d_t$	diameter of thermocouple
d	diameter of cavity
D	thickness of the disk
Q	constant heat flux
T	temperature
t	real time
x	X/D
X, Y	coordinate along and normal to the heated surface
y	Y/D
Error %	$[\theta \text{ (at cavity base)} - \theta \text{ (at edge of disk)}] \times 100$ , or distorted temperature divided by $QD/\kappa_1$

## Greek letters

$\epsilon$	distance between the base of cavity to the heated surface
$\theta$	dimensionless temperature $T\kappa_1/QD$
$\tau$	dimensionless time
$\alpha$	thermal diffusivity
$\kappa$	thermal conductivity
$\rho$	density

**subscripts**

- 1        the disk material
- 2        insulation material
- 3        thermocouple material
- t        thermocouple

PART I

ON MINIMIZATION OF TEMPERATURE DISTORTION  
IN THE THERMOCOUPLE CAVITY

## SUMMARY - PART I

When a thermocouple is embedded in a solid to measure temperature a distortion of temperature is created because of the thermocouple cavity. For a given solid under measurement the temperature distortion at the base of the thermocouple cavity can be minimized by a proper choice of the thermocouple size and its material, the cavity diameter and the depth of the cavity. This optimum combination is solved in this study by a finite element analysis for the case that the thermocouple cavity is drilled into the center of a disk. The disk, initially at a uniform temperature, is then heated with a constant heat flux on one surface and is insulated on the other surface. The calculated result covers a range of thermal conductivity and diffusivity for most commonly used thermocouples. For ease in practical applications a simple formula for determining the optimum ratio of the thermocouple diameter to that of the cavity is given as a function of density - specific heat product and thermal conductivity of solid material, thermocouple, and insulating material.

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## INTRODUCTION

A direct measurement of transient surface temperature and heat flux is often difficult. For example, a surface involves two modes of heat transfer, say, radiative and convective heat transfer. In this case if the measuring probe has a different radiative property than that of the surface, erroneous measurements will result. A piston or projectile sliding over the cylindrical surface is another case where the direct measurement at the surface is difficult. A surface involving melting or ablation is also difficult to make direct measurement. Therefore, indirect estimation by inverting the temperature history inside the heat conducting solid as measured by a thermocouple is often used for prediction of the surface temperature and heat flux. Beck [1], Hernning and Parker [2], Frank [3], Imber and Khan [4], Stolz [5], Chen and Thomsen [6] have developed different inversion solutions for this purpose. All of these solutions assumed that the cavity drilled into the solid does not distort the true temperature distribution. Thus, it is important that the temperature measurement by an interior probe is accurate and involves least distortion or error. Theoretically Beck [7], Masters and Stein [8], Burnett [9], and Chen and Li [10] studied the distortion of the temperature field in the presence of a thermocouple and its cavity. Experimentally Chen and Danh [11] showed that appreciable distortion, say 10%, of temperature field may exist for a normal implant of the thermocouple into a solid body. From studies of Chen and Li [10] and Beck [7] they found that with a proper combination of the thermocouple cavity diameter, its depth, and the thermocouple material and its diameter,

the distortion of temperature field with respect to space or time can be minimized if not eliminated. In this report we study the optimum combination of these parameters such that at a given situation one knows what is the best combination to use and what is the magnitude of the temperature distortion.

#### FORMULATION OF PROBLEM

In the present study we consider a disk depicted in Figure 1 which has a thickness  $D$  and is drilled a cavity of a diameter to a depth of  $\epsilon$  distance from the heating surface. The heat flux  $Q$  is assumed to be constant and the upper surface of the disk is assumed to be insulated. A thermocouple of a diameter  $d_t$  is then welded on the cavity base. Furthermore, the disk may be thought to approximate a cut out from a hollow cylinder if the radius of the disk is small compared with the radius of the cylinder. The diameter of the disk is chosen to be  $D$  where the temperature distortion due to the thermocouple cavity becomes negligible. For this to be true one needs to restrict the ratio of cavity diameter to the disk diameter  $d/2D$  be small. The unfilled cavity can be air or insulating material.

The basic idea to minimize or to eliminate the temperature distortion is based on a proper choice of the thermocouple size and the material which has a higher thermal conductivity than that of the disk so as to conduct more heat away at the cavity base balancing the insulation effect of the insulator in the cavity.

Let  $X$  and  $Y$  be respectively the coordinate along and normal to the heated disk surface and the  $Y$  axis coincide with the axis of the cavity.

The thermal conductivity is assumed to be constant and the temperature distribution is axisymmetrical. The governing equations for the transient heat conduction in dimensionless form are:

for the disk (subscript "1", see Figure 2)

$$\frac{\partial \theta}{\partial \tau} = \frac{2}{\alpha_1^2} \left( \frac{\partial^2 \theta}{\partial x^2} + \frac{1}{x} \frac{\partial \theta}{\partial x} + \frac{\partial^2 \theta}{\partial y^2} \right) \quad (1)$$

for the insulating material in the cavity (subscript "2")

$$\frac{\partial \theta}{\partial \tau} = \frac{\alpha_2}{\alpha_1} \left( \frac{\partial^2 \theta}{\partial x^2} + \frac{1}{x} \frac{\partial \theta}{\partial x} + \frac{\partial^2 \theta}{\partial y^2} \right) \quad (2)$$

for the thermocouple (subscript "3")

$$\frac{\partial \theta}{\partial \tau} = \frac{\alpha_3}{\alpha_1} \left( \frac{\partial^2 \theta}{\partial x^2} + \frac{1}{x} \frac{\partial \theta}{\partial x} + \frac{\partial^2 \theta}{\partial y^2} \right) \quad (3)$$

where  $\tau = \alpha_1 t/D^2$  is the dimensionless time,  $x = X/D$  the dimensionless radial coordinate,  $y = Y/D$ , the dimensionless distance normal to the heated surface.  $\alpha_1$ ,  $\alpha_2$ , and  $\alpha_3$  are respectively the thermal diffusivity of the disk, insulating material, and the thermocouple. The dimensionless temperature  $\theta$  is defined as  $T - T_1$  where  $T$  is the temperature above the initial, uniform temperature.

The initial temperature of the disk is then

$$\theta(x, y, 0) = 0 \quad (4)$$

The boundary conditions, see Figure 2, are:

a constant heat flux at the lower surface,

$$y = 0 \quad \left. \frac{\partial \theta}{\partial y} \right|_{y=0} = 1 \quad (5)$$

the insulation at the upper surface,

$$y = 0 \quad \left. \frac{\partial \theta}{\partial y} \right|_{y=0} = 0 \quad (6)$$

the zero temperature distortion at the edge of the disk,

$$x = 1 \quad \left. \frac{\partial \theta}{\partial x} \right|_{x=1} = 0 \quad (7)$$

and the axisymmetric condition at the cavity axis

$$x = 0 \quad \left. \frac{\partial \theta}{\partial x} \right|_{x=0} = 0 \quad (8)$$

The condition (7) of the zero temperature distortion was verified by Chen and Li [10] in their early calculation when the cavity diameter is one tenth of the disk diameter. In addition to the above boundary condition the temperature and heat flux at the interface of the disk, thermocouple, and insulating material are taken to be continuous.

#### ANALYSIS

There are five parameters that can be varied for the present analysis. They are (a) the dimensionless distance from the base of the cavity to the heated surface  $\epsilon/D$ , (b) the size of the cavity  $d/D$ , (c) the ratio of the thermocouple diameter to that of the cavity,  $d_t/d$ .

(d) the thermal conductivity ratio  $\kappa_2/\kappa_1$ , and  $\kappa_3/\kappa_1$  which comes from the continuity of heat flux at interfaces. (e) the ratio of the product of density and specific heat  $\rho_3 c_3/\rho_1 c_1$  (or equivalently to the ratio of thermal diffusivity  $\alpha_3/\alpha_2$  or  $\alpha_2/\alpha_1$ ). The subscripts 1, 2 and 3 denote the disk, insulation and thermocouple materials.

Because of the complexity of the geometry and the multiplicity of material the method of finite element technique as discussed by Wilson and Nikel [12] is adapted with the aid of a computer program developed by Wilson [13]. The present problem is subdivided into finite element as required by the method, see Figure 2. The dimensionless pie section is subdivided into 121 finite elements with 12 dividing lines on both coordinates. Each element is defined by four nodal points where nodal points are denoted by intersections of the dividing lines and numbered as shown in Figure 2. The material property corresponding to each element is then assigned to the program developed by Wilson [13]. The solution at each nodal with respect to time is then obtained.

For numerical calculation three typical values of the distance from heated surface to the base of the cavity  $\epsilon/D$  are chosen to be 0.04, 0.1 and 0.2. The cavity diameter is fixed at one tenth of the disk diameter. The thermocouple to cavity diameter ratio  $d_t/d$  is made to vary 0, 0.2, 0.4, 0.6, 0.8 and 1.0. Regarding the range of the ratio of the thermal conductivity and the ratio of the product of density and specific heat we surveyed these ratios for the commonly used thermocouples in Table 1 [14] and plotted in Figure 3. Figure 3b shows the ratio of

thermocouple conductivity to that of the disk  $\kappa_3/\kappa_1$  versus the ratio of density - specific heat product  $\rho_3 c_3/\rho_1 c_1$  where the value of the conductivity is taken to be the average value between 200 and 800°K. One sees that for most practical situations the ratio of density - specific heat product  $\rho_3 c_3/\rho_1 c_1$  is approximately constant at 0.7 except when the conductivity ratio  $\kappa_3/\kappa_1$  is small. Thus the value of  $\rho_3 c_3/\rho_1 c_1$  and  $\kappa_3/\kappa_1$  for calculation are chosen, as shown in triangular symbols of Figure 3b, to cover the practical range. The corresponding value for  $\rho_2 c_2/\rho_1 c_1$  and  $\kappa_2/\kappa_1$  for the insulation material are chosen to be fixed at 0.5 and 0.005 which is a typical value for Teflon insulating material and is also approximately the order of magnitude for air.

#### RESULTS AND DISCUSSIONS

Numerical results of the calculations are presented in Tables 2 to 4 and Figures 4 to 9. The percentage error of temperature is defined as the distorted temperature divided by a reference temperature defined by  $QD/\kappa_1$ . Tables 2 to 4 give the percentage error of temperature distortion) as a function of time for different values of the parameters  $d_t/d$  (0 to 1.0),  $\kappa_3/\kappa_1$  (0.5 to 10),  $\rho_3 c_3/\rho_1 c_1$  (0.5 to 1.8) and  $\epsilon/D$  (0.02 to 0.1).

Figures 4, 5 and 6 show the three typical temperature distributions in the steel disk near the thermocouple junction for the case  $\epsilon/D = 0.06$   $d/2D = 0.1$  at the time  $\tau = 0.08$ . Figure 4 is the temperature distribution when the cavity is filled entirely with the insulation material (Teflon "2"  $\kappa_2/\kappa_1 = 0.005$ ,  $\rho_2 c_2/\rho_1 c_1 = 0.5$ ). The dimensionless isotherms

$T_{k_1}/QD$  shows the distortion of temperature distribution. One sees that the insulation effect on the heat transfer near the cavity base not only creates a much higher junction temperature of  $T_{k_1}/QD = 0.342$  than the undistorted one of 0.255 at the edge of the disk giving an error of 8.7% but also causes a hot spot at the heating surface with a higher temperature of  $T_{k_1}/QD = 0.374$  over the undistorted one of 0.311. On the other hand the temperature distribution in the insulation material is much lower than the true temperature creating a large temperature gradient at the base of the cavity. Figure 5, contrary to Figure 4, is the temperature distribution when the cavity is completely filled with thermocouple material whose thermal conductivity is ten times larger than the disk material e.g., copper versus steel. Now over-conduction of heat by the thermocouple has created a cold spot at the base of the cavity giving,  $T_{k_1}/QD$  of 0.17 versus the undistorted one of 0.255 with an error of 8.5% as well as at the heating surface with  $T_{k_1}/QD$  of 0.236 versus 0.311. The temperature distribution in the thermocouple now becomes higher than the undistorted one in the disk. By properly choosing the ratio of thermocouple diameter to that of the cavity one may minimize these distortions of the temperature response at the base of the cavity. This is shown in Figure 6 where the thermocouple diameter  $d_t$  is chosen to be 0.4 of the cavity diameter  $d$  with  $\kappa_3/\kappa_1 = 10$ , e.g. copper-steel combination. Figure 6 shows that distortions at the base of the cavity and at the heating surface are almost eliminated giving the temperature  $T_{k_1}/QD$  of 0.245 and 0.320 with respectively the error of 1% and 0.9%.

In order to examine the details of the distorted temperature response at the base of the cavity we tabulated the results in Tables 2, 3 and 4 and plotted the error percentage as function of the ratio of the thermocouple diameter to that of the cavity for different cavity depth, time, and thermal conductivity in Figures 7, 8 and 9. In these Figures the  $\rho c$  ratio ranges from 0.7 to 1.3 which covers most of the practical applications.

In the case when the  $\rho c$  ratio is equal to or less than one the errors in temperature response at the base of the cavity in these figures are all positive or overheat for  $\kappa_3/\kappa_1 \leq 1$ . This is because the thermal conductivity of the thermocouple is less than that of the disk and the heat capacity  $\rho c$  of the thermocouple is also small. Therefore, no extra conduction of heat can be achieved by the thermocouple to compensate for the blocking of the heat transfer by the insulation material in the cavity. On the other hand, if  $\kappa_3/\kappa_1 > 1$  the error of temperature varies from positive value for  $d_t/d = 0$  to some negative value as  $d_t/d$  approaches 1. Thus for  $\kappa_3/\kappa_1 > 1$  a properly chosen combination of thermocouple and insulation material can minimize the error. For example, in Figure 7 a combination of  $\kappa_3/\kappa_1 = 10$ ,  $\rho_3 c_3/\rho_1 c_1 = 0.75$ ,  $\epsilon/D = 0.1$  and  $d_t/d = 0.5$  produces almost negligible error. This combination shown to be optimum at  $\tau = 1.0$  in Figure 7 is also optimum for other time periods (see Table 2). Therefore once an optimum combination of parameters is chosen it is valid throughout the entire transient period of an experiment. From Figures 7, 8 and 9 one can also see that the optimum ratio of  $d_t/d$  which gives zero temperature error decreases as the  $\kappa_3/\kappa_1$  ratio increases.

This implies that for the thermocouple with a larger thermal conductivity a smaller diameter is sufficient to eliminate the temperature distortion. The result shown in Figures 7, 8 and 9 can in general be adapted for use in practical application to choose the size of thermocouple and cavity, the thermocouple material and the depth of the cavity to be drilled.

As mentioned earlier that when  $\kappa_3/\kappa_1 \leq 1$  and  $\rho_3 c_3/\rho_1 c_1 \leq 1$  the errors of the temperature response at the base of the cavity are all positive. However we found (see Tables 2.2, 3.2 and 4.2) that if  $\rho_3 c_3/\rho_1 c_1$  ratio is made large enough during the transient period the error of temperature response at the base of the cavity may indeed become negative even when  $\kappa_3/\kappa_1 \leq 1$ . Physically although the thermal conductivity of the thermocouple  $\kappa_3$  is smaller than that of the disk material, but with a larger heat capacitance  $\rho_3 c_3/\rho_1 c_1 > 1$  the thermocouple is still capable of absorbing extra heat flux and hence eliminates the temperature distortion at the cavity base during the transient period. To illustrate this fact we examine Figure 7 (or see Table 4.2) for the data of  $\kappa_3/\kappa_1 = 1$  and  $\rho_3 c_3/\rho_1 c_1 = 1.3$ . One sees that when  $d_t/d = 1$  the temperature distortion can indeed be negative. Therefore, if  $d_t/d$  are chosen between 0.8 and 1 the error can be minimized. However one must keep in mind that the elimination of error by heat capacitance can work during the transient period only, for once a steady state conduction is established the heat capacity  $\rho c$  will no longer have any effect and overheating at the cavity base eventually will develop. This can be seen best from the governing equation (1) that at steady state the unsteady term which contains  $\rho c$  product is zero and is not a parameter affecting the distortion.

Another important fact that should be mentioned is that in general the optimum choice of  $d_t/d$  ratio for given  $\kappa_3/\kappa_1$  and  $\rho c$  ratio does not vary very much with the variation of  $\epsilon/D$  ratio. The insensitivity of the optimum  $d_t/d$  ratio to the  $\epsilon/D$  ratio ranging from 0.02 to 0.1 means that the distortion of the temperature is insensitive to the cavity depth or the thickness of the disk. This fact was already pointed out by Chen and Danh [11] in their experiment that the temperature distortion at the base of the cavity is more sensitive to the variation of the cavity diameter than the depth of the cavity drilled.

As an example of a practical application, let us consider a measurement of the transient temperature response of an engine block made of aluminum. From Figure 3b we know that aluminum has high thermal conductivity. Therefore copper-constantan thermocouple which has a higher thermal conductivity than aluminum should be chosen. For this material combination we have  $\kappa_3/\kappa_1 = 1.69$   $\rho_3 c_3/\rho_1 c_1 = 1.3$ . Now if the thermocouple cavity is drilled such that  $\epsilon/D = 0.1$  then from Figure 7 interpolating between  $\kappa_3/\kappa_1 = 2$  and 1 for  $\rho_3 c_3/\rho_1 c_1 = 1.3$  we find that the optimum  $d_t/d$  for  $\kappa_3/\kappa_1 = 1.69$  is approximately 0.7

One disadvantage of invoking finite element analysis is that the result does not give a clear functional relation among the parameters involved. In an attempt to obtain a simple and useful relation to relate the various parameters we note the following fact and result:

- (a) the optimum  $d_t/d$  ratio for zero temperature distortion is a strong function of  $\kappa_3/\kappa_1$  and  $\rho c$  ratio but is relatively insensitive to the  $\epsilon/D$  ratio, (b) from the theoretical reasoning the  $d_t/d$  ratio is independent

of  $\rho c$  ratio if the problem is steady state. A simple steady one dimensional analysis in which the thermocouple and the insulation material in the cavity is made to conduct the same amount of heat that would be transferred without the cavity gives the relation

$$d_t/d = \sqrt{(\kappa_1 - \kappa_2)/(\kappa_3 - \kappa_2)} \quad (9)$$

Using the above equation as a base we find that for the transient heat conduction as calculated by the finite element method the following equation (10) correlates very well with the optimum  $d_t/d$  ratio.

$$d_t/d = (\rho_3 c_3 / \rho_1 c_1)^{0.3} \sqrt{(\kappa_1 - \kappa_2)/(\kappa_3 - \kappa_2)} \quad (10)$$

Equation (10) gives an error or distortion of no more than two percentage points. In practice equation (10) may be used as a rule of thumb.

#### CONCLUSION

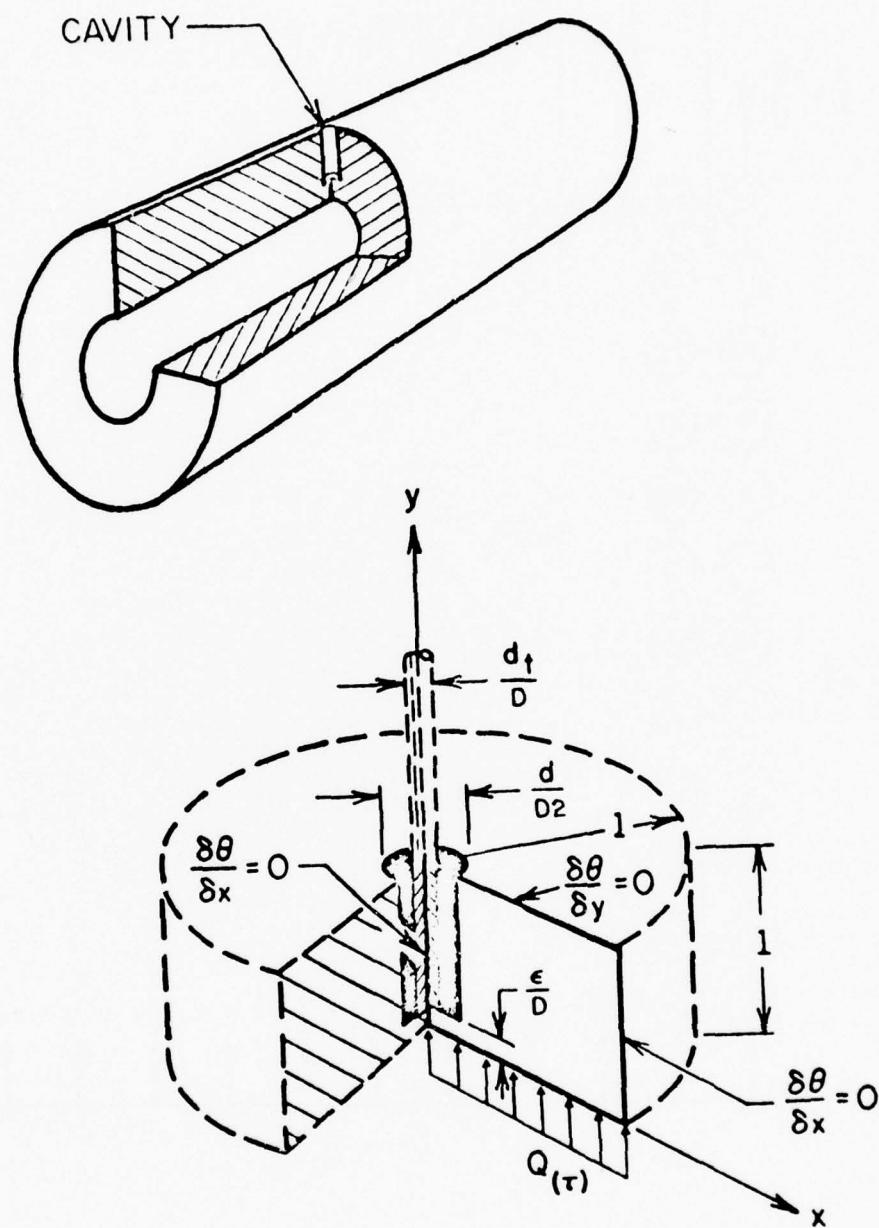
An analysis of the temperature distortion caused by the cavity drilled into a disk to accommodate the thermocouple has been studied. The calculation is carried out for the case of constant heat flux. It is shown that the temperature at the base of the cavity distorted from that without a cavity can be eliminated by a properly chosen combination of the ratio of the thermocouple diameter to the cavity diameter,  $d_t/d$  and the thermocouple material  $\kappa_3/\kappa_1$ . The optimum ratio of  $d_t/d$  can be found from Figures 7, 8 and 9 or Tables 2, 3 and 4, or approximately from equation (10). As a rule the thermocouple must be chosen

to have a higher thermal conductivity than that of the heat conducting solid. The cavity diameter should be as small as practically possible. For the case of time dependent surface heat flux the present result may be also used as a general guide.

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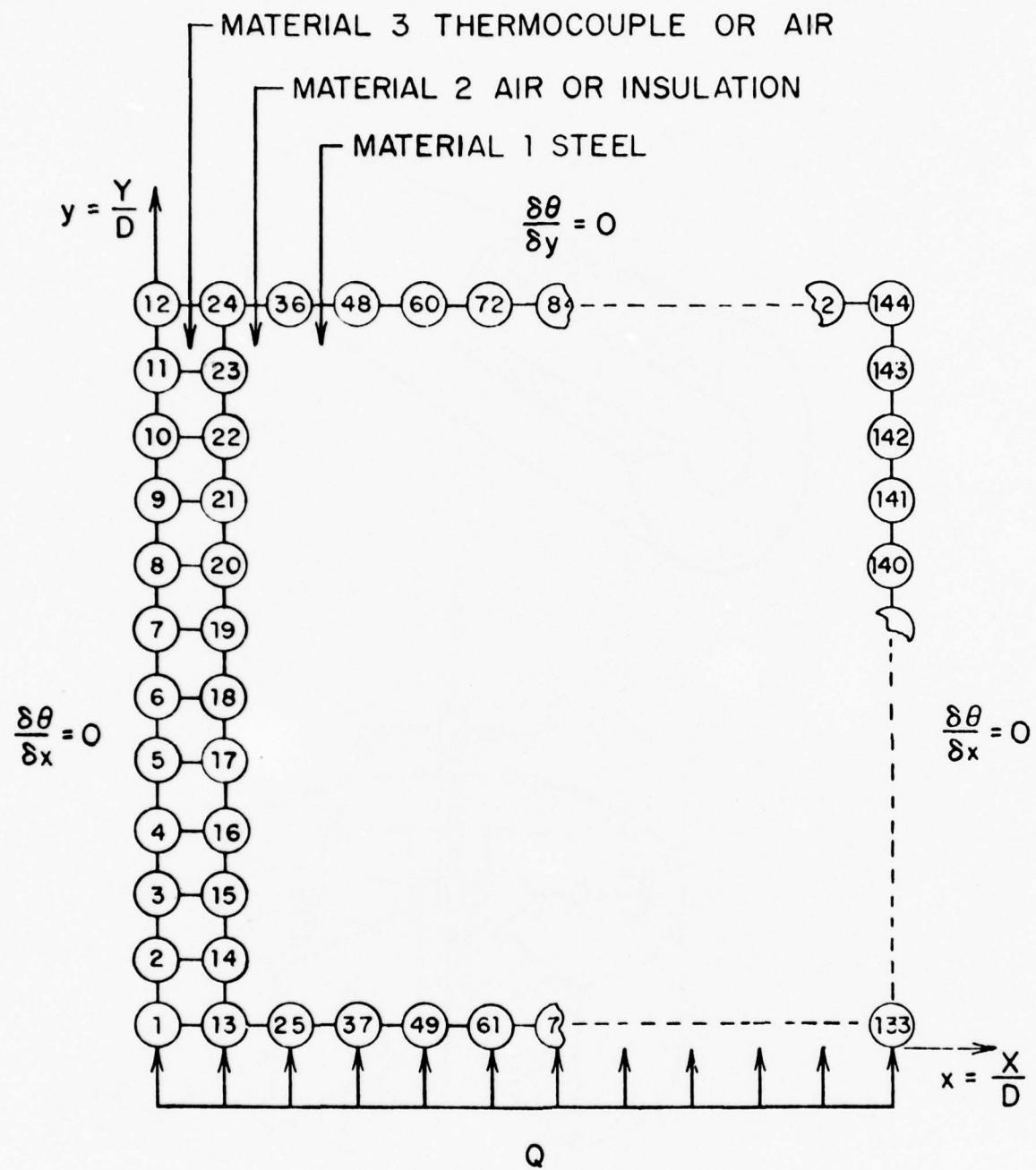
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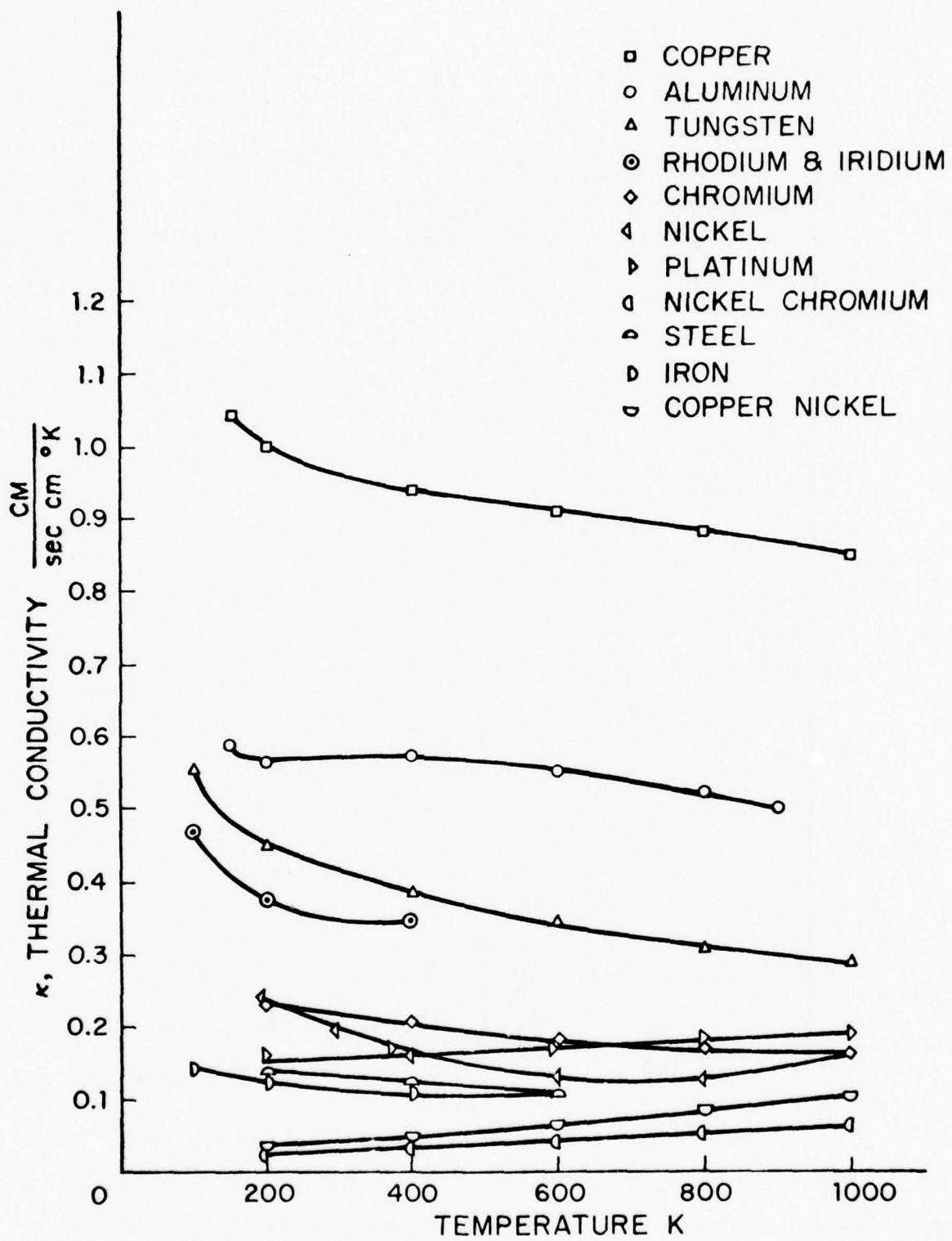


- Steel
- ▨ Thermocouple
- Air Infiltration

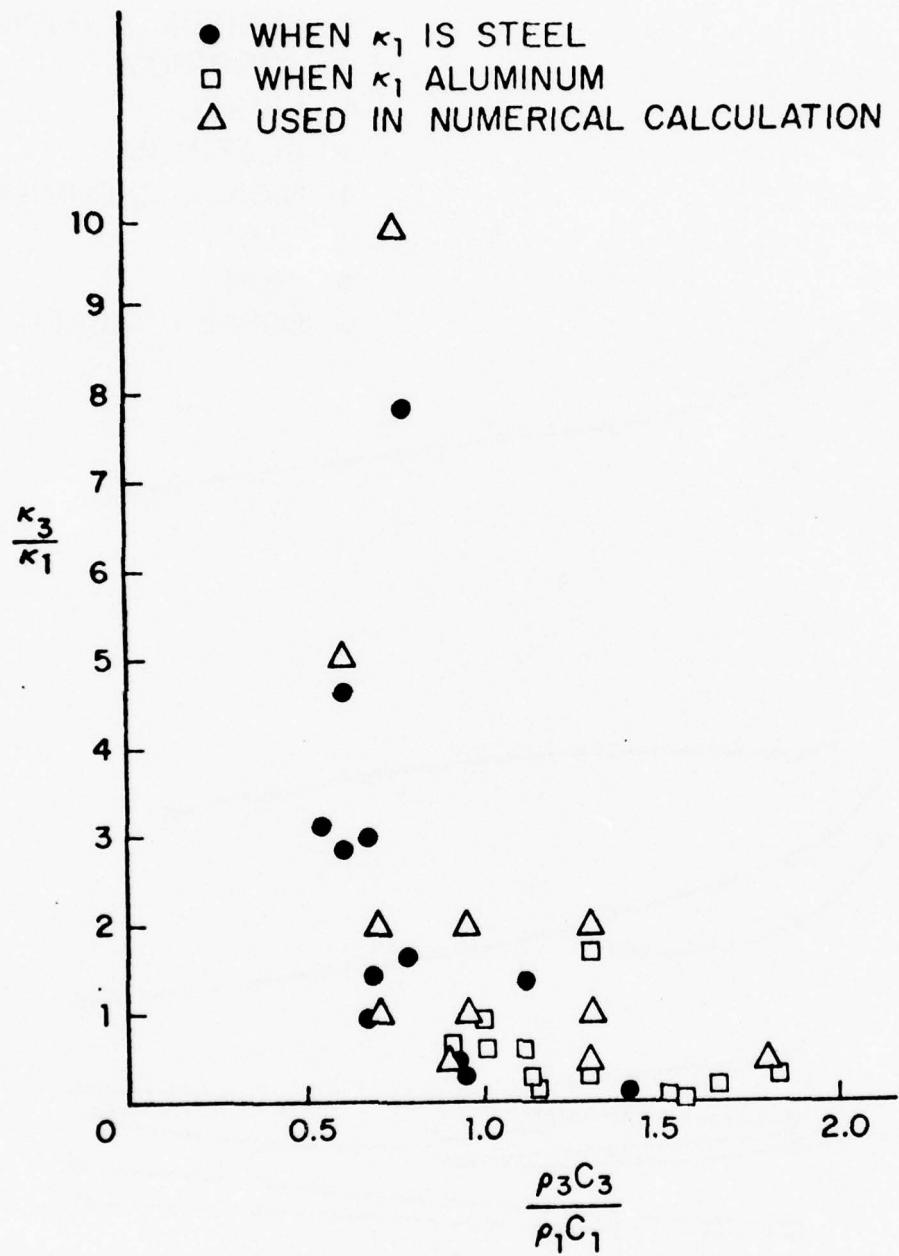
PART I Figure 1 Geometric Representation of Problem  
19



PART I Figure 2 Finite Element Idealization

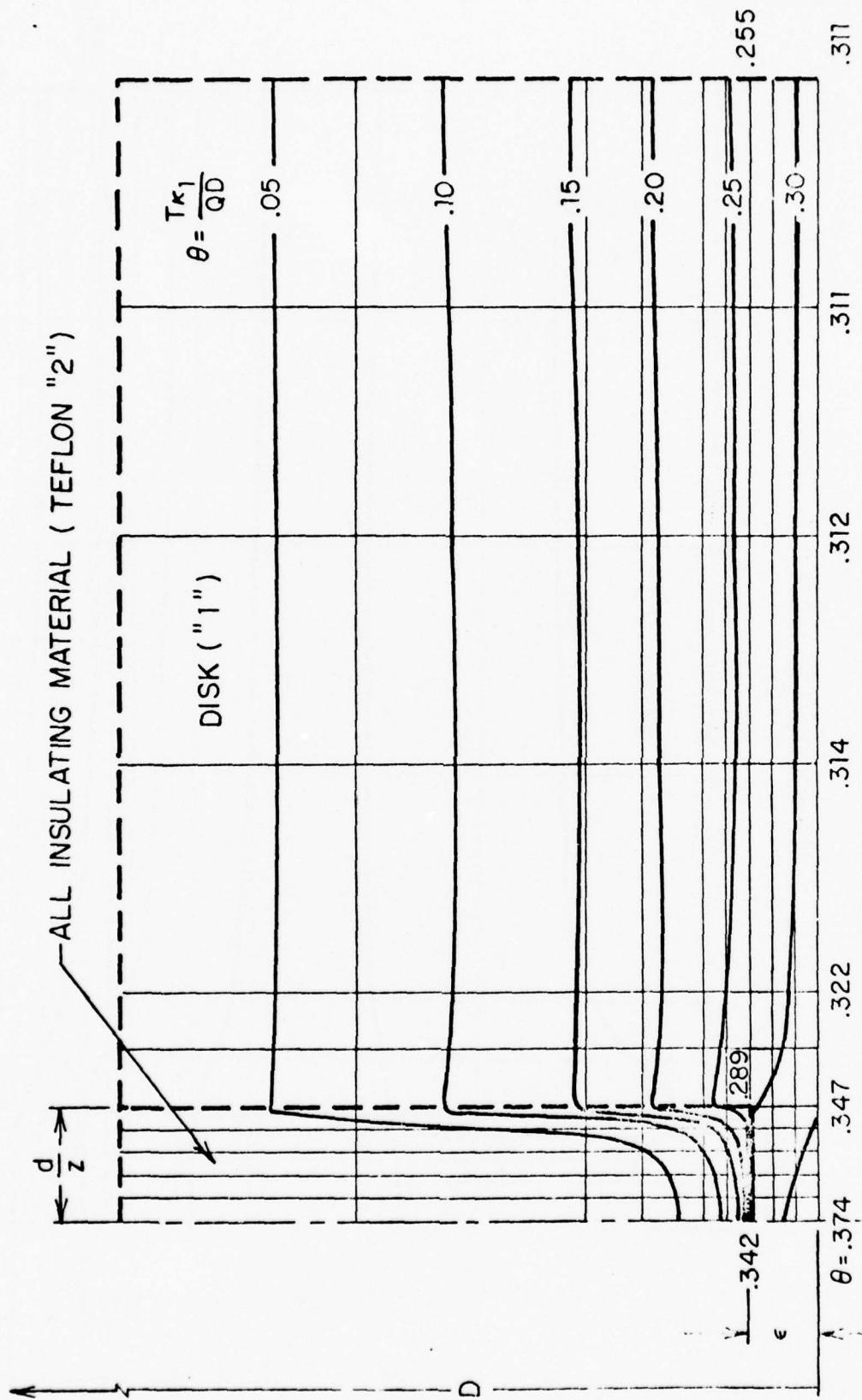


PART I Figure 3a Thermal Conductivity of Thermocouple Materials



PART I Figure 3b Variation of Thermal Conductivity Ratio and Density - Specific Heat Ratio

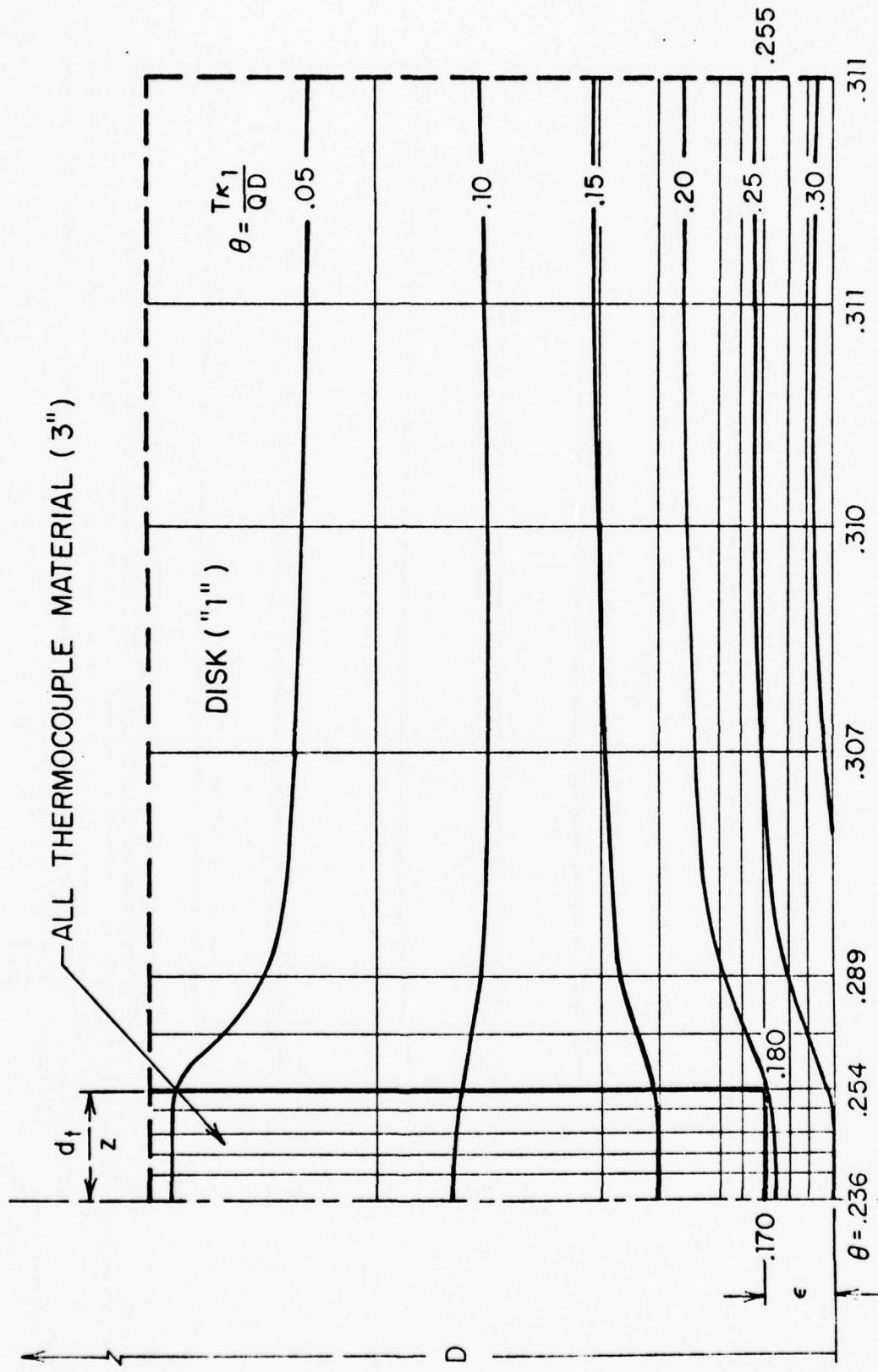
$$\tau = .08, \frac{\kappa_2}{\kappa_1} = .005, \frac{\rho c_2}{\rho c_1} = .5, \frac{d}{D} = 0, \frac{\epsilon}{D} = .06, \frac{d}{2D} = .1$$



the Cavity

PART I Figure 4 Temperature Distribution with Insulation Material Filled

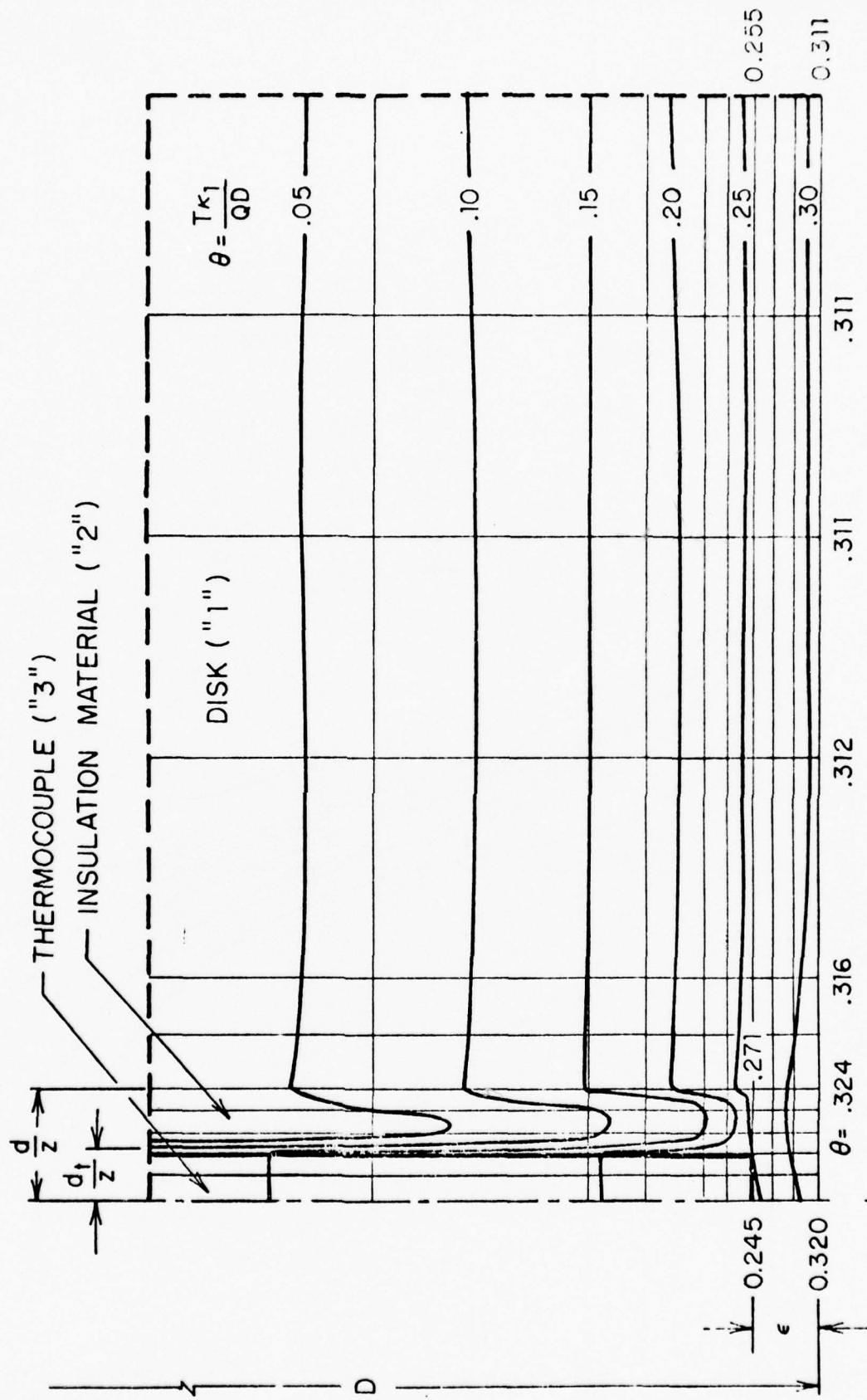
$$\tau = .08, \frac{k_3}{\kappa_1} = 10, \frac{\rho c_3}{\rho c_1} = .75, \frac{d_t}{d} = 1.0, \frac{\epsilon}{D} = .06, \frac{d}{2D} = .1$$



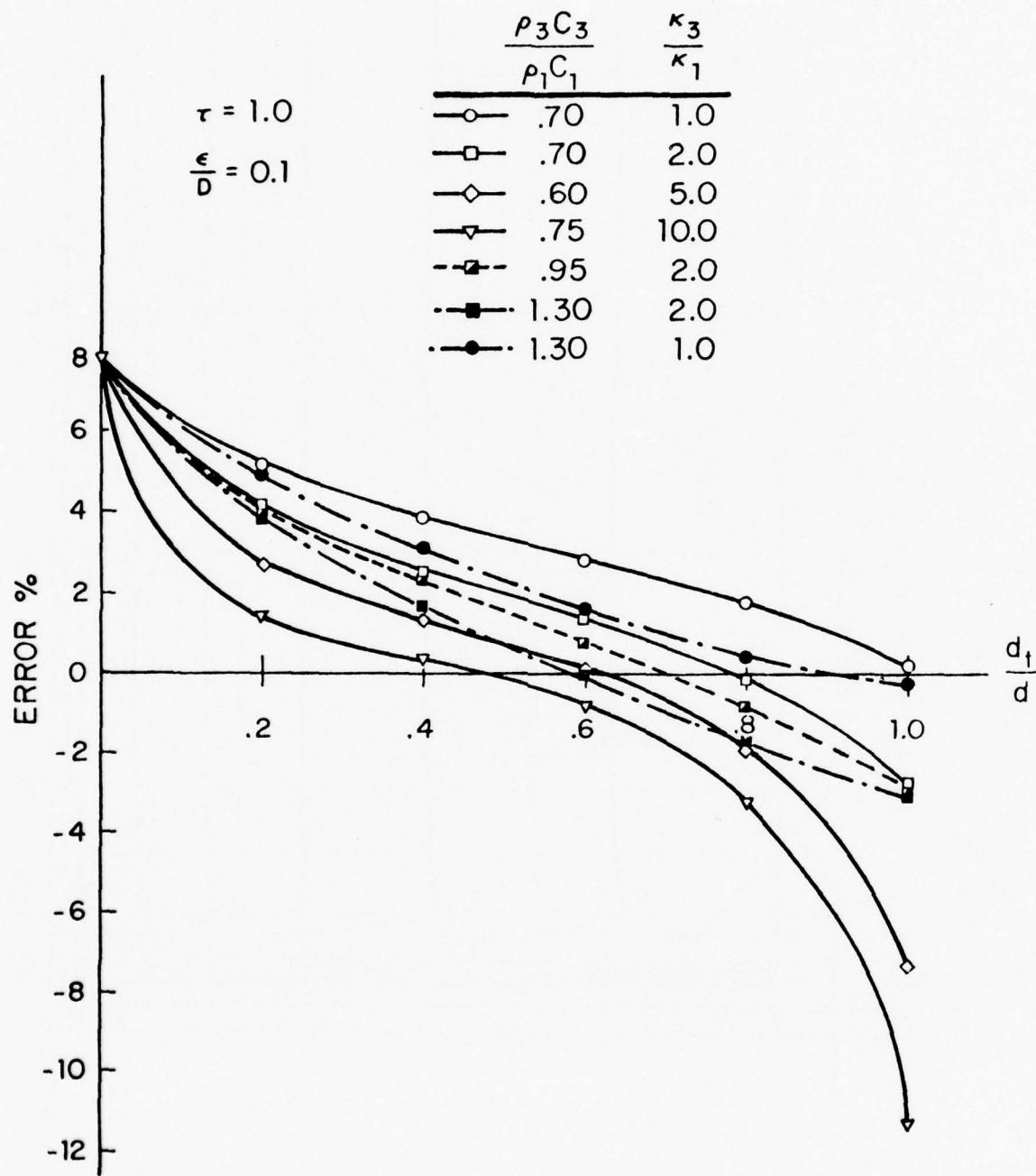
PART I Figure 5 Temperature Distribution with Thermocouple Material Filled the Cavity

$$\tau = .08, \frac{\kappa_2}{\kappa_1} = 10, \frac{\rho c_3}{\rho c_1} = .75, \frac{d_1}{d} = .4, \frac{\epsilon}{D} = .06 \frac{d}{2D} = .1$$

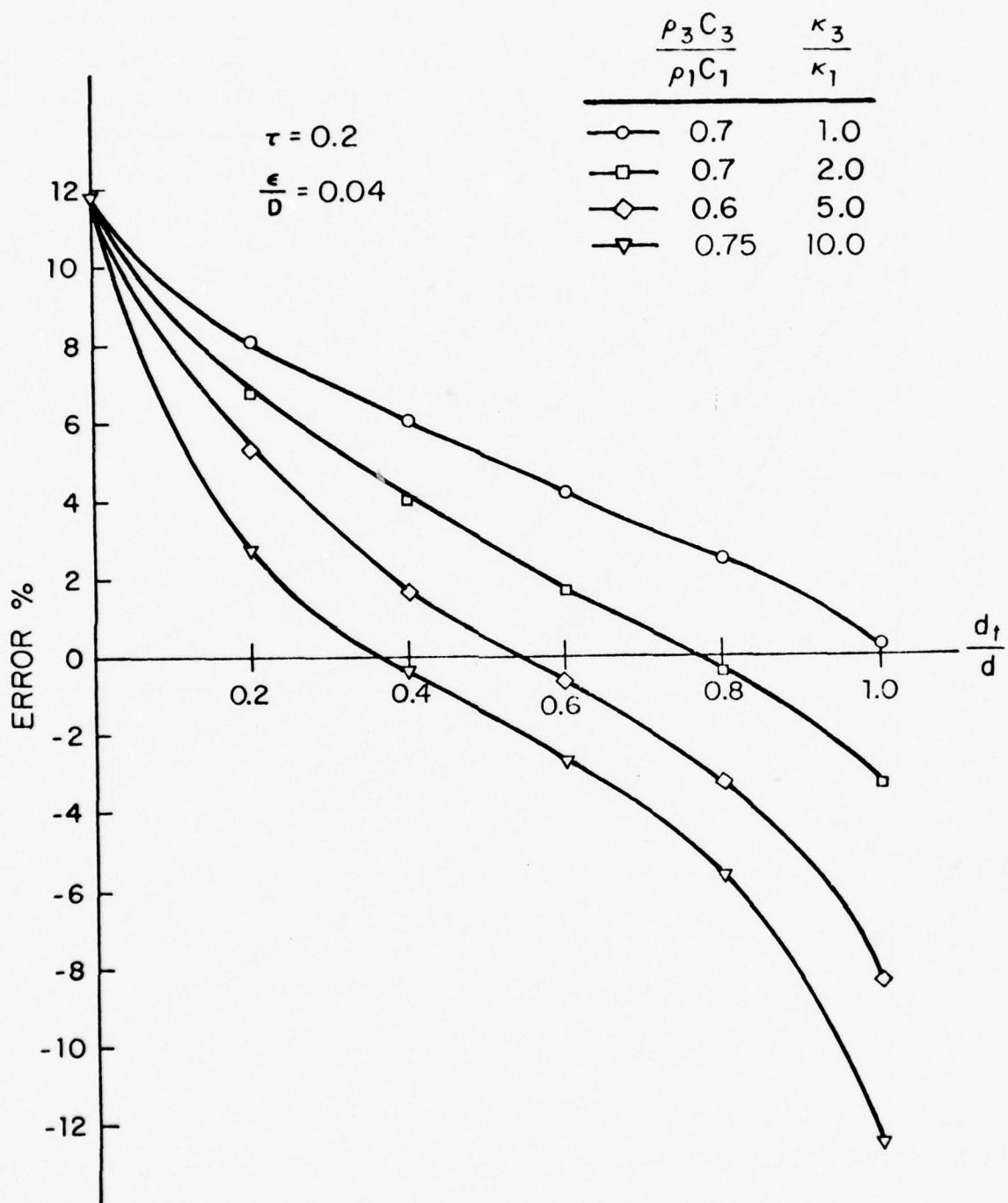
$$\frac{\kappa_2}{\kappa_1} = .005, \frac{\rho c_2}{\rho c_1} = .5$$



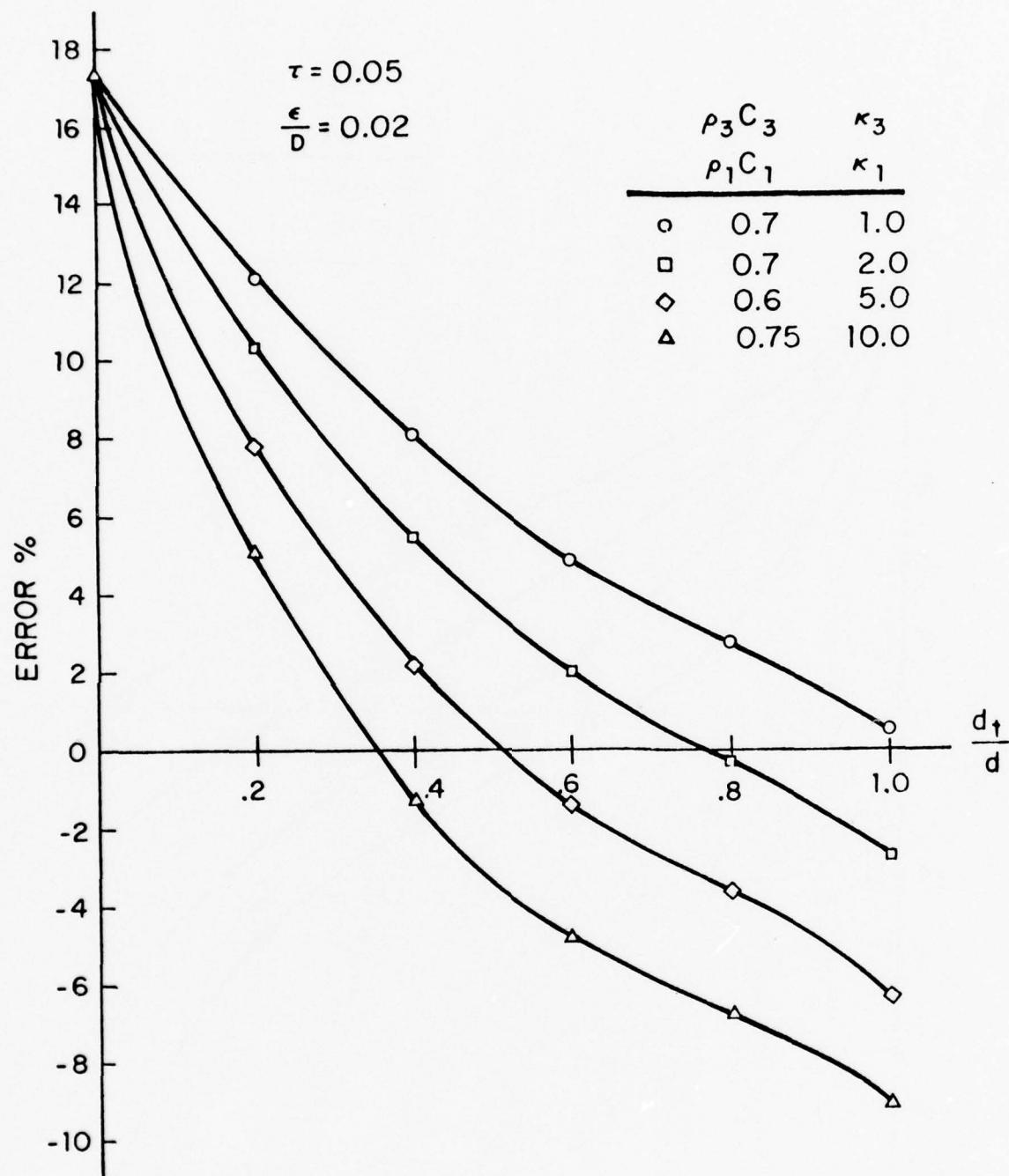
**PART I Figure 6** Temperature Distribution with Thermocouple Material Partially Filled the Cavity



PART I Figure 7 Percentage Error vs.  $d_t/d$  Ratio for Various  $\kappa_3/\kappa_1$ , and  $\rho c$   
 Ratios  $\epsilon/D = .1$



PART I Figure 8 Percentage Error vs.  $d_t/d$  Ratio for Various  $k_3/k_1$  and  $\rho c$   
Ratios  $\epsilon/D = 0.04$



PART I Figure 9 Percentage Error vs.  $d_t/d$  Ratio for Various  $k_3/k_1$  and  $\rho_3 C_3 / \rho_1 C_1$   
 Ratios  $\epsilon/D = 0.02$

PART I Table 1a  
Commonly Used Thermocouples

Types of Thermocouples	Temperature Ranges	
	°F	°C
Copper/Constantan	-300/750	-148/398
Iron/Constantan	-300/1600	-148/871
Chromel/Alumel	-300/2300	-148/1260
Chromel/Constantan	32/1800	0/982
Platinum 10% Rhodium/Platinum	32/2800	0/1537
Platinum 13% Rhodium/Platinum 6% Rh	100/3270	37/1798
Platinel 1813 Platinel 1503	32/2372	0/1300
Iridium/Iridium 60% Rhodium	2552/3326	1400/1830
Tungsten 3% Rhenium/Tungsten 25% Rhenium	50/4000	10/2204
Tungsten/Tungsten 26% Rhenium	60/5072	15/2800
Tungsten 5% Rhenium/Tungsten 26% Rhenium	32/5000	0/2760

PART I Table 1b  
 Thermal Properties of Thermocouple Materials

Metal, Insulator (subscript 3)	$K_3$	$\frac{cal}{sec. cm^2 K}$	$\frac{\kappa_3}{\kappa_{steel}}$	$\frac{\kappa_3}{\kappa_{aluminum}}$	$\rho_3 c_3$	$\frac{cal}{cm^3}$	$\frac{\rho_3 c_3}{\rho_{steel}}$	$\frac{\rho_3 c_3}{\rho_{aluminum}}$
Aluminum	0.554	4.66	1.00	0.65	0.60	1.0		
Copper	0.935	7.87	1.69	0.84	0.78	1.30		
Chromium	0.201	1.70	0.36	0.84	0.78	1.30		
Nickel	0.160	1.35	0.29	1.20	1.12	1.86		
Platinum	0.174	1.46	0.31	0.74	0.69	1.14		
Steel	0.119	1.0	0.21	1.08	1.00	1.66		
Tungsten	0.373	3.14	0.67	0.59	0.55	0.92		
Iridium	0.345	2.91	0.62	0.66	0.61	1.01		
Rhodium	0.360	3.03	0.65	0.72	0.67	1.12		
Rhenium	0.111	0.94	0.20	0.73	0.68	1.14		
Nickel-Chromium	0.040	0.34	0.07	1.01	0.94	1.56		
Copper-Nickel	0.059	0.50	0.11	0.99	0.92	1.52		
Teflon	$6 \times 10^{-4}$	0.005	0.001	0.54	0.5	0.83		
Air	$1.26 \times 10^{-4}$	0.001	0.0002	0.53	0.49	0.82		

Note: The value quoted is the averaged value over 200 to 800 °K whenever data are available.

## PART I Table 2

Percentage Error of Temperature at Cavity Base (  $/D = 0.1$  )

Table 2.1

$\rho_3 c_3 / \rho_1 c_1$	$d_t / d$	$\kappa_3 / \kappa_1 = 0.5$		$\epsilon / D = 0.1$			
		0.0	0.2	0.4	0.6	0.8	1.0
		$\tau \times 10^2$				Error %	
0.9	1	1.9	1.4	1.0	0.8	0.6	0.5
	2	3.4	2.5	1.9	1.4	1.1	1.0
	4	5.1	3.8	2.9	2.2	1.7	1.5
	8	6.4	4.7	3.7	2.9	2.3	1.9
	12	6.9	5.1	4.1	3.3	2.6	2.2
	20	7.4	5.5	4.5	3.7	2.9	2.4
	40	7.8	5.9	4.8	4.0	3.2	2.6
	60	7.9	6.0	4.9	4.0	3.2	2.6
	80	7.9	6.0	4.8	3.9	3.2	2.6
	100	7.9	6.0	4.8	3.9	3.2	2.6
1.3	1	1.3	0.9	0.6	0.4	0.3	
	2	2.4	1.7	1.1	0.8	0.6	
	4	3.7	2.6	1.8	1.2	1.1	
	8	4.6	3.4	2.4	1.6	1.5	
	12	5.0	3.8	2.7	1.9	1.8	
	20	5.4	4.2	3.1	2.2	2.0	
	40	5.8	4.5	3.4	2.5	2.3	
	60	5.9	4.5	3.4	2.5	2.3	
	80	5.9	4.5	3.3	2.5	2.3	
	100	5.8	4.5	3.3	2.5	2.3	
1.8	1	1.2	0.7	0.4	0.1	0.1	
	2	2.3	1.4	0.8	0.4	0.3	
	4	3.6	2.3	1.4	0.7	0.6	
	8	4.5	3.0	1.9	0.9	1.0	
	12	4.9	3.4	2.1	1.1	1.2	
	20	5.3	3.8	2.5	1.5	1.6	
	40	5.6	4.1	2.8	1.7	1.9	
	60	5.7	4.1	2.7	1.7	1.9	
	80	5.7	4.1	2.6	1.6	1.9	
	100	5.7	4.0	2.6	1.6	1.9	

Table 2.2

$\rho_3 c_3 / \rho_1 c_1$	$d_t / d$	$\kappa_3 / \kappa_1 = 1.0$			$\epsilon / D = 0.1$		
		0.0	0.2	0.4	0.6	0.8	1.0
$\tau \times 10^2$		Error %					
0.7	1	1.2	0.8	0.5	0.3	0.2	
	2	2.3	1.5	1.0	0.6	0.2	
	4	3.4	2.4	1.6	1.0	0.3	
	8	4.3	3.1	2.1	1.3	0.3	
	12	4.6	3.3	2.3	1.5	0.3	
	20	4.9	3.6	2.6	1.6	0.3	
	40	5.2	3.8	2.8	1.7	0.2	
	60	5.2	3.9	2.8	1.7	0.2	
	80	5.2	3.9	2.8	1.8	0.2	
	100	5.2	3.9	2.8	1.8	0.2	
0.95	1	1.2	0.7	0.4	0.2	0.0	
	2	2.2	1.4	0.8	0.3	0.0	
	4	3.3	2.2	1.2	0.6	0.0	
	8	4.2	2.8	1.7	0.8	0.1	
	12	4.5	3.1	1.9	0.9	0.0	
	20	4.8	3.4	2.2	1.1	0.0	
	40	5.1	3.6	2.3	1.2	0.0	
	60	5.1	2.6	2.3	1.2	0.0	
	80	5.1	3.6	2.3	1.2	0.0	
	100	5.1	3.6	2.3	1.2	0.0	
1.3	1	1.1	0.6	0.2	-0.0	-0.1	
	2	2.1	1.1	0.5	-0.0	-0.2	
	4	3.2	1.8	0.8	0.0	-0.3	
	8	4.0	2.4	1.2	0.2	-0.3	
	12	4.4	2.7	1.4	0.3	-0.3	
	20	4.7	3.0	1.6	0.4	-0.3	
	40	5.0	3.2	1.7	0.4	-0.2	
	60	5.0	3.2	1.7	0.4	-0.2	
	80	5.0	3.2	1.6	0.4	-0.2	
	100	5.0	3.1	1.6	0.4	-0.2	

Table 2.3

$\rho_3 c_3 / \rho_1 c_1$	$d_t / d$	$\kappa_3 / \kappa_1 = 5.0$		$\epsilon / D = 0.1$			
		0.0	0.2	0.4	0.6	0.8	1.0
		$\tau \times 10^2$					
0.6	1	0.7	0.1	-0.3	-0.5	-1.0	
	2	1.3	0.1	-0.4	-0.9	-1.8	
	4	2.1	0.5	-0.6	-1.4	-3.0	
	8	2.5	0.6	-0.7	-1.9	-4.5	
	12	2.5	0.7	-0.7	-2.0	-5.4	
	20	2.6	0.9	-0.4	-2.0	-6.4	
	40	2.6	1.2	-0.1	-1.9	-7.2	
	60	2.7	1.4	-0.2	-1.8	-7.3	
	80	2.7	1.4	-0.2	-1.8	-7.3	
	100	2.7	1.4	-0.2	-1.8	-7.3	

Table 2.4

		$\kappa_3 / \kappa_1 = 10$		$\epsilon / D = 0.1$			
		0.3	-0.4	-0.8	-1.0	-1.4	
0.75	1	0.3	-0.4	-0.8	-1.0	-1.4	
	2	0.7	-0.6	-1.4	-1.9	-2.7	
	4	1.1	-0.8	-2.0	-2.9	-4.6	
	8	1.3	-0.9	-2.4	-3.7	-6.8	
	12	1.3	-0.8	-2.3	-3.8	-8.3	
	20	1.2	-0.4	-1.8	-3.7	-9.9	
	40	1.4	0.2	-1.0	-3.3	-11.1	
	60	1.4	0.3	-0.9	-3.2	-11.3	
	80	1.4	0.4	-0.8	-3.2	-11.3	
	100	1.4	0.4	-0.8	-3.2	-11.3	

## PART I Table 3

Percentage Error of Temperature at Cavity ( $\epsilon/D = 0.04$ )

$\rho_3 c_3 / \rho_1 c_1$	$d_t / d$	$\kappa_3 / \kappa_1 = 0.5$		$\epsilon / D = 0.04$			
		0.0	0.2	0.4	0.6	0.8	1.0
		$\tau \times 10^2$		Error %			
0.9	.2	1.5	0.9	0.7	0.5	0.4	0.4
	.4	3.2	2.0	1.5	1.1	0.8	0.7
	.8	5.6	3.9	2.8	2.0	1.5	1.3
	1.6	8.2	6.0	4.5	3.3	2.4	1.9
	2.4	9.3	7.0	5.3	3.9	2.9	2.3
	4.0	10.3	7.8	6.1	4.5	3.4	2.7
	8.0	11.2	8.5	6.7	5.2	3.9	3.1
	12.0	11.5	8.8	7.1	5.5	4.2	3.3
	16.0	11.8	9.0	7.3	5.8	4.5	3.4
	20.0	11.9	9.1	7.5	6.0	4.6	3.5
1.3	.2		0.8	0.5	0.3	0.2	0.2
	.4		1.9	1.2	0.8	0.5	0.4
	.8		3.7	2.5	1.6	1.0	0.8
	1.6		5.8	4.0	2.6	1.6	1.3
	2.4		6.8	4.9	3.2	2.1	1.6
	4.0		7.6	5.6	3.8	2.5	2.0
	8.0		8.4	6.3	4.4	2.9	2.5
	12.0		8.6	6.6	4.7	3.2	2.7
	16.0		8.8	6.9	5.0	3.4	2.9
	20.0		9.0	7.0	5.2	3.6	3.0
1.8	.2		0.7	0.4	0.2	0.1	0.0
	.4		1.7	0.9	0.5	0.2	0.1
	.8		3.4	2.1	1.1	0.5	0.3
	1.6		5.5	3.5	2.0	0.9	0.6
	2.4		6.5	4.3	2.5	1.2	0.9
	4.0		7.4	5.1	3.1	1.6	1.2
	8.0		8.2	5.8	3.6	1.9	1.7
	12.0		8.5	6.1	3.9	2.1	2.0
	16.0		8.7	6.3	4.1	2.3	2.2
	20.0		8.8	6.5	4.3	2.5	2.4

Table 3.2

 $\kappa_3/\kappa_1 = 1.0$  $\epsilon/D = 0.04$ 

$\rho_3 c_3 / \rho_1 c_1$	$d_t / d$	0.0	0.2	0.4	0.6	0.8	1.0
		$\tau \times 10^2$					
		Error %					
0.7	0.2	0.8	0.5	0.3	0.2	0.1	
	0.4	1.8	1.1	0.7	0.4	0.3	
	0.8	3.4	2.2	1.4	0.8	0.4	
	1.6	5.4	3.6	2.2	1.3	0.4	
	2.4	6.3	4.3	2.7	1.6	0.4	
	4.0	7.1	5.0	3.2	1.9	0.4	
	8.0	7.6	5.5	3.7	2.2	0.4	
	12.0	7.9	5.8	4.0	2.4	0.4	
	16.0	8.0	5.9	4.1	2.4	0.4	
	20.0	8.1	6.1	4.2	2.5	0.3	
0.95	0.2	0.7	0.4	0.2	0.1	0.0	
	0.4	1.6	0.9	0.5	0.2	0.0	
	0.8	3.3	1.9	1.0	0.4	0.1	
	1.6	5.2	3.3	1.8	0.7	0.1	
	2.4	6.1	3.9	2.2	0.9	0.1	
	4.0	6.9	4.6	2.6	1.1	0.1	
	8.0	7.5	5.1	3.1	1.4	0.1	
	12.0	7.7	5.4	3.3	1.6	0.1	
	16.0	7.9	5.6	3.5	1.6	0.1	
	20.0	8.0	5.7	3.6	1.7	0.1	
1.3	0.2	0.6	0.3	0.0	-0.1	-0.1	
	0.4	1.5	0.7	0.2	-0.1	-0.2	
	0.8	3.1	1.6	0.6	-0.1	-0.3	
	1.6	5.0	2.8	1.2	0.1	-0.4	
	2.4	5.9	3.5	1.5	0.2	-0.4	
	4.0	6.7	4.1	1.0	0.3	-0.4	
	8.0	7.3	4.6	2.2	0.4	-0.4	
	12.0	7.6	4.9	2.5	0.5	-0.4	
	16.0	7.7	5.1	2.7	0.6	-0.4	
	20.0	7.9	5.2	2.8	0.7	-0.3	

Table 3.3

$\rho_3 c_3 / \rho_1 c_1$	$d_t / d$	$\kappa_3 / \kappa_1 = 5$		$\epsilon / D = 0.04$			
		0.0	0.2	0.4	0.6	0.8	1.0
$\tau \times 10^2$		Error %					
0.6	0.2	0.3	-0.1	-0.3	-0.4	-0.6	
	0.4	0.9	0.0	-0.5	-0.8	-1.1	
	0.8	2.0	0.4	-0.6	-1.2	-2.1	
	1.6	3.4	0.9	-0.7	-1.8	-3.2	
	2.4	4.1	1.2	-0.8	-2.1	-4.0	
	4.0	4.6	1.4	-0.9	-2.5	-5.0	
	8.0	4.7	1.4	-1.1	-3.0	-6.4	
	12.0	4.7	1.4	-1.0	-3.2	-7.3	
	16.0	4.7	1.6	-0.8	-3.2	-7.9	
	20.0	4.7	1.7	-0.6	-3.2	-8.3	

Table 3.4

	$\kappa_3 / \kappa_1 = 10$	$\epsilon / D = 0.04$					
		0.2	0.4	0.6	0.7	0.9	1.7
0.75	0.2	0.0	-0.4	-0.6	-0.7	-0.9	
	0.4	0.3	-0.6	-1.1	-1.4	-1.7	
	0.8	1.0	-0.7	-1.7	-2.3	-3.0	
	1.6	2.1	-0.8	-2.5	-3.5	-4.6	
	2.4	2.6	-0.8	-2.9	-4.2	-5.7	
	4.0	2.9	-0.9	-3.4	-5.0	-7.2	
	8.0	2.9	-1.0	-3.7	-5.8	-9.5	
	12.0	2.8	-0.8	-3.4	-5.8	-11.0	
	16.0	2.8	-0.5	-3.1	-5.7	-11.9	
	20.0	2.8	-0.3	-2.7	-5.6	-12.5	

## PART I Table 4

Percentage Error of Temperature at Cavity Base ( $\epsilon/D = 0.02$ )

$\rho_3 c_3 / \rho_1 c_1$	$d_t/d$	$\kappa_3 / \kappa_1 = 0.5$			$\epsilon/D = 0.02$		
		0.0	0.2	0.4	0.6	0.8	1.0
		$\tau \times 10^2$					
					Error %		
1.3	0.05	0.3	0.1	0.1	0.1	0.1	0.0
	0.10	0.8	0.4	0.2	0.2	0.2	0.1
	0.20	2.0	1.2	0.7	0.5	0.4	
	0.40	4.5	2.7	1.6	1.0	0.8	
	0.60	6.5	4.0	2.3	0.4	1.1	
	1.00	8.9	5.6	3.3	2.0	1.5	
	2.00	11.4	7.6	4.7	2.8	2.2	
	3.00	12.4	8.5	5.4	3.3	2.6	
	4.00	12.9	9.0	5.8	3.6	2.8	
	5.00	13.1	9.3	6.0	3.7	3.0	

Table 4.2

$\rho_3 c_3 / \rho_1 c_1$	$d_t / d$	$\kappa_3 / \kappa_1 = 1$		$\epsilon / D = 0.02$			
		0.0	0.2	0.4	0.6	0.8	1.0
$\tau \times 10^2$		Error %					
0.7	0.05	0.3	0.2	0.1	0.1	0.1	0.1
	0.10	0.8	0.5	0.3	0.2	0.2	0.2
	0.20	1.9	1.1	0.7	0.4	0.3	0.3
	0.40	4.3	2.5	1.5	0.8	0.5	0.5
	0.60	6.1	3.6	2.1	1.1	0.5	0.5
	1.00	8.3	5.1	2.9	1.5	0.6	0.6
	2.00	10.6	6.8	3.9	2.1	0.6	0.6
	3.00	11.4	7.5	4.4	2.4	0.5	0.5
	4.00	11.8	7.8	4.7	2.5	0.5	0.5
	5.00	12.1	8.1	4.9	2.8	0.5	0.5
0.95	0.05	0.2	0.1	0.0	0.0	0.0	0.0
	0.10	0.7	0.3	0.1	0.1	0.0	0.0
	0.20	1.8	0.9	0.4	0.1	0.0	0.0
	0.40	4.0	2.1	1.0	0.3	0.1	0.1
	0.60	5.8	3.1	1.5	0.5	0.1	0.1
	1.00	8.0	4.5	2.2	0.8	0.1	0.1
	2.00	10.3	6.1	3.1	1.2	0.1	0.1
	3.00	11.2	6.8	3.5	1.4	0.1	0.1
	4.00	11.6	7.2	3.8	1.5	0.1	0.1
	5.00	11.8	7.4	4.0	1.7	0.1	0.1
1.3	0.05	0.2	0.0	-0.0	-0.1	-0.1	-0.1
	0.10	0.5	0.1	-0.1	-0.1	-0.1	-0.1
	0.20	1.6	0.6	0.1	-0.2	-0.3	-0.3
	0.40	3.7	1.6	0.5	-0.2	-0.4	-0.4
	0.60	5.4	2.5	0.8	-0.1	-0.5	-0.5
	1.00	7.6	3.8	1.4	-0.0	-0.5	-0.5
	2.00	10.0	5.4	2.1	0.2	-0.5	-0.5
	3.00	10.9	6.0	2.5	0.3	-0.5	-0.5
	4.00	11.3	6.4	2.7	0.4	-0.5	-0.5
	5.00	11.5	6.7	2.9	0.4	-0.5	-0.5

Table 4.3

$\rho_3 c_3 / \rho_1 c_1$	$d_t / d$	$\kappa_3 / \kappa_1 = 10$		$\epsilon / D = 0.02$					
		0	0.2	0.4	0.6	0.8	1.0		
		$\tau \times 10^2$							
				Error %					
0.75	0.05	0.7	0.0	-0.26	-0.33	-0.37	-0.41		
	0.10	1.7	0.0	-0.48	-0.66	-0.76	-0.86		
	0.20	4.9	0.2	-0.72	-0.93	-1.40	-1.61		
	0.40	9.0	1.2	-0.85	-1.78	-2.28	-2.71		
	0.60	11.6	2.0	-0.86	-2.20	-2.89	-2.50		
	1.00	13.2	3.2	-0.84	-2.78	-3.76	-4.66		
	2.00	15.8	4.5	-0.88	-3.66	-5.06	-6.44		
	3.00	16.7	4.9	-0.99	-4.20	-5.84	-7.56		
	4.00	17.1	5.1	-1.10	-4.57	-6.38	-8.42		
	5.00	17.4	5.1	-1.20	-4.80	-6.75	-9.12		

Table 4.4

		$\kappa_3 / \kappa_1 = 5.0$		$\epsilon / D = 0.02$					
		0	0.2	0.4	0.6	0.8	1.0		
		$\tau \times 10^2$							
				Error %					
0.6	0.05	0.7	0.05	-0.00	-0.17	-0.20	-0.25		
	0.10	1.7	0.24	-0.15	-0.32	-0.43	-0.55		
	0.20	4.9	0.87	-0.06	-0.53	-0.80	-1.09		
	0.40	9.0	1.33	0.34	-0.70	-1.29	-1.89		
	0.60	11.6	3.53	0.71	-0.78	-1.61	-2.48		
	1.00	13.2	5.12	1.25	-0.88	-2.04	-3.33		
	2.00	15.8	6.87	1.92	-1.00	-2.62	-4.59		
	3.00	16.7	7.48	2.15	-1.12	-2.97	-5.33		
	4.00	17.1	7.72	2.21	-1.25	-3.26	-5.88		
	5.00	17.4	7.80	2.20	-1.36	-3.49	-6.32		

PART II

IMPROVED ACCURACY IN THE PREDICTION OF SURFACE HEAT FLUX  
AND TEMPERATURE BY AN INTRINSIC THERMOCOUPLE

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## INTRODUCTION

In the study of transient heat transfer, many experimental difficulties may arise if heat flux sensors or thermocouples are installed direct at the surface of a body. For example, a probe may be damaged by a piston or a projectile sliding over a cylinder or barrel. A probe on a melting and ablative surface of heat shield can be easily destroyed because of high temperature. Furthermore a surface probe exposed to both radiative and convective environment may measure an erroneous surface heat flux and temperature if the probe has a different radiative property from that of the measured surface. In these circumstances, calculation of the transient surface heat flux and the surface temperature can be achieved by inverting a temperature history measured at some location inside the body.

In general, the prediction of a surface heat flux and temperature by the measured data at some location interior to a body is known as the "inverse problem". Many configurations, such as spheres, cylinders, and slabs, had been studied by many workers and many methods such as numerical, graphical, series, convolution integral, and Laplace transforms were used. Stolz [1], Beck [2] and Williams and Curry [3], considered the numerical inversion of the integral solution for semi-infinite and other bodies. In this method, care is required in selecting a time interval in order to achieve a stable solution. Carslaw and Jaeger [4], Burggraf [5], Koveryanov [6], and Shumakov [7], respectively considered different series approaches in which generally the local heat flux at an interior location and their higher derivatives are required. However, it is difficult to measure experimentally or to process the measured data for the derivative

of the temperature. Sparrow, Haji-Sheikh, and Lundgren [8], Imber and Kahn [9], Imber [10], Sabherwal [11], Masket and Vastano [12], Deverall and Channapragada [13] and Chen and Thomsen [14] applied the transform method. In these works, the solution is represented in either an integral form after some manipulation of the contour integral from the inverse transform, or in a series form after an expansion of the solution for small and large times. Using Laplace transformation Chen and Thomsen [14] introduced a polynomial in terms of an error function to represent the response of thermocouple measurement and the inversion is accomplished for any transient surface heat flux at the inner surface of a cylindrical tube. In their study, the cylindrical thickness was assumed to be relatively thick such that the temperature at a large distance from the heating surface remains constant. Therefore, only one interior temperature response near the surface was needed in the experimental measurement. Their inversion solution however was valid only for a short duration due to the asymptotic expansion of the modified Bessel function in the inverse Laplace transform. Chen and Chiou [15] studied the inversion problem for the case of a semi-infinite slab or a thick slab using a Laplace transformation. The exact solution was obtained from the inverse Laplace transform for any time interval. It was then shown that their analysis may be approximately applied to the case of the hollow cylinder if the interior temperature response is measured at a location close to the inner wall.

This report presents (a) the improved numerical solution of the inversion solution reported by Chen and Chiou [15] and (b) a further demonstration of the capability of the solution.

The theoretical analysis of Chen and Chiou [15] is recapitulated in Appendix A in which the surface heat flux and temperature is predicted by inverting a temperature history measured at some location inside the solid body. The inversion solution is obtained by invoking Laplace transformation. Both the surface heat flux and temperature are given by Eqs. (19) and (20) in Appendix A.

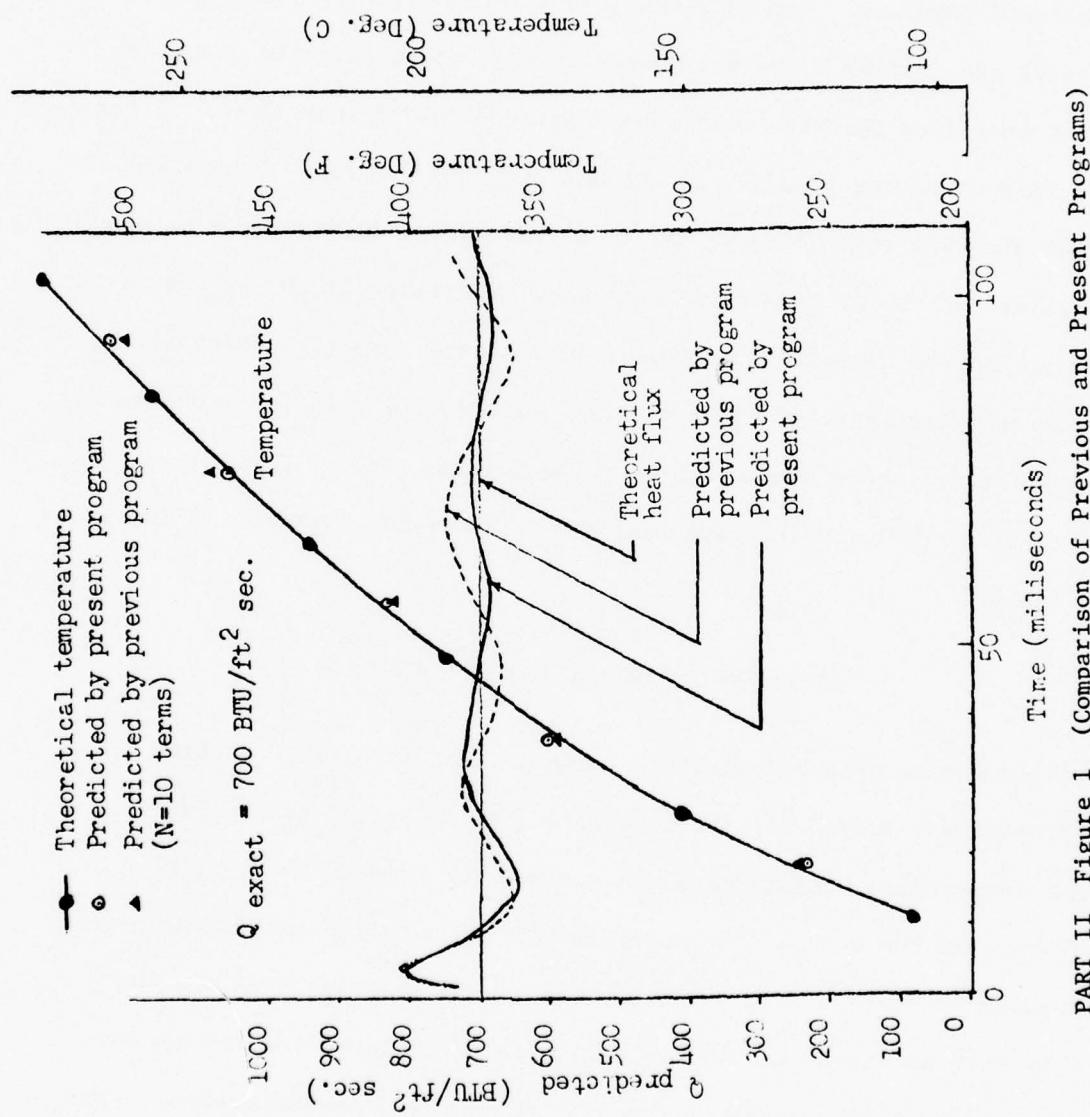
It was thought that the accuracy of the computer program generated for the solution in the previous report by Chen and Chiou [15] can be improved further for the following reasons. First, the coefficients  $b_n$  (see Appendix A Eq.(11)) in the previous formulation has a dimension of temperature. Therefore the determination of the coefficients depends on the temperature range of each particular experiment. It was found that the absolute value of the coefficients  $b_n$  in some cases can become as large as an order of  $10^{44}$ . Therefore during the subsequently numerical manipulation in the computer program error due to round off and the standard fixed up when an underflow occurred may become appreciable. To remedy this difficulty the dimensionless formulation is introduced in the analysis (Appendix A) in which the coefficient  $b_n$  is also made dimensionless. As a result the magnitude of the coefficient  $b_n$  can be greatly reduced. Secondly, the double precision format was not used throughout the previous computer program. It is felt that further accurate results may be obtained if the double precision format is adopted in the program.

In the following section the new computer solution is shown to be indeed more accurate. Later the solution is shown to be capable of predicting a case involving a periodic surface heat flux or periodic temperature variation.

## RESULTS OF THE IMPROVED COMPUTER PROGRAM

The previous computer program of Chen and Chiou [15] was recast in dimensionless form and written in the double precision format. The new computer program is listed in Appendix B. The results predicted by the new and previous computer program are given in Appendix C and shown in Figure 1 for the case of the constant surface heat flux. This is the case in which a steel slab initially at a uniform temperature is suddenly subjected to a constant heat flux  $Q$  at one of the surfaces and kept at the initial temperature on the other surface. Figure 1 shows the solution predicted by inverting the temperature response at an interior of the slab from the new and previous computer program. This solution predicted by the new and previous programs used the ten term representation for the thermocouple response. The comparison clearly shows the improvement of the new solution over the previous one. Except for the short time duration the solution with the new program reduces the error to only one half of the error of the previous program i.e., an error of less than one percent. In the short time period the solution exhibits a Gibbs phenomenon\* because of the discontinuity of the surface temperature gradient occurred at initial condition. The solution shows a 17% of initial overshoot of heat flux and then a 7.8% of undershot before the solution approaches the constant heat flux. It should be remarked that Gibbs phenomenon is artificially

\* Gibbs phenomena [3]: for a sequence of transformation  $T_n(t)$ ,  $n = 1, 2, \dots$  of a function  $q(t)$  (here  $q(t) = \text{constant}$ ) if the interval  $\lim_{t \rightarrow t_0} \inf_{n \rightarrow \infty} T_n(t)$ ,  $\lim_{t \rightarrow t_0} \sup_{n \rightarrow \infty} T_n(t)$  contains points outside the interval  $[\lim_{t \rightarrow t_0} \inf q(t), \lim_{t \rightarrow t_0} \sup q(t)]$  then the sequence is said to exhibit a Gibbs phenomena.



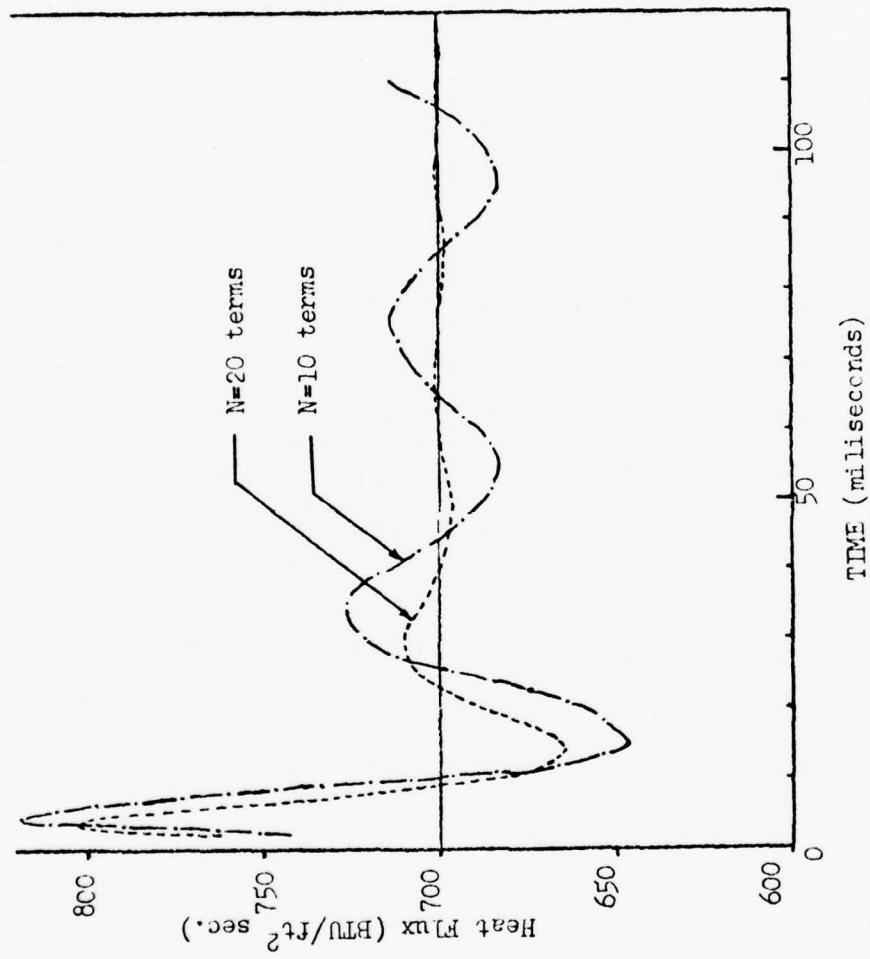
PART II Figure 1 (Comparison of Previous and Present Programs)

introduced due to the idealization of the initial condition. In most practical situations the surface heat flux will be continuous. Therefore, Gibbs phenomenon will not appear.

Figure 2 (see Table 1 also) shows the comparison between the solutions for the constant heat flux case with 10 and 20 term representation for the thermocouple response. One sees that the solution with 20 term representation after the initial Gibbs phenomenon quickly approaches the expected constant heat flux solution with a negligible error of less than 0.14 percent. This shows the accuracy of the new computer program. From Figure 2 one also observes that both overshoot and undershoot of Gibbs phenomenon are smaller for the 20 term representation. Additionally the points of the overshoot and the undershoot have moved to near the zero time which agrees with the characteristic of Gibbs phenomenon. According to Gibbs phenomenon the point of overshoot should approach the initial zero if the number of terms of the series which represent the thermocouple response is increased to infinite.

#### VERIFICATION OF OSCILLATORY SOLUTION

As a measure of applicability of the present inversion solution, a test problem was solved for the case of a slab subjected to a periodic surface temperature variation on one surface and held to the initial temperature on the other. The analytic solution for the problem is given in Appendix IV where a more suitable form of the solution than the one given by Carslaw and Jaeger [4] is derived and tabulated for the thermocouple response at one tenth of the slab thickness from the surface. The



PART II Figure 2 Comparison of Different Polynomial Representations

PART II TABLE 1

## Comparision of Inversion Prediction and Exact Solution

N	t	f(t)	$\theta(0, t)$	$\theta(0, t)$	ERROR <sup>*</sup> %	$\frac{\partial \theta}{\partial x}(0, t)$	$\frac{\partial \theta}{\partial x}(0, t)$	ERROR %
			$\theta(1, t)$	EXACT		PREDICTED	EXACT	
20	0.1	0.0624	0.3758	0.3773	+0.399	1.0000	0.9873	-1.27
	0.2	0.1628	0.5315	0.5314	+0.019	1.0000	1.0014	+0.14
	0.4	0.3395	0.7516	0.7516	0.00	1.0000	1.0001	+0.01
	0.6	0.4882	0.9205	0.9205	0.00	1.0000	1.0000	0.00
	0.8	0.6183	1.0629	1.0628	-0.009	1.0000	0.9996	-0.04
	1.0	0.7353	1.1884	1.1878	-0.05	1.0000	1.9986	-0.14
10	0.1	0.0624	0.3758	0.4048	+7.71	1.0000	1.0796	+7.96
	0.2	0.1628	0.5315	0.5284	-0.58	1.0000	0.9391	-6.09
	0.4	0.3395	0.7516	0.7516	0.00	1.0000	1.0179	+1.79
	0.6	0.4882	0.9205	0.9205	0.00	1.0000	0.9923	-0.77
	0.8	0.6183	1.0629	1.0626	-0.028	1.0000	1.0043	+0.43
	1.0	0.7353	1.1884	1.1887	+0.025	1.0000	0.9963	-0.37

\*ERROR % = ((PREDICTED)-(EXACT))/(EXACT)

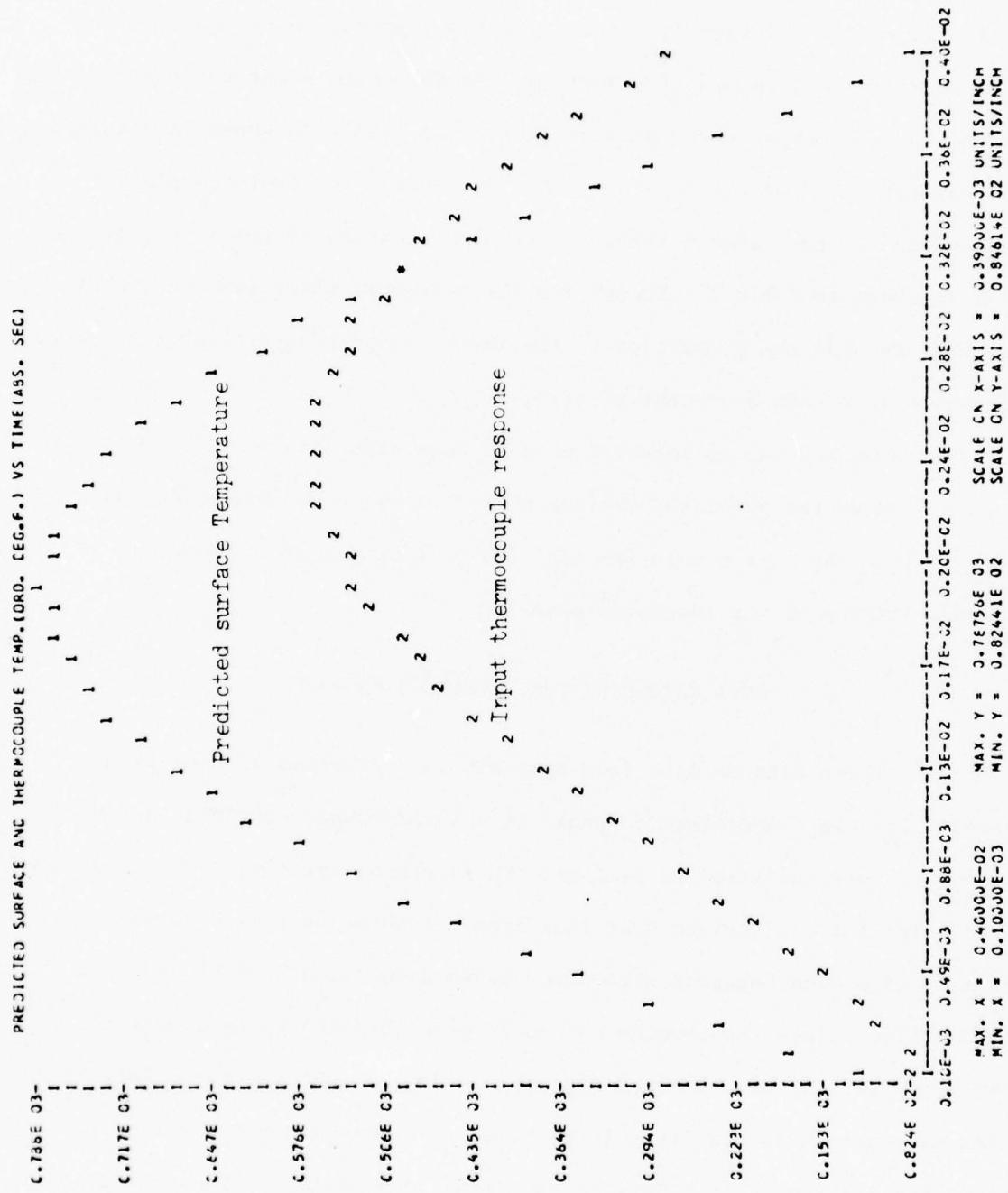
surface is subjected to a periodic temperature variation with a period of 8 milliseconds. Fifteen data points of the temperature response at the thermocouple location are then input to the inversion program for prediction of the surface temperature and heat flux. The result is shown in Figure 3 and Appendix C where the data symbol "2" denotes the thermocouple response and "1" the surface temperature. The accuracy of the inversion program is shown in Table 2. Except for the extremely short time period of 0.4 milliseconds the prediction by the inversion program with 15 term representation is within 2 percent of error.

The accuracy can be improved more if more data points are used.

Figure 4 shows the predicted surface heat flux which we were unable to compute from the series solution (Eq. (7) of Appendix D). This demonstrates the versatility of the inversion solution.

#### APPLICATION OF THE INVERSION PROGRAM

Three sets of data (see Appendix C) provided by Rock Island Arsenal for the temperature response of a thermocouple embedded in a M60 gun barrel were utilized to evaluate the inversion solution. The inversion prediction for the surface heat flux from all three sets of data were extremely high when compared with other known data calculated by Chen and Chiou [15]. Since the program correctly predicted the surface heat flux for other sets of experimental data it was judged that the three sets of data may contain inaccurate initial time. For most experimentations the recording instrument is likely to experience some delay in responding to the extremely fast transient heat flux typical in gun bores. Therefore an advanced shift of time of 1 millisecond in the data was tested. The



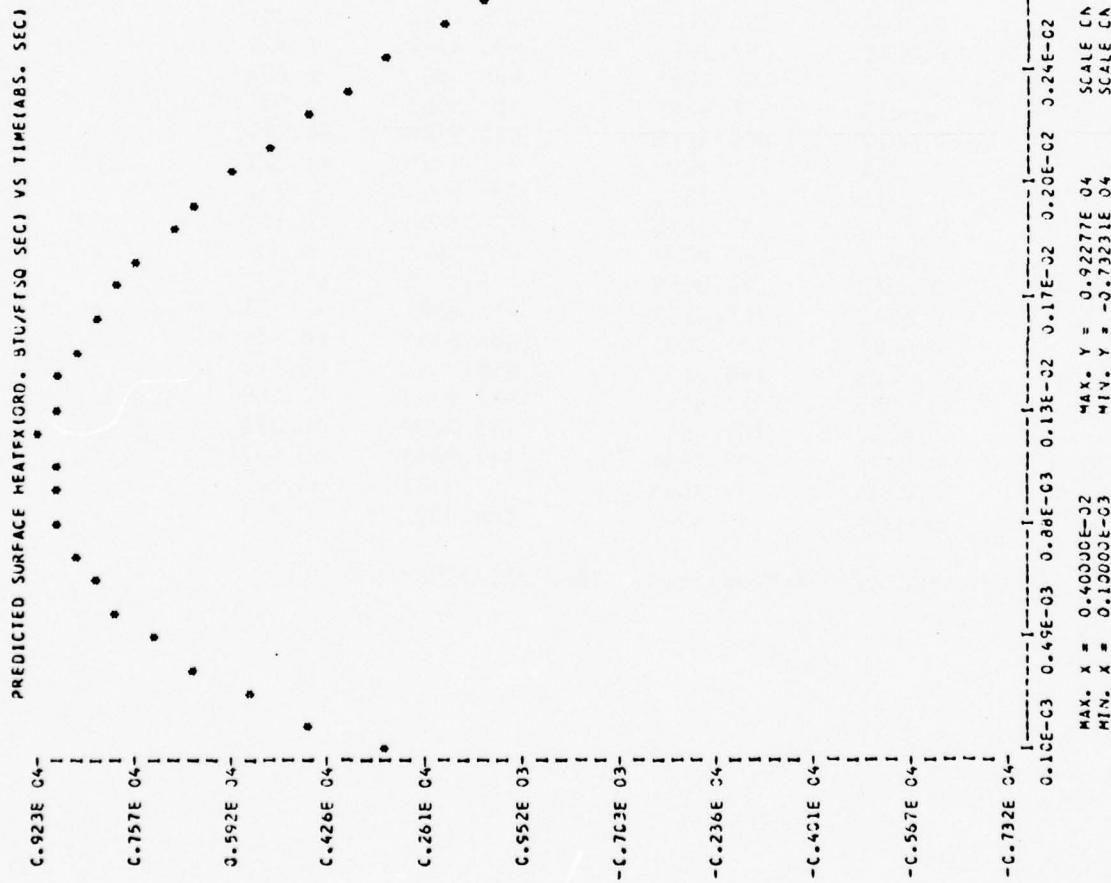
PART II Figure 3 Predicted Surface Temperature

## PART II Table 2

## Comparison of Inversion Prediction and Exact Solution

Time (sec.)	Surface Temperature (theoretical)	Surface Temperature (predicted)	Error* %
0.0002	189.5041	193.4060	+3.563
0.0004	296.3119	301.4311	+2.367
0.0006	397.7934	403.8998	+1.921
0.0008	491.4497	498.3988	+1.689
0.0010	574.9747	582.5683	+1.534
0.0012	646.3119	654.3358	+1.417
0.0014	703.7046	711.9396	+1.321
0.0016	745.7396	753.9650	+1.236
0.0018	771.3818	779.3802	+1.157
0.0020	780.0000	787.5615	+1.08
0.0022	771.3818	778.3093	+1.00
0.0024	745.7396	751.8527	+0.918
0.0026	703.7046	408.8443	+0.874
0.0028	646.3119	650.3441	+0.712
0.0030	574.9747	577.7933	+0.569
0.0032	491.6697	492.9790	+0.372
0.0034	397.7934	397.9893	+0.0616
0.0036	296.3119	295.2761	-0.679
0.0038	189.5041	188.8522	-0.595

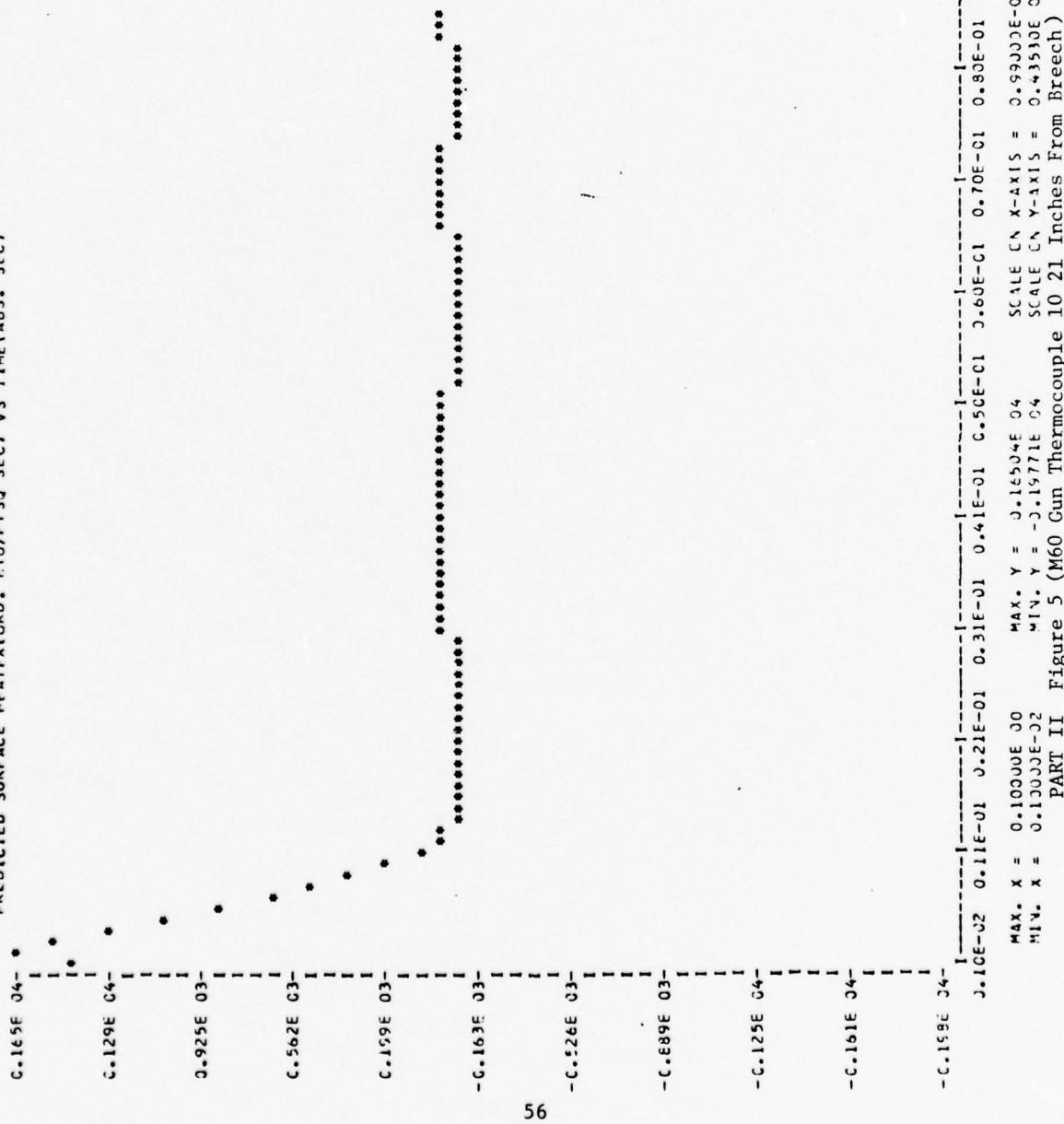
\*Error = ((Predicted - (Exact))/(Exact)) x 100



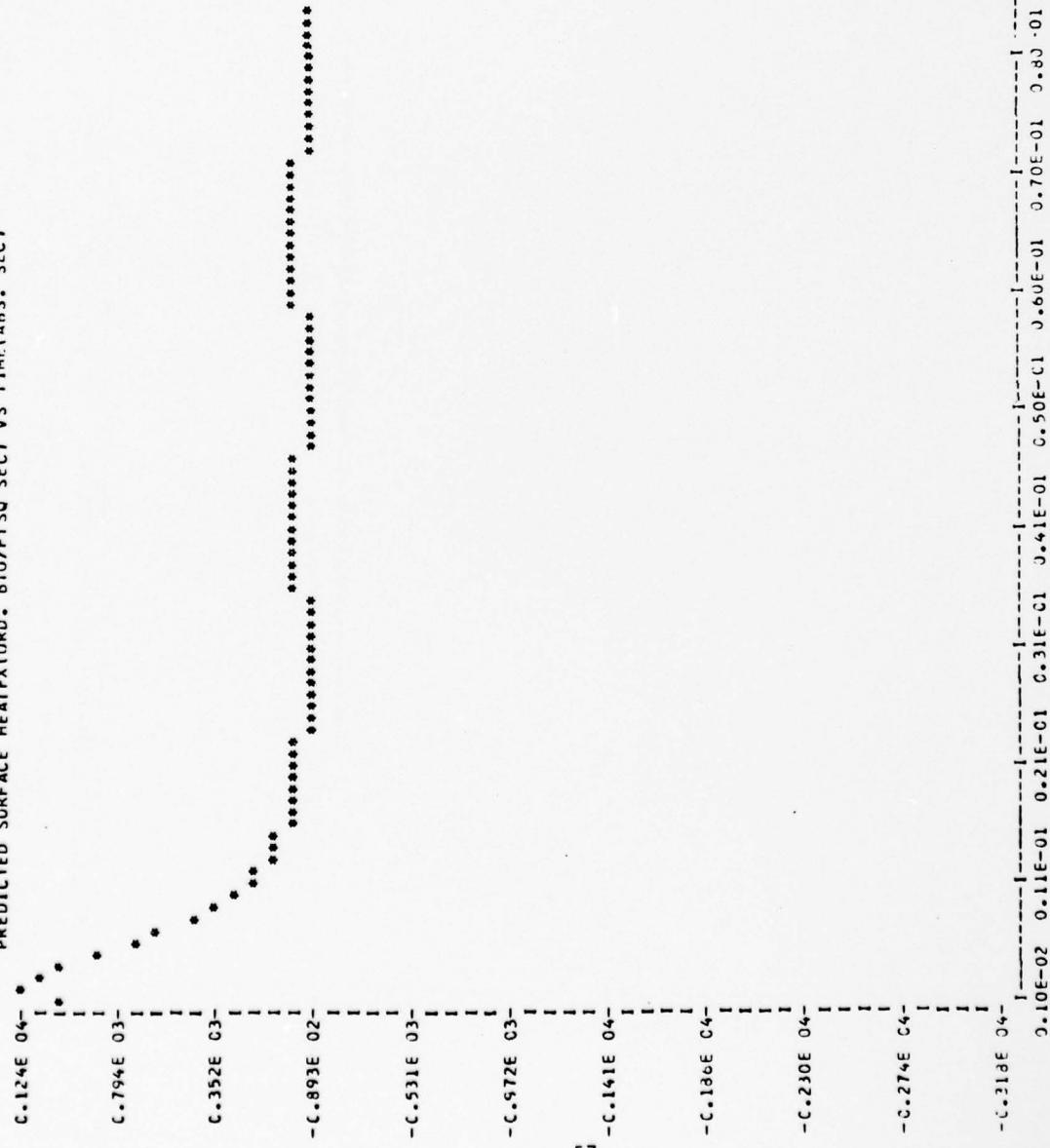
PART II Figure 4 Predicted Surface Heat Flux

result is shown in Appendix C-3. Figures 5, 6 and 7 respectively show the predicted maximum heat flux of 1650, 1235 and 2695 Btu/ft<sup>2</sup> sec. approximately at 2 milliseconds after the firing, while Figures 8, 9 and 10 show that the surface temperatures are 273, 234 and 396 degree F. approximately at 6 milliseconds after the firing.

PREDICTED SURFACE HEATFLUX (BTU/FT<sup>2</sup> SEC) VS TIME (ABS. SEC)



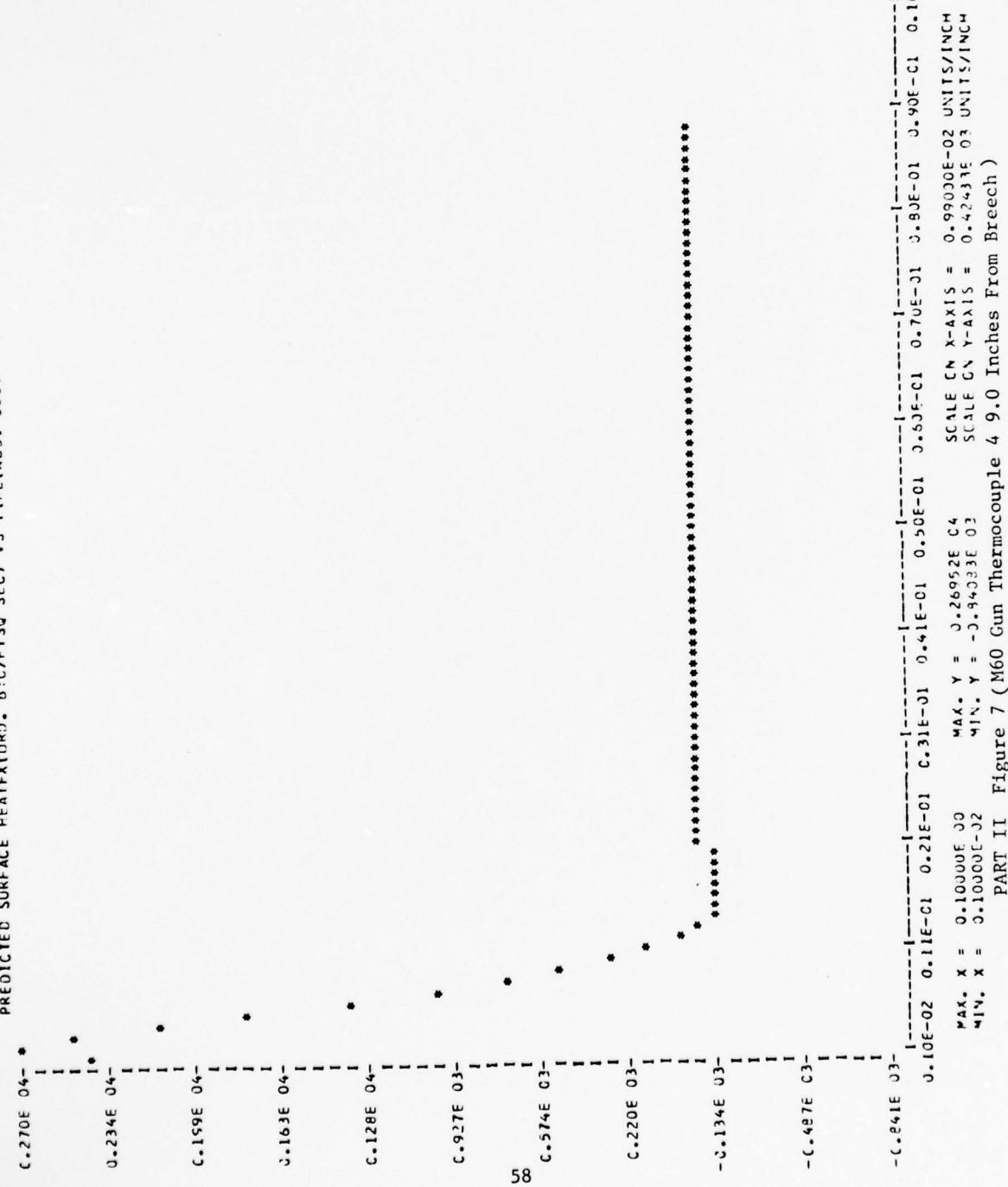
PREDICTED SURFACE HEATFLUX (W/D<sup>2</sup>) VS TIME (ABS. SEC.)



MAX. X = 0.10000E 00  
MIN. X = 0.10000E-02  
MAX. Y = 0.12354E 04  
MIN. Y = -0.31002E 04  
SCALE CN X-AXIS = 0.9910E-02 UNITS/INCH  
SCALE CN Y-AXIS = 0.527E 03 UNITS/INCH

PART II Figure 6 (M60 Gun Thermocouple 7 15.0 Inches From Breech)

PREDICTED SURFACE HEATFX(0RD. BYL/FTSQ SEC) VS TIME(ABS. SEC)

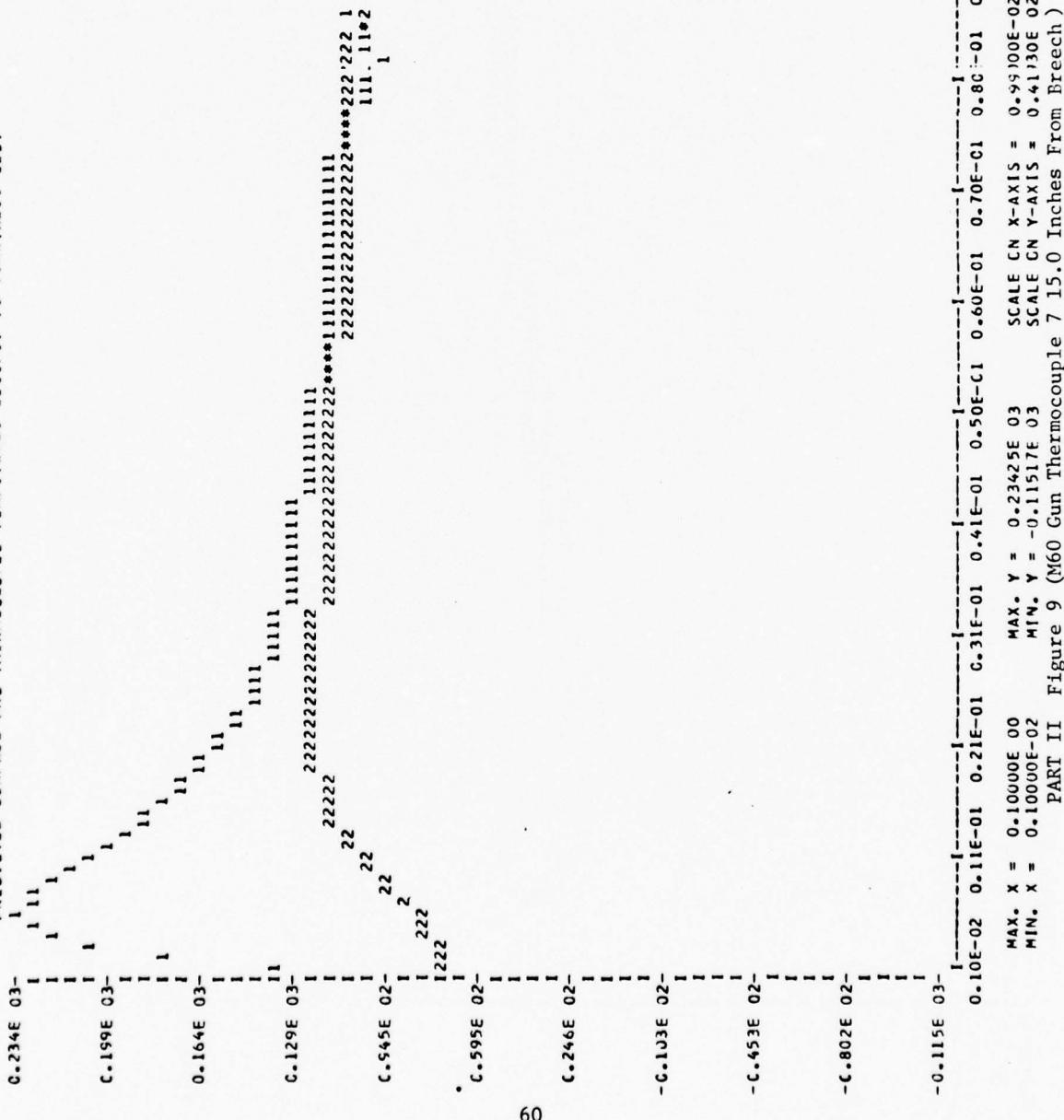


PART II Figure 7 (M60 Gun Thermocouple 4 9.0 Inches From Breech)

MAX. X = 0.10000E-09  
MIN. X = 0.10000E-32  
MAX. Y = 0.26952E-04  
MIN. Y = -0.94033E-03  
SCALE CN X-AXIS = 0.99000E-02  
SCALE CN Y-AXIS = 0.42431E-03 UNITS/INCH

PREDICTED SURFACE AND THERMOCHIMICAL TEMPERATURES VS TIME (A.S.T. SEC.)

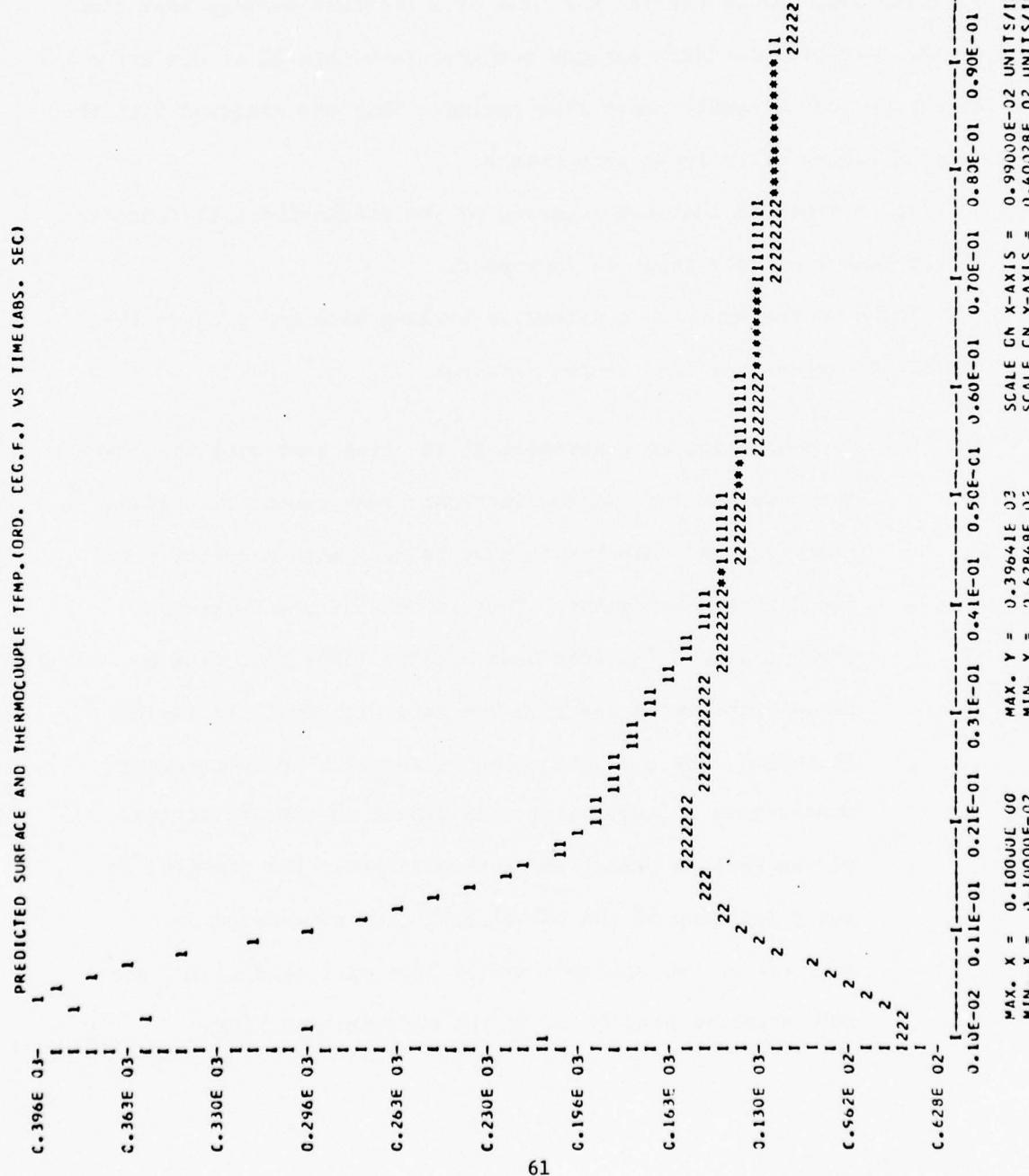
PREDICTED SURFACE AND THERMOCOUPLE TEMP.(ORD. DEC.F.) VS TIME (ABS. SEC)



MAX. X = 0.10000E 00  
MIN. X = 0.10000E-02  
MAX. Y = 0.23425E 03  
MIN. Y = -0.11517E 03  
SCALE CN X-AXIS = 0.9900E-02 UNITS/INCH  
SCALE CN Y-AXIS = 0.41130E 02 UNITS/INCH

PART II Figure 9 (M60 Gun Thermocouple 7 15.0 Inches From Breech)

PREDICTED SURFACE AND THERMOCOUPLE TEMP.(ORD. CEG.F.) VS TIME(ABS. SEC.)



PART II Figure 10 (M60 Gun Thermocouple 4 9.0 Inches From Breech)

## CONCLUSIONS AND SUGGESTIONS

The new inversion computer program was thoroughly tested for its applicability and accuracy. The program can invert an intrinsic temperature response to predict the case of a constant surface heat flux or the case of a periodic surface temperature within 2% of deviation except at the extremely short time period. This was achieved with the maximum number of 20 input data points.

It is expected that the accuracy of the prediction will increase if the number of data input is increased.

Based on the experience gained in working with the program the following suggestions are thought relevant.

- (1) In conducting an experiment it is vital that both the temperature and the time in the intrinsic measurement of surface heat flux and temperature must be much more accurate than the direct measurement. This is because the inversion problem always involves predicting a large heat flux or temperature variation from the data with small variation. Therefore, for a slight error in the time or temperature measurement a large error will result in the prediction of the surface heat flux or temperature. For example, in the prediction of gun barrel heat flux considered in an error of two milliseconds in time will lead to 100 percent error in prediction of the surface heat flux.

(2) In selecting the data points for input to the computer program care must be exercised not to create a locally abrupt jump in the data. An abrupt change in data points will often introduce an abnormal fitting of a curve in its neighborhood and hence resulting in an incorrect prediction of the surface heat flux and temperature. If indeed the abrupt jump of the data must be used, then more data points in its neighborhood must also be chosen.

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## APPENDIX A ANALYSIS OF THE INVERSION PROBLEM

Consider a slab, having a sufficient wall thickness,  $L$ , such that the outer surface temperature has a negligible response when the inner surface is exposed to a transient heat flux. A probe, for example a thermocouple, is located at  $X = X_1$  and it is normally desirable, as reported by Chen and Li [16], to be close to the heating surface since a better transient response and more accurate experimental measurements can be obtained to reduce error amplification in the mathematical inversion program. Under these circumstances we thus assume in the analysis  $L/X_1 \gg 1$ . The governing equation for the transient heat conduction may be written in a dimensionless form as

$$\frac{\partial \theta}{\partial t} = \frac{\partial^2 \theta}{\partial x^2} \quad (1)$$

with the initial and boundary conditions

$$\theta(x, 0) = 0 \quad (2)$$

$$\theta(\infty, t) = 0 \quad (3)$$

$$\theta(1, t) = f(t) \quad (4)$$

Where  $x = X/X_1$  is a dimensionless distance from the heating surface and  $t = \alpha\tau/X_1^2$  is a dimensionless time or Fourier number with  $\alpha$  being the thermal diffusivity and  $\tau$  the real time.  $\theta = (T - T_0)/T_0$  is the dimensionless temperature above the initially uniform temperature  $T_0$ .  $f(t) = (F(\tau) - T_0)/T_0$  is the dimensionless measured temperature response at  $x = 1$  with  $F(\tau)$  being the measured absolute temperature. The

inversion problem is then given the interior temperature  $f(t)$  to predict the surface temperature  $\theta(0,t)$  and the surface heat flux per unit area  $q(0,t)$  or  $\frac{\partial \theta}{\partial x}\Big|_0 = -q(0,t)x_1/T_o$ . Here  $\kappa$  is the thermal conductivity of the solid.

The above problem may be solved by Laplace Transformation. Let the transformation be:

$$\bar{\theta}(x,s) = \int_0^\infty \theta(x,t)e^{-ts} dt \quad (5)$$

where  $\theta$  is continuous otherwise satisfied the Dirichlet's condition. The temperature function  $\theta$  is recovered by inversion of the Laplace Transformation as:

$$\theta(x,t) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \bar{\theta} e^{st} ds \quad (6)$$

where  $c$  is a suitable positive value. Equation (1) and the initial condition (2) under transformation (5) become:

$$\frac{d^2\bar{\theta}}{dx^2} = s\bar{\theta} \quad (7)$$

which has a solution, with integration constants  $A$  and  $B$ ,

$$\bar{\theta}(x,s) = A e^{\sqrt{s}x} + B e^{-\sqrt{s}x} \quad (8)$$

The transformation of boundary conditions (3) and (4) into the Laplace

plane give  $\bar{\theta}(\infty, s) = 0$ ,  $\bar{\theta}(1, s) = \bar{f}(s)$ . Substitution of the boundary conditions into equation (8), we get  $\bar{\theta}$  and its derivative as:

$$\bar{\theta}(x, s) = \bar{f}(s) e^{\sqrt{s}(1-x)} \quad (9)$$

$$-\frac{\partial \bar{\theta}}{\partial x} = \sqrt{s} \bar{f}(s) e^{\sqrt{s}(1-x)} \quad (10)$$

According to Chen and Thomsen [14], we choose the temperature response as measured by a probe to be represented by a polynomial

$$f(t) = \sum_{n=1}^N b_n (4t)^n \Gamma(n+1) i^{2n} \operatorname{erfc} \left( \frac{1}{2\sqrt{t}} \right) \quad (11)$$

or in Laplace plane

$$\bar{f}(s) = e^{-\sqrt{s}} \sum_{n=1}^N \Gamma(n+1) \frac{b_n}{s^{1+n}} \quad (12)$$

The  $b_n$ 's are coefficients of the expansion to be determined such that the  $N$  term polynomial describes the temperature response  $f(t)$  measured at  $x = 1$ . With equation (12), equations (9) and (10) can be simplified to

$$\bar{\theta}(x, s) = e^{-x\sqrt{s}} \sum_{n=1}^N \Gamma(n+1) \frac{b_n}{s^{1+n}} \quad (13)$$

$$-\frac{\partial \bar{\theta}}{\partial x} = e^{-x\sqrt{s}} \sum_{n=1}^N \Gamma(n+1) \frac{b_n}{s^{1/2+n}} \quad (14)$$

The inversion of equations (13) and (14) at  $x = 0$  give:

$$\theta(0, t) = \sum_{n=1}^N b_n t^n \quad (15)$$

$$-\frac{\partial \theta(x, t)}{\partial x} \Big|_{x=0} = \sum_{n=1}^N b_n t^{n-1/2} \frac{\Gamma(n+1)}{\Gamma(n+1/2)} \quad (16)$$

where  $\theta(0, t)$  gives the surface temperature and  $-\frac{\partial \theta(0, t)}{\partial x}$  gives the surface heat flux  $\frac{qX_1}{\kappa T_0}$  as a function of time.

The integral of the error function in (11) is defined as:

$$i^{2n} \operatorname{erfc}\left(\frac{1}{2\sqrt{t}}\right) = \int_{\left(\frac{1}{2\sqrt{t}}\right)}^{\infty} i^{2n-1} \operatorname{erfc}(y) dy \quad (17)$$

with  $n = 0$ ,  $\operatorname{erfc}(y) = \frac{2}{\sqrt{\pi}} \int_y^{\infty} e^{-x^2} dx$ .  $\Gamma(n)$  in equations (11) and (16) is the gamma function or Euler's integral function of the second kind.

$$\Gamma(n) = \int_0^{\infty} e^{-\omega} \omega^{n-1} d\omega \quad (18)$$

It should be remarked that the choice of the particular form (11) is to ensure the convergence of the solution on the Laplace plane and an analytic inversion back to the physical plane. With  $b_n$  coefficients determined from equation (11) and the experimental measurement of the temperature response  $f(t)$  at  $x = 1$ , the surface temperature,  $T_w(t)$ , is obtained from equation (15) as

$$T_w(t) = T_0 \left( 1 + \sum_{n=1}^N b_n t^n \right) \quad (19)$$

and the heat flux,  $q(t) = \frac{-T_o \kappa}{X_1} \left. \frac{\partial T}{\partial x} \right|_{x=0}$ , from equation (16) as:

$$q(t) = \frac{-T_o \kappa}{X_1} \sum_{n=1}^N b_n t^{n-1/2} \frac{\Gamma(n+1)}{\Gamma(n+1/2)} \quad (20)$$

The above solution is the exact solution for predicting the transient surface heat flux and temperature valid for both short and long time durations as long as the slab is thick enough such that the outer surface maintains its initial temperature. The feature of the present solution is the polynomial (11) which on the Laplace plane gives a term in equation (12)  $\exp(-\sqrt{s})$  to cancel the term  $\exp(\sqrt{s})$  in equations (9) and (10). This polynomial (11) as suggested by Chen and Thomsen [15] makes the present solution simpler than many inversion solutions derived in the past and valid in any time.

If the fluid temperature,  $T_g(t)$ , away from the surface of the slab is known, then the instantaneous heat transfer coefficients,  $h(t)$ , can be determined from Newton's cooling law as:

$$h(t) = \frac{q(t)}{T_g(t) - T_w(t)} \quad (21)$$

where heat flux,  $q(t)$ , and wall temperature  $T_w(t)$  are given by equations (19) and (20), and the average heat transfer coefficient up to time  $t$ ,  $\bar{h}(t)$ , can be defined as:

$$\bar{h}(t) = \frac{\int_0^t q(t') dt'}{\int_0^t [T_g(t') - T_w(t')] dt'} \quad (22)$$

APPENDIX B

IMPROVED COMPUTER PROGRAM

Cartesian Inversion Problem

```

FORTRAN IV G LEVEL 21          DATE * 76157          PAGE 0001
                                15/33/00

C      *** DATA INPUT ***
C
C      TITLE OF DATA(BLANK FIRST 20 SPACE)          FIRST CARD
C      ACRD=8CRE RADIUS OF CYLINDER (FT.)          SECOND CARD
C      OUTR= OUTER RADIUS (FT.)                   THIRD CARD
C
C      THICK=OUTER RADIUS-BURRED (FT.)          FOURTH CARD
C      DIS=ACRF TO THERMOCOUPLE DISTANCE (FT.)      FIFTH CARD
C
C      TEMPO=INITIAL THERMOCOUPLE TEMP (ATREF.)      SIXTH CARD
C      TG(1)= INITIAL GAS CORE TEMPERATURE AT NORM TIME 0.00001
C
C      ALP= THERMAL DIFFUSIVITY OF MATERIAL (FT2/SEC)      SEVENTH CARD
C      THCCN= THERMAL CONDUCTIVITY OF MATERIAL(BTU/FT, SEC, DEG, F)
C
C      NP=NUMBER OF TIME-TEMPERATURE DATA PAIRS UP TO 25 PAIRS
C
C      NB= NUMBER OF B COEFF. TO BE FITTED          EIGHTH CARD
C      FROM L1TH (N AR PAIRS OF TIME(L1) AND TEMP(L1))      NINTH CARD
C
C      TIME(L1)=TIME FROM THERMOCOUPLE CURVE(SEC.)      TENTH CARD
C      TEMP(L1)= TEMPERATURE FROM THERMOCOUPLE RESPONSE (DEG. F)
C      TSHTFT= TIME SHIFTED FROM DATA READ (NSEC)      LAST CARD
C
C
C      DECLARATION STATEMENTS      START COMPUTER PROGRAM
C
C
C      IMPLICIT REAL*8(A-H,O-Z)
C      DIMENSION TG(1210),QM(1210),DTWH(1210),HM(1210),TG(1210),
C      LTWH(1210),TOT(1210),TEM0(1210),HT(1210),EQP(1210),DTEMR(1210)
C      DIMENSION B(126),REFC(52),TEMP(26),TIME(26),T(26),X(126),GM(128),
C      IC(125,26),FN(26),
C      G(126,1210),H(126,1210),EQUE(1210),DTEMP(1210)
C      2,TC(26),TEM(1210),TRN(126,26),Q(26,26)
C      DATA EPS/.10-15/
C
C      READ DATA CARDS INTO PROGRAM
C
C
C      CALL ERSET(208,0,-1,1,1)
C      199  FCRMAT( 20X,50H
C
C      0002  1
C      200  FCRMAT(120,10)
C
C      201  FORGET(121)
C
C      1  REFO(5,119,END=9991)
C      READ(5,200) QRRAD
C      READ(5,200) C1TR
C      READ(5,200) DIS
C      READ(5,200) TEMPO
C      READ(5,200) TG(1)
C      READ(5,200) ALP
C      READ(5,200) TMCN
C
C      2  READ(5,201) NB
C
C      3  READ(5,201) NB
C      WRITE(6,198)
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C      0003  1
C      198  WRITE(6,199)
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FCRTRAN IV G LEVEL 21          MAIN          DATE = 76157      15/33/00
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C056          DO 40 J=1,N8
C055          J2=2+J
C056          C11=J*(4.000*T(111)*J*HERFC(J2)*0GAMMA(J+1.000)
C060          40
C061          CONTINUE
C
C-----NOW B COEFFICIENT FOR THE TEMPERATURE RESPONSE AT
C-----THE THERMOCPILE ARE SOLVED EITHER BY EXACT OR LEAST SQUARE
C
C062          NT=NP+1
C063          IF (NP = N8) 42, 43, 42
C064          CALL LTSQR (NP,N8,C,B,EPS,TIME,NT,TRAN,Q1Q,TEMPO)
C065          GO TO 44
C066          43  DET = SIMUL (NP, C, B, FPS, 1,11)
C067          44  CCNTINUE
C
C-----NOW CALCULATE HEAT FLUX FUNCTION AND
C-----CALCULATE TEMPERATURE FUNCTION
C
C068          TEM = TEMIN
C069          DO 55 I=1,200
C070          DO 50 J=1,N8
C071          G(J,I) = TEM**J
C072          H(J,I) = -TEM**((J-0.5)*0.1*0GAMMA(J+1.000)/0GAMMA(J+0.500)
C073          50  CONTINUE
C074          TEM = TEM + TEMIN
C075          55  CCNTINUE
C
C-----CALCULATE TEMPERATURE AND HEAT FLUX
C
C076          DO 75 I=1,200
C077          DELT=0.000
C078          QUE=0.000
C079          DO 70 J=1,N8
C080          DELT=DELT*(1(J)*G(J,I))
C081          QUE=QUE+R(J)*H(J,I)
C082          70  DTEMP(I)=DELT
C083          DTEMP(I)=DTEMP(I)*(460.000+TEMPO)+TEMPO
C084          EQUE(I)=QUE
C085          75  EQUF(I)=EQF(I)*(460.000+TEMPO)+THCQN/DIS
C
C-----CALCULATE INST. AND MEAN QUANTITIES FOR TEMP. AND HEAT T.COEF.
C
C086          T(111)=0.000
C087          TGM(11)=0.000
C088          QM(11)=0.000
C089          DTWM(11)=0.000
C090          TW(11)=0.000
C091          HM(11)=0.000
C092          SUMTM=0.000

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FCRTRAN IV G LEVEL 21          LTSQER          DATE = 76157      15/33/00
0001      SUBROUTINE LTSQER (M, JJ, PHI1, B, EPS, TYPE, NT, TRAN, QTQ, TEMP0)
C
C   THIS PROGRAM CALCULATES THE LEAST SQUARES COEFFICIENTS FOR
C   MODIFIED POLYNOMIAL WHICH IS POLYNOMIAL TIMES THE INTEGRATED
C   ERROR FUNCTION
C
C   M IS THE NUMBER OF DATA PAIRS TO BE USED
C   JJ IS THE NUMBER OF DEGREE TO BE FITTED
C
C
0002      REAL*8 PHI1(26,26), FX(26), TRAN(26,26), QTQ(26,26), X(26), DET, FPX,
          X(126), TIME(26),
          CALL SQRSR(1208,0,-1,1,1)
          LD1P=50
          DO 1001 I=1,M
          N1=JJ+1
          X(1)=TIME(1)
          X(1)=TIME(1)
          FX(1)=PHI1(1,N1)
          1001 CCNTNJE
          ERROR =0.000
          C
C   FIND THE TRANSPOSE OF THE PHI1 MATRIX FOR LTSQ FITTING AND
C   PERFORM THE MULTIPLICATION AND GET JJ BY JJ MATRIX
C
0011      CALL TRANS (PHI1, TRAN,M, JJ, NT)
          CALL MULT (PHI1, TRAN,JJ, M, QTQ, NT)
          DO 4 J=1,JJ
          QTQ(J,N1)=0.000
          DO 5 K=1,M
          QTQ(J,N1)=QTQ(J,N1) + TRAN(J,K)*FX(K)
          4      WRITE (6,5) (I,QTQ(I,N1),I=1,JJ)
          5      FORMAT (I=1,12,*, QTQ(I,N1)=,016.7)
          C
C   FIND THE SOLUTION TO THE NORMAL EQUATIONS
C
0015      DET = SIMUL (JJ, QTQ,B,FPSS,1,20)
          WRITE (6,100)
          100  FORMAT (I, LEAST SQUARE COEFFICIENTS ')
          100  WRITE (6,6) (I,A(I), 'I=1,JJ)
          6      FORMAT (I= ,12, , B(I)= , 016.7)
          6      WRITE (6,9)
          9      FORMAT (I, LEAST SQ CHIT PUT ,//, KTH TIME', T15, 'TEMP. GIVEN ',
          1      LTSQ, 'FITTED', T42, 'ERRCR', T55, 'PER CENTAGE')
          DC 15 K=1,M
          FPX=TEMP3
          DC 8 I=1,JJ
          8      FPX=FPX + QTQ(I,N1)* PHI1(K,I)
          8      FX(K)= FX(K) + TEMP0
          8      FPR = PHI1(K,N1) - FPX + TEMP0
          8      IF (FPR(K,N1).NE. 0.) PER=FPR/(PHI1(K,N1)) + TEMP0*100.0
          8      IF (ABS(ER0).LE. ERRCR) GO TO 15
          8      ER0 = ABS(ER0)

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FORTRAN IV G LEVEL 21  
 FUNCTION SIMULIN,A,X,EPS,INDIC,NRC  
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C
C 1 ***GUSS JORDAN TECHNIQUE MAXIMUM PIVOT
C 1 ***PAGE 290 'APPLIED NUMERICAL METHODS' CARRAHAN
C 1
C
C002      IMPLICIT REAL*8(A-H,O-Z)
C003      DIMENSION IROW(60),JCOL(60),JORD(60),Y(60),A(26,26),X(26)
C004      CALL ERSET(208,0,-1,1,1)
C005      MAX=N
C006      IF(INDIC.EQ.0) MAX=N+1
C007      C ***IS N LARGER THAN 50
C008      IF(N.LE.50) GO TO 5
C009      WRITE(6,200)
C010      SIMUL=0.0000
C011      RETURN
C012      C BEGIN ELIMINATION PROCEDURE
C013      C DETER=1.0000
C014      C PIVCT=0.0000
C015      DO 11 I=1,N
C016      DO 11 J=1,N
C017      C ***SCAN IROW AND JCCL ARRAYS FOR INVALID PIVCT SUBSCRIPTS
C018      IF(K.EQ.1) GO TO 9
C019      DO 8 ISCAN=1,KM1
C020      IF(I.EQ.IROW(ISCAN)) GO TO 11
C021      IF(J.EQ.JCCL(ISCAN)) GO TO 11
C022      CONTINUE
C023      IF(CASSIA(I,J).LE.CABS(PIVOT)) GO TO 11
C024      PIVCT=A(I,J)
C025      IROW(K)=I
C026      JCCL(K)=J
C027      C CONTINUE
C028      C ***INSURE THAT SELECTED PIVOT IS LARGER THAN EPS
C029      IF(CABS(PIVOT).GT.EPS) GO TO 13
C030      SIMUL=0.0000
C031      C UPDATE THE DETERMINANT VALUE...
C032      C PIVCM=(PIVOT)
C033      C JCCLK=JCCL(K)
C034      C DETER=DETER*PIVOT
C035      C ***NORMALIZE PIVOT ROW ELEMENTS
C036      DO 14 J=1,MAX
C037      A(I,JCCLK,J)=A(I,JCCLK,J)/PIVOT
C038      JCCLK=JCCL(K)
C039      C ***CARRY OUT ELIMINATION AND DEVELOPE INVERSE
C040      DO 18 I=1,N
C041      A(JCK,A(I),JCCLK)=1./PIVOT
C042      IF(I.EQ.IROW(K)) GO TO 18
C043      IF(J.NE.JCCLK) A(I,J)=A(I,J)-A(JCK,A(I,JCCLK),J)
C044      C CONTINUE
C045      C ***ORDER SOLUTION VALUES AND CREATE JORD ARRAY
  
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```
0044      DO 20 I=1,N
0045      IROW1=IROW(I)
0046      JCOL1=JCOL(I)
0047      JCRL1=IPROW(I)-JCOL1
0048      IF(JCRL1.GE.0) X(JCCL1)=A(IROW1,MAX)
0049      C      ADJUST SIGN OF DETERMINANT
0050      INTCH=0
0051      NP1=N-1
0052      DO 22 I=1,NM1
0053      IP1=I+1
0054      DO 22 J=IP1,N
0055      IF(JCRL1.GE.0) JCRL1=J-TEMP
0056      JCRL1=JCRL1
0057      JCRL1=JCRL1
0058      INTCH=INTCH+1
0059      CCNTINUE
0060      C      IF(INTCH/2*2.NE.INTCH) DETER=DETER
0061      C      IF(INDIC.LE.0) GC TC 26
0062      SIMUL=DETER
0063      RETURN
0064      C      ....IF INDIC IS NEGATIVE OR ZERO, UNSCRAMBLE THE INVERSE
0065      C      ...FIRST BY ROWS
0066      DO 29 J=1,N
0067      IROW1=IROW(I)
0068      JCCL1=JCCL(I)
0069      Y(JCCL1)=A(IROW1,J)
0070      DO 28 I=1,N
0071      A(I,J)=Y(I)
0072      C      ...THEN BY COLUMNS
0073      DO 30 J=1,N
0074      JCCLJ=JCCL(J)
0075      Y(IROW1,J)=A(I,J)
0076      DO 30 J=1,N
0077      A(I,J)=Y(I,J)
0078      C      ....RETURN FOR INDIC NEGATIVE OR ZERO
0079      SIMUL=DETER
0080      RETURN
0081      C      ....FORMATIION TOO BIG
0082      CDEBUG SUBCHK
0083      END
0084
0085
0086
0087
0088
0089
0090
0091
0092
```

## APPENDIX C

### NUMERICAL RESULTS

- Appendix C - 1      Constant Heat Flux Case
- Appendix C - 2      Periodic Surface Temperature Case
- Appendix C - 3      Gun Barrel (M60) Heating

## APPENDIX C - 1

## FLAT STEEL PLATE EXPOSED TO CONST HEAT FLUX

INCHES RADIUS (FT.) = 0.05089  
 OUTER RADIUS (FT.) = 0.09494  
 CORE TO THERMOCOUPLE DISTANCE (FT.) = 0.004405  
 INITIAL THERMOCOUPLE TEMPERATURE (F.) = 67.0000  
 INITIAL GAS TEMPERATURE (F.) = 4.6167  
 THERMAL DIFFUSIVITY (FT<sup>2</sup>/SEC) = 0.000010307  
 THERMAL CONDUCTIVITY (BTU/FT<sup>2</sup> SEC. F.) = 0.0555555  
 NUMBER OF TIME TEMPERATURE PAIRS (SEC., F.) = 10  
 NUMBER OF B(I) CCF. TO BE FITTED = 10

TIME	TEMPERATURE
0.0188260648	69.1882892600
0.0376521296	84.0565859500
0.0564781544	106.9027670000
0.0753047592	132.7351606000
0.0941303241	156.4852637000
0.1129561889	186.2713064000
0.1317824537	212.0382040000
0.1506085185	238.6058830000
0.1694345933	263.9195910000
0.1882606481	288.6139635000
TIME OF DATA SHIFTED BY (SEC)	0.0
COEFFICIENTS OF B(I) = 0.124691070	I = 1
COEFFICIENTS OF B(I) = -0.231739660	I = 2
COEFFICIENTS OF B(I) = 0.257499150	I = 3
COEFFICIENTS OF B(I) = -0.164626540	I = 4
COEFFICIENTS OF B(I) = 0.643047000	I = 5
COEFFICIENTS OF B(I) = -0.158560170	I = 6
COEFFICIENTS OF B(I) = 0.247575850	I = 7
COEFFICIENTS OF B(I) = -0.237106540	I = 8
COEFFICIENTS OF B(I) = 0.126933550	I = 9
COEFFICIENTS OF B(I) = -0.250550150	I = 10

FLAT STEEL PLATE EXPOSED TO CONST. HEAT FLUX  
TEMPERATURE AND HEAT FLUX  
PREDICTED BY JFNCSHING CHAU

( $10^3$ RT <sub>m</sub> ) ( $10^4$ F <sub>s</sub> )	(REAL T <sub>m</sub> ) (SEC.)	(INR. FLUX) (DELESSI) (BTU/FT <sup>2</sup> SEC)	(REAL FLUX) (BTU/FT <sup>2</sup> SEC)	(NOR.) (TEMPER)	(REAL, SUF. T) (DEG. F.)	(TH, CO, T) (DEG. F.)
C.005212	0.0010000	0.8997	597.9591	0.0601	98.6557	67.00000000
C.010624	0.0020000	1.1185	743.4044	0.1092	124.5475	67.00000000
U.015335	0.0030000	1.2376	802.6189	0.1493	145.6677	67.00000112
0.021247	0.0040000	1.2336	819.5364	0.1819	162.8741	67.00000116
0.026555	0.0050000	1.2257	814.6777	0.2085	176.9046	67.00019875
0.031611	0.0060000	1.1597	757.3942	0.2333	188.3900	67.00140603
0.031822	0.0070000	1.1652	776.4222	0.2483	197.8656	67.00551487
C.042454	0.0080000	1.1280	749.7190	0.2633	205.7826	67.0178594
0.04786	0.0090000	1.0920	725.7677	0.2761	212.5174	67.04391977
0.053118	0.0100000	1.0593	704.0854	0.2873	218.3813	67.03817455
0.058430	0.0110000	1.0314	685.5339	0.2972	223.6281	67.16032801
J.063141	0.0120000	1.0088	670.5237	0.3064	228.6624	67.26603032
0.065053	0.0130000	0.9917	659.1532	0.3151	233.9656	67.41050573
0.07365	0.0140000	0.9799	651.3071	0.3233	237.5018	67.59939341
0.075677	0.0150000	0.9730	646.7278	0.3319	241.5249	67.83465420
0.085858	0.0160000	0.9705	635.0575	0.3301	246.3816	68.11362713
0.090360	0.0170000	0.9718	645.4273	0.3493	250.5167	68.45215395
0.095612	0.0180000	0.9763	646.8856	0.3573	255.5569	68.83515341
0.100524	0.0190000	0.9833	653.5161	0.3689	260.2143	69.2688168
J.106336	0.0200000	0.9921	655.4201	0.3761	265.1852	69.74576862
J.111547	0.0210000	1.0023	665.2036	0.3855	270.1730	70.27021864
J.115559	0.0220000	1.0133	673.5161	0.3932	275.2500	70.83331522
0.122171	0.0230000	1.0247	681.0502	0.4049	280.4001	71.44779410
0.127383	0.0240000	1.0359	698.5258	0.4119	285.6000	72.09597750
0.132455	0.0250000	1.0256	695.1120	0.4247	290.2447	72.1834613
J.138106	0.0260000	1.0568	702.4168	0.4366	296.0.86	72.50511570
0.143418	0.0270000	1.0660	708.4977	0.4445	301.2467	74.26058355
0.149130	0.0280000	1.0780	713.6059	0.4543	306.3953	75.04802121
0.155042	0.0290000	1.0807	718.3033	0.4639	311.1723	75.86538472
0.159354	0.0300000	1.0862	721.9203	0.4734	316.5882	76.11315524
0.164465	0.0310000	1.0903	724.6418	0.4826	321.3355	77.58836333
C.165577	0.0320000	1.0930	726.4744	0.4916	326.0913	78.49055763
0.175285	0.0330000	1.0945	727.4469	0.5004	330.7131	79.4181618
C.185651	0.0340000	1.0957	727.6067	0.5039	335.1930	80.37225643
0.185512	0.0350000	1.0950	727.0167	0.5171	335.5235	81.35923020
J.191224	0.0360000	1.0919	725.7520	0.5251	343.7082	82.35193129
C.195336	0.0370000	1.0891	723.8971	0.5327	347.7079	82.37662779
C.201649	0.0380000	1.0856	721.5428	0.5401	351.6262	84.42374619
0.207160	0.0390000	1.0815	718.7638	0.5472	355.2681	85.58251826
J.212171	0.0400000	1.0768	715.7159	0.5540	358.5742	86.58251826
0.217793	0.0410000	1.0719	712.4246	0.5616	362.5117	87.69282777
0.222355	0.0420000	1.0668	700.0324	0.5570	365.8194	88.92282119
0.228457	0.0430000	1.0616	705.5576	0.5732	368.0105	89.97173463
C.233715	0.0440000	1.0565	702.2130	0.5792	372.2146	91.13866154
0.239330	0.0450000	1.0516	698.5540	0.5850	375.2333	92.32357554
0.244342	0.0460000	1.0470	695.3885	0.5907	378.2194	93.52482037
C.247354	0.0470000	1.0428	693.0758	0.5962	381.2151	94.74187828
0.254666	0.0480000	1.0390	680.5659	0.6017	384.1024	95.57351049
C.262727	0.0490000	1.0357	688.4955	0.6071	386.5259	97.22005800
J.265885	0.0500000	1.0330	686.6C77	0.6125	389.7778	98.4679.9701
C.270501	0.0510000	1.0309	685.2120	0.6178	392.5873	99.75134538
0.276213	0.0520000	1.0295	684.2247	0.6231	395.3909	101.93476913

## CONSTANT HT FLUX

PORE RADIUS (FT.) = 0.05089  
 CUTER RADIUS (FT.) = 0.09594  
 BORE TO THERMOCOUPLE DISTANCE (FT.) = 0.004405  
 INITIAL THERMOCOUPLE TEMPERATURE (F.) = 67.0000  
 INITIAL GAS TEMPERATURE (F.) = 4.6167  
 THERMAL DIFFUSIVITY (FTSQ/SEC) = 0.00610307  
 THERMAL CONDUCTIVITY (BTU/FT. SEC.F.) = 0.00555555  
 NUMBER OF TIME TEMPERATURE PAIRS (SEC., F.) = 20  
 NUMBER OF ELLI COFF. TO BE FITTED = 20

TIME	TEMPERATURE
0.0188280648	69.1882892400
0.0376521296	84.0565859500
0.0564781544	106.5027670000
0.0753042592	122.7352606000
0.0941303241	159.4852627000
0.1129538889	186.27130404000
0.1317824537	212.70382046000
0.1506085185	239.60988300000
0.1694345833	263.51955100000
0.1882606481	288.61396350000
0.2070861129	312.65966650000
0.2256227777	336.15607540000
0.24447388425	359.12845740000
0.2635649073	381.52456160000
0.2823509722	403.41285100000
0.3012170370	424.81912890000
0.3200431018	445.77081370000
0.3388691666	466.22019330000
0.357692314	486.40535900000
0.3765212562	506.13253210000
TIME OF DATA SHIFTED BY (SEC)	0.0
COEFFICIENTS OF ELLI = 0.136910520	0.2
COEFFICIENTS OF ELLI = -0.326635100	0.3
COEFFICIENTS OF ELLI = 0.504844960	0.4
COEFFICIENTS OF ELLI = -0.485537750	0.5
COEFFICIENTS OF ELLI = 0.320658150	0.6
COEFFICIENTS OF ELLI = -0.149504610	0.7
COEFFICIENTS OF ELLI = 0.519634380	0.7
COEFFICIENTS OF ELLI = -0.137164890	0.8
COEFFICIENTS OF ELLI = 0.2808115540	0.8
COEFFICIENTS OF ELLI = -0.451360470	0.8
COEFFICIENTS OF ELLI = 0.574655460	0.8
COEFFICIENTS OF ELLI = -0.574052080	0.8
COEFFICIENTS OF ELLI = 0.160581160	0.8
COEFFICIENTS OF ELLI = -0.28437950	0.8
COEFFICIENTS OF ELLI = 0.160593480	0.8
COEFFICIENTS OF ELLI = -0.519558940	0.7
COEFFICIENTS OF ELLI = 0.140826370	0.7
COEFFICIENTS OF ELLI = -0.26396520	0.6
COEFFICIENTS OF ELLI = 0.305176370	0.5
COEFFICIENTS OF ELLI = -0.164086340	0.4
(COEFFICIENTS OF ELLI =	1 = 20

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IOWA INST OF HYDRAULIC RESEARCH IOWA CITY  
HEAT TRANSFER ANALYSIS FOR UNSTEADY HIGH VELOCITY PIPE FLOW. (U)  
APR 77 C J CHEN, P LI, J S CHIOU, H Y LEE

F/G 20/13

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CONSTANT HT FLUX  
TEMPERATURE AND HEAT FLUX  
PREDICTED BY JENG-SHING CHIOU

(INR.T)	(REAL HT)	(NOR. FLUX)	(REAL,FLUX)	(NOR.)	(REAL,SUF.T)	(TH,CO,T)
(DMFLSS)	(SEC.)	(DMLESS)	(BTU/FTSQ SEC)	(DMLESS)	(DEG.F.)	(CEG,F.)
0.005212	0.0010000	0.9534	633.6531	0.0642	100.3476	67.00000000
0.010624	0.0020000	1.187	763.5597	0.1141	127.1067	67.00000000
0.015535	0.0030000	1.207	802.7108	0.1528	147.5288	67.00000013
0.021267	0.0040000	1.2031	802.9512	0.1831	165.5185	67.00001257
0.026555	0.0050000	1.1823	785.8116	0.2072	176.1928	67.00021428
0.031871	0.0060000	1.1967	762.1832	0.2266	188.4310	67.00150476
0.037182	0.0070000	1.1102	737.8673	0.2427	198.5175	67.00628848
0.042456	0.0080000	1.0771	715.8933	0.2565	202.1776	67.0188639
0.047800	0.0090000	1.047	697.6728	0.2587	208.6083	67.04512806
0.053118	0.0100000	1.0287	683.6940	0.2799	218.5034	67.09185289
0.058130	0.0110000	1.0139	673.8239	0.2505	220.0748	67.16650319
0.063741	0.0120000	1.0048	667.8232	0.3007	227.82707	67.27386016
0.069553	0.0130000	1.0005	664.9527	0.3108	230.789%	67.42063610
0.074265	0.0140000	1.0002	664.7808	0.3209	230.015	67.61041071
0.079577	0.0150000	1.0009	666.5868	0.3309	241.4999	67.84599297
0.084589	0.0160000	1.0008	669.8525	0.3411	246.7570	68.12903373
0.0505200	0.0170000	1.0142	674.0833	0.3513	250.1311	68.46020091
0.055612	0.0180000	1.0214	678.2529	0.3615	257.5214	68.93936919
0.100524	0.0190000	1.0288	683.8158	0.3117	260.9114	69.26561288
0.106236	0.0200000	1.0362	688.6549	0.3819	260.2820	69.57387671
0.111547	0.0210000	1.0331	693.2009	0.3921	273.6134	70.25562661
0.116159	0.0220000	1.0453	657.4246	0.4021	278.8668	70.81567077
0.122171	0.0230000	1.0547	701.0285	0.4119	288.0851	71.41773212
0.127183	0.0240000	1.0553	704.0591	0.4216	289.1936	72.05925063
0.132755	0.0250000	1.0629	706.4368	0.4311	292.71886019	72.71886019
0.138166	0.0260000	1.0656	708.2283	0.4405	299.0578	72.15508362
0.142418	0.0270000	1.0674	709.4702	0.4495	303.0785	74.20834221
0.148170	0.0280000	1.0695	710.1831	0.4583	308.5357	74.99125905
0.154542	0.0290000	1.0689	710.4359	0.4669	313.0805	75.80850611
0.159354	0.0300000	1.0687	710.2504	0.4753	317.5021	76.6562648
0.164665	0.0310000	1.0680	709.8251	0.4835	321.6075	77.53502519
0.169477	0.0320000	1.0669	709.0568	0.4915	326.0012	78.44155978
0.175285	0.0330000	1.0655	708.1740	0.4992	330.0287	79.53765254
0.180651	0.0340000	1.0639	707.1171	0.5068	336.0766	80.33755615
0.185512	0.0350000	1.0632	705.5865	0.5142	331.9716	81.32407072
0.191224	0.0360000	1.0603	703.8139	0.5221	341.7811	82.3357552
0.195526	0.0370000	1.0537	703.6577	0.5295	345.5123	83.37009368
0.201348	0.0380000	1.0570	702.5469	0.5354	347.1724	84.42721940
0.207150	0.0390000	1.0555	701.5091	0.5423	352.7683	85.50585803
0.212471	0.0400000	1.0540	700.5657	0.5490	356.3066	86.50793584
0.217783	0.0410000	1.0528	699.7317	0.5566	359.7932	87.72379555
0.223055	0.0420000	1.0517	698.4215	0.5621	363.2337	88.86119521
0.228447	0.0430000	1.0504	697.2129	0.5686	366.6330	89.01631050
0.233715	0.0440000	1.0501	697.556	0.5719	369.955	90.12823781
0.239330	0.0450000	1.0496	697.4088	0.5813	373.3248	92.37806655
0.244342	0.0460000	1.0492	697.3790	0.5915	376.6242	93.57930328
0.249455	0.0470000	1.0490	697.2465	0.5937	379.8963	94.79621372
0.254566	0.0480000	1.0495	697.2626	0.5959	393.1430	95.02665034
0.265585	0.0500000	1.0493	697.3033	0.6121	389.5668	96.52555874
0.270551	0.0510000	1.0495	697.5625	0.6181	392.7558	97.79200505
0.275213	0.0520000	1.0499	697.7876	0.6241	395.9038	101.05915367

FLAT STEEL PLATE EXPOSED TO CONST HEAT FLUX

ECRE RADIUS (FT.) = 0.05089  
 CUTER RADIUS (FT.) = 0.05194  
 ECRE TO THERMOCOUPLE DISTANCE (FT.) = 0.002203  
 INITIAL THERMOCOUPLE TEMPERATURE (F.) = 67.00000  
 INITIAL GAS TEMPERATURE (F.) = 4.6167  
 THERMAL DIFFUSIVITY (FTSQ/SEC) = 0.00010307  
 THERMAL CONDUCTIVITY (BTU/FT·SEC·F.) = 0.005555555  
 NUMBER OF TIME TEMPERATURE PAIRS (SEC., F.) = 16  
 NUMBER OF E(11) CCOFF. TO BE FITTED = 16

TIME	TEMPERATURE
0.01382603648	99.8675803200
0.0376521296	152.8049415000
0.0564781544	201.5980147000
0.0753042592	245.9095645000
0.0941303241	286.5662660000
0.1129563889	324.2767059000
0.1317824537	359.5725166000
0.1506085185	392.8529152000
0.1694345833	424.4155064000
0.1882606481	454.5057479000
0.2070867129	483.3696407000
0.2259127777	510.9708251000
0.2447388425	537.6174433000
0.2635649073	563.3527744000
0.2823509722	588.2638204000
0.3012170370	612.4246448000
TIME OF DATA SHIFTED BY (SEC)	0.0
COEFFICIENTS OF B(11) = 0.2604458200	1 = 1
COEFFICIENTS OF B(11) = -0.9959523250	1 = 2
COEFFICIENTS OF B(11) = 0.2646403590	1 = 3
COEFFICIENTS OF B(11) = -0.1560871200	1 = 4
COEFFICIENTS OF B(11) = 0.5384521700	1 = 5
COEFFICIENTS OF B(11) = -0.4543489400	1 = 6
COEFFICIENTS OF B(11) = 0.2819266900	1 = 7
COEFFICIENTS OF B(11) = -0.1305435100	1 = 8
COEFFICIENTS OF B(11) = 0.4594062900	1 = 9
COEFFICIENTS OF B(11) = -0.1219223200	1 = 10
COEFFICIENTS OF B(11) = 0.2431336100	1 = 11
COEFFICIENTS OF B(11) = -0.358315670-01	1 = 12
COEFFICIENTS OF B(11) = 0.378302190-02	1 = 13
COEFFICIENTS OF B(11) = -0.27048370-03	1 = 14
COEFFICIENTS OF B(11) = 0.117252270-04	1 = 15
COEFFICIENTS OF B(11) = -0.232574520-06	1 = 16

FLAT STEEL PLATE EXPOSED TO CONST. HEAT FLUX  
TEMPERATURE AND HEAT FLUX  
PREDICTED BY JENG-SHING CHIOU

(NORM.) (SEC.)	(REAL TH.) (SEC.)	(NORM. FLUX) (BTU/FT <sup>2</sup> SEC)	(REAL FLUX) (BTU/FT <sup>2</sup> SEC)	(NORM.) (DEG.F.)	(REAL, S.U.F.T.) (DEG.F.)	(TH. CO.T.) (DEG.F.)
0.021247	0.0010000	0.3850	511.7559	0.0511	93.5223	67.000000252
0.042474	0.0020000	0.4909	652.504	0.0946	116.8632	67.00021528
0.063741	0.0030000	0.5442	723.3625	0.1317	136.4065	67.06951214
0.084565	0.0040000	0.5110	759.0234	0.1635	153.0213051	67.32513051
0.106226	0.0050000	0.5227	774.5318	0.1909	167.5889	67.884251188
0.127483	0.0060000	0.5853	777.556	0.2146	180.0826	68.80119885
0.148730	0.0070000	0.5923	773.306	0.2353	190.9866	70.080333053
0.169577	0.0080000	0.5767	766.5576	0.2535	209.5690	71.69244833
0.191224	0.0090000	0.5693	756.8054	0.2697	209.1318	72.60518047
0.212411	0.0100000	0.5615	756.3735	0.2813	216.8173	75.76313149
0.233175	0.0110000	0.2537	735.0173	0.2976	223.8137	78.12364280
0.254566	0.0120000	0.5465	726.4203	0.3098	230.2596	80.64387126
0.276213	0.0130000	0.5010	717.8032	0.3212	250.16188	83.25016188
0.297460	0.0140000	0.5346	710.3617	0.3319	241.9333	86.02362004
0.318767	0.0150000	0.5267	704.0537	0.3422	253.3279	88.83344223
0.339554	0.0160000	0.5258	695.0051	0.3520	257.5117	91.68592736
0.361261	0.0170000	0.5229	695.0644	0.3615	257.5313	94.56750327
0.382448	0.0180000	0.5206	692.0911	0.3708	262.4227	97.46107977
0.403676	0.0190000	0.5191	690.0337	0.3799	267.2136	100.3714738
0.424853	0.0200000	0.5180	668.7504	0.3799	271.9242	102.28071853
0.446010	0.0210000	0.5177	688.1193	0.3977	276.5694	106.18283212
0.467437	0.0220000	0.5176	688.0237	0.4064	281.1559	105.07652357
0.488368	0.0230000	0.5188	683.3546	0.4150	287.7009	111.95461715
0.509521	0.0240000	0.5183	685.0125	0.4235	290.1577	114.83024515
0.531178	0.0250000	0.5190	699.9103	0.4320	294.6519	117.69315310
0.552226	0.0260000	0.5188	668.5619	0.4403	120.53254501	
0.573473	0.0270000	0.5207	692.1257	0.4486	303.4325	123.367014
0.594520	0.0280000	0.5216	653.3224	0.4568	301.7570	126.18369836
0.616167	0.0290000	0.5225	594.5150	0.4650	312.0354	128.99318448
0.637414	0.0300000	0.5233	695.6615	0.4730	316.2659	131.78547381
0.658461	0.0310000	0.5222	656.7522	0.4809	320.4467	134.56806272
0.679508	0.0320000	0.5249	657.7489	0.4988	324.5761	137.35933401
0.701152	0.0330000	0.5217	701.1527	0.5335	348.2138	153.7562473
0.722103	0.0340000	0.5202	655.4220	0.5041	332.6750	142.35106107
0.743650	0.0350000	0.5267	700.1007	0.5117	336.627	145.55163257
0.764133	0.0360000	0.5231	701.544	0.5546	359.2952	161.80246558
0.786154	0.0370000	0.5274	701.1078	0.5264	346.4123	151.03798356
0.807157	0.0380000	0.5290	701.9225	0.5632	366.4238	167.11855702
0.828368	0.0390000	0.5210	701.8150	0.5748	367.9148	167.75534380
0.849556	0.0400000	0.5279	701.7103	0.5407	373.3591	172.38202713
0.871268	0.0410000	0.5280	701.6626	0.5477	355.6545	159.12807916
0.892615	0.0420000	0.5221	701.3847	0.5541	382.1149	177.5977988
1.019833	0.0430000	0.5275	701.1559	0.6001	383.3301	186.18904355
1.041110	0.0440000	0.5214	701.0915	0.5067	382.7059	182.76143046
1.062357	0.0500000	0.5212	700.915	0.6128	382.5641	191.32352428
1.083604	0.0510000	0.5271	700.4347	0.6189	393.1465	187.97244521
1.104551	0.0520000	0.5265	700.4554	0.6249	396.3147	190.4071979

## APPENDIX C-2

## THE SLAB WITH PERIODIC SURFACE TEMPERATURE

PLATE RADIUS (FT.) = 0.05089  
 CUTER RADIUS (FT.) = 0.10200  
 DUE TO THERMOCOUPLE DISTANCE (FT.) = 0.0002220  
 INITIAL THERMOCOUPLE TEMPERATURE (F.) = 80.0000  
 INITIAL GAS TEMPERATURE (F.) = 44.615  
 INTERNAL DIFFUSIVITY (FT.SQ/SEC) = 0.00010307  
 THERMAL CONDUCTIVITY(BTU/FT.SEC.F.) = 0.00555555  
 NUMBER OF TIME TEMPERATURE PAIRS (SEC., F.) = 15  
 NUMBER OF E(11) COFF. TO BE FITTED = 15

TIME	TEMPERATURE
0.000700000	217.4304650000
0.001000000	276.4284470000
0.001100000	335.0123870000
0.001300000	390.6554760000
0.001500000	441.2844230000
0.001700000	495.1659150000
0.001900000	520.8663420000
0.002100000	547.2406900000
0.002300000	563.4333160000
0.002500000	568.8819130000
0.002700000	583.3198170000
0.002900000	546.7748510000
0.003100000	519.6368800000
0.003300000	482.2791520000
0.003500000	435.7737270000
TIME OF DATA SHIFTED BY (SEC)	0.0
COEFFICIENTS OF B(11) = 3.515867190 00	1 = 1
COEFFICIENTS OF B(11) = -0.804359240-01	1 = 2
COEFFICIENTS OF B(11) = 0.108357980 00	1 = 3
COEFFICIENTS OF B(11) = -0.116257760 00	1 = 4
COEFFICIENTS OF B(11) = 0.792457150-01	1 = 5
COEFFICIENTS OF B(11) = -0.322792890-01	1 = 6
COEFFICIENTS OF B(11) = 0.145442150-01	1 = 7
COEFFICIENTS OF B(11) = -0.406951500-02	1 = 8
COEFFICIENTS OF B(11) = 0.862021320-03	1 = 9
COEFFICIENTS OF B(11) = -0.137916530-03	1 = 10
COEFFICIENTS OF B(11) = 0.163344430-04	1 = 11
COEFFICIENTS OF B(11) = -0.139168190-05	1 = 12
COEFFICIENTS OF B(11) = 0.805611200-07	1 = 13
COEFFICIENTS OF B(11) = -0.233646920-08	1 = 14
COEFFICIENTS OF R(11) = 0.458429940-10	1 = 15

THE SLAB WITH PERIODIC SURFACE TEMPERATURE  
TEMPERATURE AND HEAT FLUX  
PREDICTED BY HSAI-YIN LEE

(NUR.1) (DLESS)	(REALTM) (SEC.)	(INOR.FLUX) (DLESS)	(REAL.FLUX) (BTU/FTSQ SEC)	(NUR.) (DLESS)	(REAL.SUF.T) (DEG.F.)	(TH.CO.1) (CEG.F.)
2.212375	C.CC01000	0.2596	3534.6035	0.1069	137.6490	82.44126335
2.12415C	C.0002000	0.1600	4702.0712	0.2100	193.4660	93.33060384
2.637125	D.0003000	0.4343	5914.7548	0.3111	247.5886	112.9245213
2.84970C	C.CC4C00	0.4942	6729.3413	0.4101	301.4311	135.71658123
1.Cal175	C.CC05000	0.5433	7358.4433	0.5065	353.5044	161.2965414
1.274250	C.0006000	0.5835	7946.5473	0.5598	403.8978	188.7553437
1.465625	D.0007000	0.6159	8387.6363	0.6194	452.2995	217.4304550
1.69900C	D.0008000	0.6411	8730.7338	0.7148	450.3588	246.7555353
1.711375	C.CC9C00	0.6594	8982.1487	0.8554	541.9112	276.428477C
1.123149	D.0010000	0.6717	9144.3874	0.9307	582.5683	295.9423372
2.335124	D.0011000	0.6770	9227.2243	1.1032	620.1231	335.512377C
2.544855	D.0012000	0.6716	9227.7377	1.3036	654.3358	363.33093137
2.76314	D.0013000	0.6720	9156.7958	1.1234	685.0052	390.6555476C
2.573149	D.0014000	0.6608	8399.2010	1.1703	711.9396	415.71231113
2.185424	D.0015000	0.6644	8715.5C82	1.2129	734.9734	441.284433C0
2.357559	D.0016000	0.6620	8483.6003	1.2481	754.9650	464.16459851
2.410314	C.CC1C00	0.5967	8125.7368	1.2770	768.7977	485.165919C3
2.422149	D.0018000	0.5658	7765.5645	1.2931	779.3802	50.2177491
4.035124	D.0019000	0.5307	7226.7369	1.3068	795.6476	520.966342C0
4.242459	C.GC2C00	0.4915	6693.017C	1.3133	787.5615	535.27955190
4.459174	D.0021000	0.4986	6108.0798	1.3053	798.8443	547.24C05900
4.672245	C.GC22C00	0.4922	5677.6032	1.2932	776.3093	556.0517639
4.384624	D.0023000	0.3528	4804.2763	1.2726	757.2035	563.433315C1
5.055559	D.0024000	0.3906	4093.7832	1.2542	751.8527	575.52422565
5.3C3174	D.0025000	0.2661	3350.7634	1.2081	732.3635	508.831733C2
5.521745	D.0026000	0.1895	2580.25C7	1.1645	7C8.8443	567.43231449
5.734124	D.0027000	0.1313	1787.4441	1.1138	661.4492	563.319917C5
5.740458	D.0028000	0.0118	9177.4862	1.1562	656.4471570	556.4471570
6.152813	D.0029000	0.0114	155.7348	0.9921	615.7210	546.7749113
6.371268	D.0030000	-0.0996	-672.4449	0.9218	577.7933	534.47135521
6.583623	D.0031000	-0.1103	-1501.6767	0.8859	519.56358330	
6.752552	D.0032000	-0.1709	-2326.6171	0.7658	492.5790	502.13315545
7.0C4313	D.0033000	-0.2307	-3141.6700	0.6759	446.6147	492.2791263
7.22C148	D.0034000	-0.2293	-3542.2226	0.5583	357.4893	46.1162292
7.433123	D.0035000	-0.1667	-4721.4837	0.4952	347.47173	435.7737346
7.645458	D.0036000	-0.4417	-5470.221C	0.3987	295.2161	405.3595305
7.357173	D.0037000	-0.4532	-5170.5C74	0.3022	242.1060	381.16525C2
8.070248	D.0038000	-0.4933	-6785.2335	0.2016	188.8322	351.31552168
8.282433	D.0039000	-0.508	-7228.5655	0.1063	137.3508	220.2324324
8.494558	D.0040000	-0.5378	-7323.11C4	0.0215	91.5981	288.469526C7
			TEMPERATURE AND HEAT FLUX PREDICTED BY HSAI-YIN LEE			

THE SLAB WITH PERIODIC SURFACE TEMPERATURE

(NUR.1)  
(DLESS)

(REALTM)  
(SEC.)

(INOR.FLUX)  
(DLESS)

(REAL.FLUX)  
(BTU/FTSQ SEC)

(NUR.)  
(DLESS)

(REAL.SUF.T)  
(DEG.F.)

(TH.CO.1)  
(CEG.F.)

## APPENDIX C - 3

M60 GUN THERMOCOUPLE 10 21 INCHES FROM BREECH  
TEMPERATURE AND HEAT FLUX  
PREDICTED BY HSIAI-YIN LEE

(NORM.TM) (DMLESS)	(REAL.TM) (SEC.)	(NOR. FLUX) (DMLESS)	(REAL. FLUX) (BTU/FTSQ SEC)	(NOR. ) (DMLESS)	(REAL.SUF.T) (DEG.F.)	(TH.CO.T) (DEG.F.)
0.036557	0.0010000	0.8394	1504.5883	0.1524	160.8969	78.80248541
0.073914	0.0020000	0.9208	1650.3766	0.2540	215.6479	79.10233268
0.110872	0.0030000	0.8640	1548.6441	0.3169	249.5496	80.75107503
0.147829	0.0040000	0.7536	1350.6857	0.3509	267.8668	84.24427704
0.184786	0.0050000	0.6257	1121.5054	0.3638	274.8224	89.27332529
0.221743	0.0060000	0.4987	893.8086	0.3618	273.7626	95.23538854
0.258701	0.0070000	0.3818	684.3901	0.3499	267.3031	101.53891376
0.295658	0.0080000	0.2796	501.2017	0.3316	257.4543	107.70954560
0.332615	0.0090000	0.1936	347.0147	0.3098	245.7306	113.41640288
0.367572	0.0100000	0.1236	221.5337	0.2866	233.2437	118.45552591
0.406529	0.0110000	0.0685	122.6979	0.2635	220.7827	122.72522448
0.443487	0.0120000	0.0265	47.5114	0.2414	208.8817	126.20000000
0.480444	0.0130000	-0.0041	-7.4187	0.2210	197.8764	128.90746934
0.517401	0.0140000	-0.0254	-45.5354	0.2026	187.9522	130.90947620
0.554358	0.0150000	-0.0391	-70.1131	0.1863	179.1825	132.28740140
0.591316	0.0160000	-0.0469	-84.1269	0.1722	171.5612	133.13121584
0.628273	0.0170000	-0.0503	-90.1846	0.1600	165.0278	133.53167289
0.665230	0.0180000	-0.0505	-90.4995	0.1498	159.4880	133.575033C9
0.702187	0.0190000	-0.0485	-86.8935	0.1411	154.8290	133.33976926
0.739144	0.0200000	-0.0451	-80.8193	0.1339	150.9312	132.89477620
0.776102	0.0210000	-0.0409	-73.3944	0.1278	147.6772	132.29869263
0.813059	0.0220000	-0.0365	-65.4423	0.1228	144.9571	131.60000000
0.850016	0.0230000	-0.0321	-57.5369	0.1185	142.6723	130.83766159
0.886973	0.0240000	-0.0279	-50.0464	0.1150	140.7379	130.04207940
0.923931	0.0250000	-0.0241	-43.1766	0.1119	139.0826	129.23622713
0.960088	0.0260000	-0.0206	-37.0095	0.1092	137.6490	128.43684019
0.997045	0.0270000	-0.0176	-31.5388	0.1069	136.3923	127.65558013
1.034802	0.0280000	-0.0149	-26.7007	0.1048	135.2788	126.90011703
1.071759	0.0290000	-0.0125	-22.3991	0.1030	134.2837	126.17509340
1.108717	0.0300000	-0.0103	-18.5266	0.1013	133.3900	125.48295008
1.145674	0.0310000	-0.0084	-14.9803	0.0998	132.5860	124.8240683
1.182631	0.0320000	-0.0065	-11.6736	0.0985	131.8634	124.20000000
1.219588	0.0330000	-0.0048	-8.5437	0.0973	131.2163	123.6C848629
1.256546	0.0340000	-0.0031	-5.5551	0.0962	130.6394	123.04912647
1.293503	0.0350000	-0.0015	-2.7011	0.0953	130.1271	122.52086579
1.330460	0.0360000	-0.0000	-0.0017	0.0944	129.6728	122.02262934
1.367417	0.0370000	0.0014	2.5000	0.0937	129.2682	121.55335075
1.404374	0.0380000	0.0026	4.7427	0.0930	128.9033	121.11195267
1.441322	0.0390000	0.0037	6.6538	0.0924	128.5662	120.69725537
1.478289	0.0400000	0.0045	8.1551	0.0918	128.2437	120.30810907
1.515246	0.0410000	0.0051	9.1704	0.0912	127.9213	119.94292268
1.552203	0.0420000	0.0054	9.6319	0.0905	127.5841	119.60000000
1.589161	0.0430000	0.0053	9.4859	0.0899	127.2172	119.27729167
1.626118	0.0440000	0.0049	8.6978	0.0891	126.8066	118.972409C1
1.663075	0.0450000	0.0040	7.2564	0.0882	126.3397	118.68262348
1.700032	0.0460000	0.0029	5.1763	0.0872	125.8062	118.40489326
1.736589	0.0470000	0.0014	2.4992	0.0861	125.1984	118.13591670
1.773547	0.0480000	-0.0004	-0.7061	0.0848	124.5119	117.87221030
1.810504	0.0490000	-0.0024	-4.3452	0.0834	123.7461	117.61020774
1.847661	0.0500000	-0.0046	-8.3012	0.0819	122.9038	117.34637497
1.884818	0.0510000	-0.0069	-12.4387	0.0802	121.9919	117.07733556
1.921776	0.0520000	-0.0093	-16.6090	0.0784	121.0208	116.80000000

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1.958733	0.0530000	-0.0115	-20.6564	0.0765	120.0045	116.51169214
1.995690	0.0540000	-0.0136	-24.4246	0.0745	118.9596	116.2102603
2.032647	0.0550000	-0.0155	-27.7644	0.0726	117.9049	115.89420669
2.069605	0.0560000	-0.0170	-30.5408	0.0706	116.8606	115.56270890
2.106562	0.0570000	-0.0182	-32.6399	0.0688	115.8475	115.21572882
2.143519	0.0580000	-0.0190	-33.9759	0.0670	114.8858	114.85400433
2.180476	0.0590000	-0.0192	-34.4966	0.0653	113.9941	114.47904109
2.217433	0.0600000	-0.0191	-34.1882	0.0638	113.1885	114.09306290
2.254391	0.0610000	-0.0185	-33.0787	0.0625	112.4817	113.69892615
2.291348	0.0620000	-0.0174	-31.2390	0.0614	111.8818	113.30000000
2.328305	0.0630000	-0.0161	-28.7829	0.0605	111.3918	112.90001541
2.365262	0.0640000	-0.0144	-25.8643	0.0598	111.0094	112.50288784
2.402220	0.0650000	-0.0126	-22.6726	0.0593	110.7258	112.11252006
2.439177	0.0660000	-0.0108	-19.4252	0.0589	110.5267	111.73259289
2.476134	0.0670000	-0.0091	-16.3580	0.0586	110.3920	111.36635320
2.513051	0.0680000	-0.0077	-13.7137	0.0585	110.2968	111.01640954
2.551048	0.0690000	-0.0065	-11.7275	0.0583	110.2122	110.68454671
2.587006	0.0700000	-0.0059	-10.6115	0.0581	110.1071	110.37157113
2.623563	0.0710000	-0.0059	-10.5372	0.0578	109.9494	110.07719914
2.660920	0.0720000	-0.0065	-11.6178	0.0574	109.7091	109.80000000
2.697077	0.0730000	-0.0077	-13.8892	0.0567	109.3599	109.53740473
2.734835	0.0740000	-0.0096	-17.2932	0.0558	108.8823	109.28579062
2.771792	0.0750000	-0.0121	-21.6612	0.0547	108.2665	109.04064932
2.808149	0.0760000	-0.0149	-26.7011	0.0533	107.5151	108.79684356
2.845706	0.0770000	-0.0178	-31.9890	0.0517	106.6461	108.54895463
2.882663	0.0780000	-0.0206	-36.9667	0.0499	105.6949	108.29171788
2.919621	0.0790000	-0.0223	-40.9470	0.0481	104.7166	108.02053902
2.956578	0.0800000	-0.0241	-43.1290	0.0464	103.7867	107.73207721
2.993535	0.0810000	-0.0238	-42.6264	0.0449	103.0010	107.42487414
3.030492	0.0820000	-0.0215	-38.5100	0.0439	102.4737	107.10000000
3.067450	0.0830000	-0.0167	-29.8692	0.0437	102.3342	106.76167727
3.104407	0.0840000	-0.0089	-15.8947	0.0444	102.7200	106.41733289
3.141364	0.0850000	0.0022	4.0146	0.0463	103.7685	106.08051657
3.178321	0.0860000	0.0168	30.1143	0.0497	105.6030	105.76611245
3.215278	0.0870000	0.0347	62.1450	0.0548	108.3163	105.49524511
3.252236	0.0880000	0.0553	99.1241	0.0615	111.9472	105.29228652
3.289193	0.0890000	0.0776	139.0949	0.0699	116.4528	105.18432393
3.326150	0.0900000	0.0998	178.8266	0.0796	121.6722	105.19945055
3.363107	0.0910000	0.1191	213.4559	0.0900	127.2828	105.36421092
3.400065	0.0920000	0.1317	236.0913	0.1001	132.7477	105.70000001
3.437022	0.0930000	0.1324	237.2856	0.1085	137.2510	106.21821278
3.473579	0.0940000	0.1141	204.5369	0.1129	139.6254	106.91388248
3.510536	0.0950000	0.0678	121.5264	0.1104	138.2573	107.75753978
3.547053	0.0960000	-0.0180	-32.2719	0.0969	130.9886	108.68498008
3.584051	0.0970000	-0.1579	-283.0449	0.0672	114.9924	109.58459482
3.621808	0.0980000	-0.3699	-662.9825	0.0145	86.6343	110.28138452
3.658765	0.0990000	-0.6760	-1211.6418	-0.0696	41.3115	110.52073128
3.695722	0.1000000	-1.1030	-1977.0840	-0.1959	-26.7312	109.94096599

TEMPERATURE AND HEAT FLUX  
PREDICTED BY HSAI-YIN LEE

(NORM.TM) (REAL.TM) (NUR.FLUX) (REAL.FLUX) (NOR. ) (REAL.SUF.T) (TH.CO.T)  
(DMLESS) (SEC.) (DMLESS) (BTU/FTSQ SEC) (DMLESS) (DEG.F.) (CFG.F.)

M60 GUN THERMOCOUPLE 10 21 INCHES FROM BREACH

BEST AVAILABLE COPY

M60 GUN THERMOCOUPLE 7 15.0 INCHES FROM BREECH  
BORE SURFACE TEMPERATURE AND HEAT FLUX PROGRAM  
NUMBER OF B(I) COFF. TO BE FITTED = 10

TIME TEMPERATURE

0.0100000000	110.6000000000
0.0200000000	124.5000000000
0.0300000000	123.1000000000
0.0400000000	119.6000000000
0.0500000000	116.9000000000
0.0600000000	114.1000000000
0.0700000000	112.5000000000
0.0800000000	110.3000000000
0.0900000000	107.6000000000
0.1000000000	106.8000000000

TIME OF DATA SHIFTED BY (SEC) -0.002000

BORE RADIUS (FT.) = 0.01250

CUTER RADIUS (FT.) = 0.04417

BORE TO THERMOCOUPLE DISTANCE (FT) = 0.001830

INITIAL THERMOCOUPLE TEMPERATURE (F.) = 78.8000

INITIAL GAS TEMPERATURE (F.) = 4.4937

TERMAL DIFFUSIVITY (FTSQ/SEC) = 0.00010307

TERMAL CONDUCTIVITY(BTU/FT.SEC.F.) = 0.00555555

NUMBER OF TIME TEMPERATURE PAIRS (SEC.,F.)=10

NUMBER OF B(I) COFF. TO BE FITTED = 10

TIME TEMPERATURE

0.0120000000	110.6000000000
0.0220000000	124.5000000000
0.0320000000	123.1000000000
0.0420000000	119.6000000000
0.0520000000	116.9000000000
0.0620000000	114.1000000000
0.0720000000	112.5000000000
0.0820000000	110.3000000000
0.0920000000	107.6000000000
0.1020000000	106.8000000000

TIME OF DATA SHIFTED BY (SEC) -0.002000

COEFFICIENTS OF B(I) = 0.42998293D 01 I = 1

COEFFICIENTS OF B(I) = -0.237957C8D 02 I = 2

COEFFICIENTS OF B(I) = 0.64874208D 02 I = 3

COEFFICIENTS OF B(I) = -0.10652420D 03 I = 4

COEFFICIENTS OF B(I) = 0.11235543D 03 I = 5

COEFFICIENTS OF B(I) = -0.77590114D 02 I = 6

COEFFICIENTS OF B(I) = 0.34776113D 02 I = 7

COEFFICIENTS OF B(I) = -0.97172330D 01 I = 8

COEFFICIENTS OF B(I) = 0.15343C23D 01 I = 9

COEFFICIENTS OF B(I) = -0.10434213D 00 I = 10

BEST AVAILABLE COPY

M60 GUN THERMOCOUPLE 7 15.0 INCHES FROM BREECH  
TEMPERATURE AND HEAT FLUX  
PREDICTED BY HSAI-YIN LEE

(NORTM) (DMLESS)	(REALTM) (SEC.)	(NOR. FLUX) (DMLESS)	(REAL. FLUX) (BTU/FTSQ SEC)	(NOR. ) (DMLESS)	(REAL.SUF.T) (DEG.F.)	(TH.CO.T) (DEG.F.)
C.030777	0.0010000	0.6763	1106.1522	0.1116	138.9277	78.80037178
C.061555	C.0200000	0.7553	1235.3923	0.1882	180.2072	78.89254605
C.092332	0.0030000	0.7260	1187.5234	0.2382	207.1313	79.56894741
C.1231C9	0.0040000	0.6538	1069.4664	0.2682	223.3065	81.26193308
C.153E86	0.0050000	0.5665	926.6396	0.2836	231.5973	83.98321767
C.184664	0.0060000	0.4779	781.6665	0.2885	234.2521	87.49324244
C.215441	C.0070000	0.3950	646.0729	0.2862	233.0109	91.48181303
C.246218	0.0080000	0.3211	525.3048	0.2791	229.1991	95.66106957
C.276555	0.0090000	0.2576	421.3085	0.2691	223.8069	99.80078458
C.307773	C.0100000	0.2042	334.0054	0.2575	217.5565	103.73480391
C.338550	0.0110000	0.1603	262.1869	0.2453	210.9596	107.35437364
C.369227	C.0120000	0.1248	204.0755	0.2330	204.3639	110.60030000
C.400105	0.0130000	0.0964	157.6791	0.2212	197.9933	113.44515669
C.430E82	0.0140000	0.0740	121.0127	C.2101	191.9790	115.89130675
C.461E59	0.0150000	0.0564	92.2309	0.1997	186.3861	117.95651340
C.492436	C.0160000	0.0426	69.7006	0.1901	181.2338	119.66963617
C.523214	0.0170000	0.0318	52.0336	0.1913	176.5111	121.06507221
C.553991	0.0180000	0.0233	38.0914	0.1733	172.1898	122.17914841
C.584768	0.0190000	0.0165	26.9727	0.1660	168.2326	123.04762728
C.615545	0.0200000	0.0110	17.9899	0.1592	164.6001	123.70412166
C.646773	0.0210000	0.0065	10.6397	0.1530	161.2549	124.17920042
C.677100	0.0220000	0.0028	4.5708	0.1473	158.1641	124.50000000
C.707E77	0.0230000	-0.0003	-0.4481	0.1420	155.3006	124.69018933
C.738E54	0.0240000	-0.0028	-4.5583	0.1371	152.6436	124.77016617
C.769432	0.0250000	-0.0048	-7.8388	0.1325	150.1780	124.75738986
C.800209	0.0260000	-0.0063	-10.3301	0.1292	147.8935	124.66677881
C.830586	0.0270000	-0.0074	-12.0536	0.1243	145.7836	124.51112060
C.861764	0.0280000	-0.0080	-13.0274	0.1207	143.8440	124.30145876
C.892541	0.0290000	-0.0081	-13.2779	0.1174	142.0720	124.04743335
C.923318	C.0300000	-0.0079	-12.8478	0.1144	140.4646	123.75756275
C.954095	0.0310000	-0.0072	-11.8000	0.1118	139.0180	123.43946213
C.984E73	0.0320000	-0.0062	-10.2193	0.1094	137.7265	123.10000000
I.015E50	0.0330000	-0.0050	-8.2111	0.1072	136.5822	122.74535846
I.046427	0.0340000	-0.0036	-5.8975	0.1054	135.5745	122.38128553
I.0772C4	0.0350000	-0.0021	-3.4145	0.1037	134.6903	122.01270981
I.107582	0.0360000	-0.0006	-0.9020	0.1023	133.9141	121.64412808
I.138759	0.0370000	0.0009	1.4986	0.1010	133.2281	121.27937668
I.169536	0.0380000	0.0022	3.6537	0.0999	132.6132	120.92163655
I.200314	C.0390000	0.0033	5.4431	0.0988	132.0494	120.57340112
I.231C91	0.0400000	0.0041	6.7664	0.0978	131.5166	120.23645429
I.261E68	0.0410000	0.0046	7.5481	0.0969	130.9956	119.91186465
I.292E45	0.0420000	0.0047	7.7416	0.0959	130.4685	119.60000000
I.323423	0.0430000	0.0045	7.3318	0.0949	129.9200	119.30056472
I.354200	0.0440000	0.0039	6.3368	0.0938	129.3377	119.01266071
I.384577	C.0450000	0.0029	4.8074	0.0926	128.7128	118.73487111
I.415754	0.0460000	0.0017	2.8256	0.0914	128.0407	118.46536456
I.446532	0.0470000	0.0003	0.5012	0.0901	127.3208	118.20201648
I.4773C9	0.0480000	-0.0012	-2.0324	0.0886	126.5572	117.94254276
I.503C86	0.0490000	-0.0028	-4.6242	0.0872	125.7579	117.68464069
I.538E64	C.0500000	-0.0043	-7.1125	0.0856	124.9352	117.42613058
I.567E41	C.0510000	-0.0057	-9.3323	0.0841	124.1045	117.16505487
I.600418	0.0520000	-0.0068	-11.1238	0.0826	123.2840	116.90000000

BEST AVAILABLE COPY

1.631195	0.0530000	-0.0075	-12.3407	0.0811	122.4936	116.62980900
1.661973	0.0540000	-0.0079	-12.8587	0.0797	121.7543	116.35406521
1.692750	0.0550000	-0.0077	-12.5835	0.0785	121.0865	116.07295073
1.723527	0.0560000	-0.0070	-11.4581	0.0774	120.5091	115.78731264
1.754304	0.0570000	-0.0058	-9.4691	0.0765	120.0384	115.49855460
1.785082	0.0580000	-0.0041	-6.6513	0.0759	119.6867	115.20909215
1.815859	0.0590000	-0.0019	-3.0911	0.0755	119.4612	114.92127132
1.846636	0.0600000	0.0007	1.0727	0.0753	119.3632	114.63825189
1.877413	0.0610000	0.0035	5.6503	0.0753	119.3871	114.36335794
1.908151	0.0620000	0.0064	10.4044	0.0756	119.5200	114.10000000
1.938568	0.0630000	0.0092	15.0566	0.0760	119.7415	113.85147476
1.969745	0.0640000	0.0118	19.2967	0.0765	120.0235	113.62074965
2.000523	0.0650000	0.0139	22.7950	0.0771	120.3310	113.41024106
2.031300	0.0660000	0.0154	25.2162	0.0776	120.6229	113.22159622
2.062077	0.0670000	0.0160	26.2372	0.0780	120.8533	113.05548985
2.092854	0.0680000	0.0156	25.5661	0.0783	120.9735	112.91144729
2.123632	0.0690000	0.0140	22.9636	0.0782	120.9340	112.73770657
2.154409	0.0700000	0.0112	18.2652	0.0777	120.6875	112.68113169
2.185186	0.0710000	0.0070	11.4040	0.0768	120.1913	112.58718931
2.215963	0.0720000	0.0015	2.4333	0.0754	119.4133	112.50000000
2.246741	0.0730000	-0.0052	-8.4524	0.0734	118.3307	112.41247419
2.277518	0.0740000	-0.0128	-20.8992	0.0708	116.9386	112.31654053
2.308295	0.0750000	-0.0210	-34.3797	0.0677	115.2512	112.20347218
2.339073	0.0760000	-0.0295	-48.1849	0.0640	113.3062	112.36431290
2.369850	0.0770000	-0.0376	-61.4238	0.0601	111.1669	111.39040074
2.400127	0.0780000	-0.0446	-73.0337	0.0559	108.9251	111.67398156
2.431404	0.0790000	-0.0500	-81.8020	0.0518	106.7017	111.40800207
2.462182	0.0800000	-0.0528	-86.4042	0.0480	104.6464	111.09135344
2.492959	0.0810000	-0.0522	-85.4602	0.0448	102.9352	110.72065125
2.523736	0.0820000	-0.0474	-77.6121	0.0426	101.7655	110.30000000
2.554513	0.0830000	-0.0377	-61.6278	0.0418	101.3482	109.83720247
2.585291	0.0840000	-0.0223	-36.5344	0.0429	101.8957	109.34525107
2.616668	0.0850000	-0.0011	-1.7851	0.0460	103.6061	108.64272839
2.646645	0.0860000	0.0260	42.5338	0.0517	106.6418	108.35392925
2.677623	0.0870000	0.0584	95.4527	0.0600	111.1017	107.90859956
2.708400	0.0880000	0.0946	154.8176	0.0709	116.9867	107.34117322
2.739177	0.0890000	0.1326	216.9337	0.0842	124.1565	107.28935705
2.769554	0.0900000	0.1683	276.1452	0.0993	132.2773	107.19191523
2.801732	0.0910000	0.1983	324.3462	0.1150	140.7586	107.28545342
2.831509	0.0920000	0.2142	350.4135	0.1297	148.6775	107.60000000
2.862286	0.0930000	0.2076	339.5534	0.1408	154.6894	108.15314365
2.893063	0.0940000	0.1666	272.5553	0.1450	156.9235	108.94246036
2.923841	0.0950000	0.0764	124.9405	0.1375	152.8606	109.93593128
2.954618	0.0960000	-0.0819	-134.0018	0.1121	139.1916	111.06001889
2.985355	0.0970000	-0.3315	-542.3027	0.0610	111.6550	112.18503324
3.016172	0.0980000	-0.7C09	-1146.5363	-0.0259	64.8502	113.10738105
3.046550	0.0990000	-1.2247	-2003.1892	-0.1611	-7.9752	113.52824958
3.077727	0.1000000	-1.9442	-3100.1900	-0.3600	-115.1663	113.02823602

TEMPERATURE AND HEAT FLUX  
PREDICTED BY HSAT-YIN LEE

(NGRTM)	(REALTM)	(NOR.FLUX)	(REAL.FLUX)	(NUR. )	(REAL.SUF.T)	(TH.CO.T)
(DMLFSS)	(SEC.)	(DMLESS)	(BTU/FTSQ SEC)	(DMLESS)	(DEG.F.)	(DEG.F.)

M60 GUN THERMOCOUPLE 7 15.0 INCHES FROM BREECH

BEST AVAILABLE COPY

M60 GUN THERMOCOUPLE 4 9.0 INCHES FROM BREACH  
 BORE SURFACE TEMPERATURE AND HEAT FLUX PROGRAM  
 NUMBER OF B(I) COFF. TO BE FITTED = 10

TIME TEMPERATURE

0.0100000000	143.2000000000
0.0200000000	157.3000000000
0.0300000000	151.0000000000
0.0400000000	144.4000000000
0.0500000000	137.5000000000
0.0600000000	132.8000000000
0.0700000000	129.0000000000
0.0800000000	125.9000000000
0.0900000000	122.8000000000
0.1000000000	119.6000000000

TIME OF DATA SHIFTED BY(SEC) -0.002000

BORE RADIUS (FT.) = 0.01250

CUTTER RADIUS (FT.)= 0.05000

BORE TO THERMOCOUPLE DISTANCE (FT)= 0.001830

INITIAL THERMOCOUPLE TEMPERATURE (F.)= 78.8000

INITIAL GAS TEMPERATURE (F.)= 4.4937

THERMAL DIFFUSIVITY (FTSQ/SEC) = 0.00010307

THERMAL CONDUCTIVITY(BTU/FT.SEC.F.)= 0.00555555

NUMBER OF TIME TEMPERATURE PAIRS (SEC.,F.)=10

NUMBER OF B(I) COFF. TO BE FITTED = 10

TIME TEMPERATURE

0.0120000000	143.2000000000
0.0220000000	157.3000000000
0.0320000000	151.0000000000
0.0420000000	144.4000000000
0.0520000000	137.5000000000
0.0620000000	132.8000000000
0.0720000000	129.0000000000
0.0820000000	125.9000000000
0.0920000000	122.8000000000
0.1020000000	119.6000000000

TIME OF DATA SHIFTED BY(SEC) -0.002000

COEFFICIENTS OF B(I)= 0.97265186D 01 I = 1

COEFFICIENTS OF B(I)= -0.56818786D 02 I = 2

COEFFICIENTS OF B(I)= 0.15078051D 03 I = 3

COEFFICIENTS OF B(I)= -0.22863973D 03 I = 4

COEFFICIENTS OF B(I)= 0.21619564D 03 I = 5

COEFFICIENTS OF B(I)= -0.13191471D 03 I = 6

COEFFICIENTS OF B(I)= 0.52022149D 02 I = 7

COEFFICIENTS OF B(I)= -0.12814655D 02 I = 8

COEFFICIENTS OF B(I)= 0.17933680D 01 I = 9

COEFFICIENTS OF B(I)= -0.10887593D 00 I = 10

BEST AVAILABLE COPY

M60 GUN THERMOCOUPLE 4.9.0 INCHES FROM BREECH  
 TEMPERATURE AND HEAT FLUX  
 PREDICTED BY HSAI-YIN LEE

(NORTM) (DAFLSS)	(REALTM) (SEC.)	(NUR.FLUX) (DMLESS)	(REAL.FLUX) (BTU/FTSQ SEC)	(NOR. ) (DMLESS)	(REAL.SUF.T) (DEG.F.)	(TH.CO.T) (CEG.F.)
C.030777	0.0010000	1.5068	2464.6174	0.2497	213.3551	78.80093965
0.061555	0.0020000	1.6478	2695.2439	0.4155	302.6683	79.00826216
C.092332	0.0030000	1.5392	2517.6782	0.5171	357.4214	80.52178452
C.123109	0.0040000	1.3339	2181.8196	0.5708	386.3264	84.27817903
C.153886	0.0050000	1.0978	1755.6545	0.5895	396.4145	90.24714983
C.184664	0.0060000	0.8645	1414.0284	0.5837	393.2918	97.83164519
C.215441	0.0070000	0.6512	1065.1228	0.5616	381.3640	106.28706625
C.246218	0.0080000	0.4661	762.4626	0.5294	364.0339	114.93619558
C.276555	0.0090000	0.3123	510.7522	0.4920	343.8737	123.24990929
C.307773	0.0100000	0.1892	309.4211	0.4528	322.7744	130.86054314
C.338550	0.0110000	0.0946	154.8016	0.4144	302.0751	137.545+5604
C.369327	0.0120000	0.0254	41.5359	0.3784	282.6736	143.20000000
C.400105	0.0130000	-0.0223	-36.4968	0.3458	265.1212	147.80889574
C.430082	0.0140000	-0.0523	-85.5691	0.3172	249.7018	151.42011603
C.461659	0.0150000	-0.0683	-111.7083	0.2927	236.4994	154.12282411
C.492426	0.0160000	-0.0736	-120.4482	0.2722	225.4530	156.02968789
C.523214	0.0170000	-0.0713	-116.6880	0.2554	216.4023	157.26331105
C.553551	0.0180000	-0.0640	-104.6240	0.2419	209.1240	157.94627789
C.584768	0.0190000	-0.0536	-87.7361	0.2312	203.3609	158.19423324
C.615545	0.0200000	-0.0421	-68.8098	0.2228	198.8445	158.11142795
C.646223	0.0210000	-0.0306	-49.9833	0.2162	195.3114	157.78820971
C.677100	0.0220000	-0.0201	-32.8107	0.2111	192.5163	157.30000000
C.707877	0.0230000	-0.0112	-18.3334	0.2068	190.2392	156.70737747
C.738E54	0.0240000	-0.0044	-7.1563	0.2032	188.2913	156.05693060
C.769432	0.0250000	0.0003	0.4764	0.1999	186.5168	155.38263042
C.800269	0.0260000	0.0028	4.6083	0.1967	184.7935	154.70750441
C.830586	0.0270000	0.0034	5.5020	0.1935	183.0314	154.34549480
C.861764	0.0280000	0.0022	3.5769	0.1900	181.1703	153.40306061
C.892541	0.0290000	-0.0004	-0.6465	0.1863	179.1763	152.78137602
C.923318	0.0300000	-0.0040	-6.5545	0.1823	177.0384	152.17748532
C.954055	0.0310000	-0.0084	-13.6808	0.1781	174.7637	151.58594730
C.984873	0.0320000	-0.0130	-21.3395	0.1737	172.3743	151.00000000
1.015650	0.0330000	-0.0178	-29.0511	0.1691	169.9022	150.41255172
1.046427	0.0340000	-0.0222	-36.3615	0.1644	167.3862	149.81695499
1.077204	0.0350000	-0.0262	-42.8552	0.1597	164.8685	149.20757430
1.107582	0.0360000	-0.0296	-48.3621	0.1551	162.3913	148.58016428
1.138759	0.0370000	-0.0321	-52.5605	0.1507	159.9947	147.93207954
1.169536	0.0380000	-0.0339	-55.3751	0.1465	157.7144	147.26233979
1.200214	0.0390000	-0.0347	-56.7716	0.1425	155.5802	146.57157527
1.231091	0.0400000	-0.0347	-56.7889	0.1389	153.6153	145.86187822
1.261668	0.0410000	-0.0339	-55.5290	0.1356	151.8353	145.13658279
1.292645	0.0420000	-0.0325	-53.1448	0.1326	150.2481	144.40000000
1.323423	0.0430000	-0.0305	-49.8284	0.1300	148.8545	143.65712497
1.354200	0.0440000	-0.0280	-45.7978	0.1278	147.6482	142.91333763
1.384977	0.0450000	-0.0252	-41.2845	0.1259	146.6170	142.17411192
1.415754	0.0460000	-0.0223	-36.5215	0.1242	145.7437	141.44474732
1.446532	0.0470000	-0.0194	-31.7327	0.1229	145.0071	140.73013289
1.477309	0.0480000	-0.0166	-27.1228	0.1217	144.3834	140.03455198
1.508066	0.0490000	-0.0140	-22.8698	0.1207	143.8474	139.36152249
1.538864	0.0500000	-0.0117	-19.1187	0.1198	143.3738	138.71374553
1.569641	0.0510000	-0.0098	-15.9770	0.1190	142.9382	138.09295266
1.600418	0.0520000	-0.0083	-13.5124	0.1183	142.5183	137.50000000

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1.631195	0.0530000	-0.0072	-11.7523	0.1175	142.0945	136.93485581
1.661913	0.0540000	-0.0065	-10.6851	0.1167	141.6511	136.39668628
1.692750	0.0550000	-0.0063	-10.2629	0.1158	141.1763	135.88396353
1.723527	0.0560000	-0.0064	-10.4069	0.1148	140.6627	135.39459839
1.754304	0.0570000	-0.0067	-11.0122	0.1138	140.1072	134.92609049
1.785092	0.0580000	-0.0073	-11.9555	0.1127	139.5111	134.47568749
1.815859	0.0590000	-0.0080	-13.1020	0.1115	138.8794	134.04054543
1.846636	0.0600000	-0.0088	-14.3143	0.1103	138.2204	133.6178259
1.877413	0.0610000	-0.0095	-15.4597	0.1090	137.5451	133.20511968
1.908191	0.0620000	-0.0100	-16.4184	0.1078	136.8660	132.80000000
1.938968	0.0630000	-0.0104	-17.0509	0.1065	136.1964	132.40068430
1.969745	0.0640000	-0.0106	-17.4039	0.1053	135.5493	132.00531616
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2.031300	0.0660000	-0.0103	-16.8187	0.1031	134.3678	131.22657740
2.062277	0.0670000	-0.0097	-15.9424	0.1022	133.8496	130.84204064
2.092254	0.0680000	-0.0090	-14.7517	0.1013	133.3850	130.46152121
2.123232	0.0690000	-0.0082	-13.3445	0.1005	132.9727	130.08592073
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2.492559	0.0810000	-0.0140	-22.9714	0.0904	127.5129	126.19517697
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3.046550	0.0990000	-0.3103	-507.5151	0.0236	91.5341	121.90966125
3.077727	0.1000000	-0.5140	-840.8277	-0.0296	62.8478	121.55763221

TEMPERATURE AND HEAT FLUX  
PREDICTED BY HSAI-YIN LEE

(NORM.TM) (REAL.TM) (NUR.FLUX) (REAL.FLUX) (NOR. ) (REAL.SUF.T) (TH.CO.?)  
(DMLESS) (SEC.) (DMLESS) (BTU/FTSQ SEC) (DMLESS) (DEG.F.) (DEG.F.)

M60 GUN THERMOCOUPLE 4 9.0 INCHES FROM BREECH

BEST AVAILABLE COPY

M60 GUN THERMOCOUPLE 10 21 INCHES FROM BORE  
 BORE SURFACE TEMPERATURE AND HEAT FLUX PROGRAM  
 NUMBER OF B(I) COFF. TO BE FITTED = 10

TIME TEMPERATURE

0.0100000000	126.2000000000
0.0200000000	131.6000000000
0.0300000000	124.2000000000
0.0400000000	119.6000000000
0.0500000000	116.8000000000
0.0600000000	113.3000000000
0.0700000000	109.8000000000
0.0800000000	107.1000000000
0.0900000000	105.7000000000
0.1000000000	104.2000000000

TIME OF DATA SHIFTED BY(SEC) -0.002000

BORE RADIUS (FT.) = 0.01250

CUTER RADIUS (FT.)= 0.03562

BORE TO THERMOCOUPLE DISTANCE (FT)= 0.001670

INITIAL THERMOCOUPLE TEMPERATURE (F.)= 78.8000

INITIAL GAS TEMPERATURE (F.)= 4.4937

TERMAL DIFFUSIVITY (FTSQ/SEC) = 0.00010307

TERMAL CONDUCTIVITY(BTU/FT SEC F.)= 0.00555555

NUMBER OF TIME TEMPERATURE PAIRS (SEC.,F.)=10

NUMBER OF B(I) COFF. TO BE FITTED = 10

TIME TEMPERATURE

0.0120000000	126.2000000000
0.0220000000	131.6000000000
0.0320000000	124.2000000000
0.0420000000	119.6000000000
0.0520000000	116.8000000000
0.0620000000	113.3000000000
0.0720000000	109.8000000000
0.0820000000	107.1000000000
0.0920000000	105.7000000000
0.1020000000	104.2000000000

TIME OF DATA SHIFTED BY(SEC) -0.002000

COEFFICIENTS OF B(I)= 0.49358586D 01 I = 1

COEFFICIENTS OF B(I)= -0.23981321D 02 I = 2

COEFFICIENTS OF B(I)= 0.53463543D 02 I = 3

COEFFICIENTS OF B(I)= -0.69864894D 02 I = 4

COEFFICIENTS OF B(I)= 0.57893886D 02 I = 5

COEFFICIENTS OF B(I)= -0.31318033D 02 I = 6

COEFFICIENTS OF B(I)= 0.11016604D 02 I = 7

COEFFICIENTS OF B(I)= -0.24257631D 01 I = 8

COEFFICIENTS OF B(I)= 0.30329015D 00 I = 9

COEFFICIENTS OF B(I)= -0.16413678D-01 I = 10

## APPENDIX D THE CASE OF OSCILATORY SURFACE TEMPERATURE

Consider a slab with a sufficient thickness,  $\ell$ , such that when a surface is subjected to a periodic surface temperature variation with a frequency  $w$  the other surface is held at the initial temperature  $T_0$ . If properties are assumed constant the governing equation for the problem can be written as

$$\frac{\partial v}{\partial t} = \alpha \frac{\partial^2 v}{\partial x^2} \quad (1)$$

with initial and boundary conditions as

$$v(x, 0) = 0 \quad (2)$$

$$v(\ell, t) = \sin wt \quad (3)$$

$$v(0, t) = 0 \quad (4)$$

where  $v = (T - T_0)/(T_{\max} - T_0)$  and  $\alpha$  is the thermal diffusivity.

The solution of the problem according to Carslaw and Jaeger [4] can be written as

$$v = \frac{2\alpha\pi}{\ell^2} \sum_{n=1}^{\infty} (-1)(-1)^n n e^{-\alpha n^2 \pi^2 \tau / \ell^2} \sin\left(\frac{n\pi x}{\ell}\right) \int_0^t e^{\alpha n^2 \pi^2 \lambda / \ell^2} \sin w\lambda d\lambda \quad (5)$$

with

$$\int e^{ax} \sin bx dx = e^{ax} (a \sin bx - b \cos bx) / (a^2 + b^2) \quad (6)$$

we have

$$v = 2\alpha\pi \sum_{n=1}^{\infty} (-1)(-1)^n n \left[ (\alpha n^2 \pi^2 \sin \omega t - \omega^2 \cos \omega t) + \omega^2 e^{-\alpha n^2 \pi^2 t/\ell^2} \right] \cdot \sin \left( \frac{n\pi x}{\ell} \right) / [\alpha^2 n^4 \pi^4 + \omega^2 \ell^2] \quad (7)$$

PART III

PREDICTION OF TRANSIENT SURFACE HEAT FLUX AND  
TEMPERATURE ON A HOLLOW CYLINDER

## INTRODUCTION

In the study of transient heat transfer many efforts have been made on the so-called "inverse problem" [1,2] where a surface heat flux and temperature is to be predicted by the measured data at some location interior to a body.

In the previous works [1-6] the solution is represented in either an integral form after some manipulation of the contour integral from the inverse transform, or in a series form after the expansion of the solution for small and large times. Using Laplace transformation Chen and Thomsen [6] introduced a polynomial in terms of the error function to represent the response of thermocouple measurement and the inversion is accomplished for any transient surface heat flux at the inner surface of a cylindrical tube. However, their inversion solution is valid only for a short duration due to the asymptotic expansion of the modified Bessel function in the inverse Laplace transform. In this study an exact solution obtained from the inverse Laplace transform by the convolution method is given for the case of hollow cylinder. The solution is valid for both constant and variable heat flux and for both short and long time duration.

## ANALYSIS

Consider a long hollow cylinder with sufficient wall thickness such that the outer surface temperature has a negligible response when the inner surface is exposed to a thermal pulse of a transient process. This condition considerably simplifies the theoretical analysis as the outer boundary may be assumed to be infinite, and only one interior probe of the cylinder

is required in the experimental measurement. The material of the cylinder is considered to be homogeneous and isotropic with constant thermal diffusivity,  $\alpha$ . Let  $R_i$  and  $R_o$  be, respectively, the inner and outer surface radii.  $R_i$  the radius of the probe location and  $t$  the dimensionless time. If the temperature of the cylinder is initially uniform at  $T_o$ , the mathematical problem governing the temperature  $T$ , may be written as

$$\frac{\partial \theta}{\partial t} = \frac{\partial^2 \theta}{\partial r^2} + \frac{1}{r} \frac{\partial \theta}{\partial r} \quad 1 < r < r_o = \infty \quad (1)$$

$$\theta(r, 0) = 0 \quad (2)$$

$$\theta(\infty, t) = 0 \quad (3)$$

$$\theta(r_i, t) = f(t) \quad 1 < r_i < \infty \quad (4)$$

where  $\theta = T - T_o$ ,  $r = R/R_i$ ,  $t = \alpha r^2/R_i^2$ , and  $f(t)$  is the interior temperature response of the thermocouple measured at  $r = r_i$  at the dimensionless time  $t$ . The problem is to predict the surface temperature  $\theta(1, t)$  and heat flux per unit area

$$q = - (K/R_i) (\partial \theta / \partial r) \Big|_{r=1} \quad (5)$$

where  $K$  is the thermal conductivity.

The problem can be solved by Laplace transformation. Let the transformation be

$$\bar{\theta}(r, s) = \int_0^{\infty} \theta e^{-ts} dt \quad (6)$$

When  $\theta$  satisfies the Dirichlet's condition the temperature function  $\theta$  is recovered by inversion of the Laplace transformation as

$$\theta(r, t) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \theta e^{st} ds \quad (7)$$

where  $c$  is a suitable positive value. Equation (1) and (2) under transformation (6) becomes

$$\frac{d^2\bar{\theta}}{dr^2} + \frac{1}{r} \frac{d\bar{\theta}}{dr} = s\bar{\theta}$$

which has a solution of the form

$$\bar{\theta} = A I_0(pr) + B K_0(pr) \quad (8)$$

where  $I_0$  and  $K_0$  are modified Bessel functions of the first and second kind with  $p = (s)^{1/2}$ . With the boundary conditions (3) and (4), (8) becomes

$$\bar{\theta} = \bar{f}(s) [K_0(pr)/K_0(pr_1)] \quad (9)$$

where  $\bar{f}(s)$  is the Laplace transform of the boundary condition (4).

The temperature response measured at  $r = r_1$  can be expressed by a polynomial or numerous other suitable functions. In the present analysis, for reasons to be explained later,  $f(t)$  will be represented as

$$f(t) = \sum_{n=1}^N b_n \int_0^t F_1(\tau) F_n(t - \tau) d\tau \quad (10)$$

If we choose  $F_1(t) = \frac{1}{2\tau} e^{-\frac{r_1^2}{4\tau}}$  and  $F_n(t - \tau)$  being any arbitrary function

depending on  $n$ , for example  $(t - \tau)^n$  etc. then the Laplace transform of Eq. (10) gives

$$\bar{f}(s) = \sum_{n=1}^N b_n K_o(pr_1) \bar{F}_n(s) \quad (11)$$

where  $\bar{F}_n(s)$  is the Laplace transform of  $F_n(t)$ . Substituting Eq. (11) into Eq. (19) we have

$$\bar{\theta}(s, r) = \sum_{n=1}^N b_n K_o(pr) \bar{F}_n(s) \quad (12)$$

It is noted that the  $K_o(pr_1)$  in Eq. (9) has been cancelled by this substitution which explains the choice of  $F(t) = \frac{1}{4K} e^{-\frac{r_1^2}{4K}}$  in Eq. (10).

The inversion of Eq. (12) gives

$$\theta(t, r) = \sum_{n=1}^N b_n \int_0^t \frac{1}{2\tau} e^{-\frac{r^2}{4\tau}} F_n(t - \tau) d\tau \quad (13)$$

At surface  $r = 1$

$$\theta(t, 1) = \sum_{n=1}^N b_n \int_0^t \frac{1}{2\tau} e^{-\frac{1}{4\tau}} F_n(t - \tau) d\tau \quad (14)$$

The temperature gradient and hence heat flux at surface is

$$\left. \frac{\partial \theta}{\partial r} \right|_{r=1} = - \sum_{n=1}^N b_n \int_0^t \frac{1}{4\tau^2} e^{-\frac{1}{4\tau}} F_n(t - \tau) d\tau \quad (15)$$

some examples of  $F_n(t - \tau)$  function and the representation of the thermo-couple response are:

$$\text{Case 1} \quad \text{If } \bar{F}_n(s) = \frac{1}{s^2 + n^2}$$

$$\text{then } f(t) = \sum_{n=1}^N b_n \int_0^t \frac{1}{2\tau} e^{-\frac{r_1^2}{4\tau}} \frac{1}{n} \sin n(t-\tau) d\tau \quad (16)$$

$$\text{Case 2} \quad \text{If } \bar{F}_n(s) = \frac{s}{s^2 + n^2}$$

$$\text{then } f(t) = \sum_{n=1}^N b_n \int_0^t \frac{1}{2\tau} e^{-\frac{r_1^2}{4\tau}} \cos n(t-\tau) d\tau \quad (17)$$

$$\text{Case 3} \quad \text{If } \bar{F}_n(s) = \frac{1}{(s+a)^n}$$

$$\text{then } f(t) = \sum_{n=1}^N b_n \int_0^t \frac{e^{-\frac{r_1^2}{4\tau}} (t-\tau)^{n-1} e^{-a(t-\tau)}}{2\tau(n-1)} d\tau \quad (18)$$

$$\text{Case 4} \quad \text{If } \bar{F}_n(s) = s^{-(n+1/2)}$$

$$\text{then } f(t) = \sum_{n=1}^N b_n \int_0^t \frac{e^{-\frac{r_1^2}{4\tau}} 2^n (t-\tau)^{n-1/2}}{2\tau 1.3.5\dots(2n-1)\sqrt{\pi}} d\tau \quad (19)$$

$$\text{Case 5} \quad \text{If } \bar{F}_n(s) = \frac{1}{s^n}$$

$$\text{then } f(t) = \sum_{n=1}^N b_n \int_0^t \frac{e^{-\frac{r_1^2}{4\tau}}}{2} \frac{(t-\tau)^{n-1}}{(n-1)!} d\tau \quad (20)$$

## DISCUSSION

In practical application one normally has a measurement of an interior thermocouple response preferably near the heating surface. The temperature measured by the thermocouple can then be fitted with a suitable polynomial given in Eq. (16) through (20). From Chen and Thomsen's (6) work the minimum degree of polynomial N should be about 10 for a transient heating such as in gun barrel and engine. The coefficients  $b_n$  obtained from the fitting is then substituted into Eq. (14) for the surface temperature and Eq. (15) for the surface heat flux.

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port properties and size. The second is the refinement of the presently available computer program for pre-		dition of the surface temperature and heat flux at the (cont.) over	dition of the surface temperature and heat flux at the (cont.) over	dition of the surface temperature and heat flux at the (cont.) over		dition of the surface temperature and heat flux at the (cont.) over	dition of the surface temperature and heat flux at the (cont.) over
PIPE FLOW		PIPE FLOW	PIPE FLOW	PIPE FLOW		PIPE FLOW	PIPE FLOW

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